

OPTIMALITY OF REFRACTION STRATEGIES FOR A CONSTRAINED DIVIDEND PROBLEM

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Abstract

We consider de Finetti's problem for spectrally one-sided Lévy risk models with control strategies that are absolutely continuous with respect to the Lebesgue measure. Furthermore, we consider the version with a constraint on the time of ruin. To characterize the solution to the aforementioned models, we first solve the optimal dividend problem with a terminal value at ruin and show the optimality of threshold strategies. Next, we introduce the dual Lagrangian problem and show that the complementary slackness conditions are satisfied, characterizing the optimal Lagrange multiplier. Finally, we illustrate our findings with a series of numerical examples.

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1. Introduction

In de Finetti's optimal dividend problem, the aim is to maximize the total expected discounted dividends accumulated until ruin. Intuitively, as the risk of ruin must be considered, dividends should be paid only when there is a sufficient amount of surplus available. With this conjecture and under the assumption of stationary increments of the underlying process (in the Lévy cases), the optimality of a barrier strategy that pays out any amount above a certain barrier has been pursued in various papers. Because the resulting controlled process becomes a classical reflected process, existing fluctuation theoretical results have been efficiently applied to solve explicitly the problem, at least under suitable conditions. See, among others, Avram *et al.* [2] for the spectrally negative case and Bayraktar *et al.* [4] for the spectrally positive (dual) case.

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Despite these important works, there are several disputes about the classical model in the sense that the set of admissible strategies is too large and contains those that are in reality impossible to implement. For example, under the barrier strategy that is shown to be optimal, ruin must occur almost surely, and this is rather an undesirable conclusion. In addition, companies in reality cannot focus on dividend maximization and instead need to take into account the impacts of bankruptcy to themselves and also to the market. For these reasons, in the past decade, several extensions have been considered so as to obtain a more realistic model, by considering more restricted sets of admissible strategies.

Motivated by these, in this paper we focus on the model with the absolutely continuous condition on the dividend strategy and additional condition on the time of ruin. We consider both cases driven by spectrally negative and positive Lévy processes.

Regarding the absolutely continuous condition, it is assumed that the rate at which dividends are paid is bounded. More specifically, the dividend strategy must be absolutely continuous with respect to the Lebesgue measure with its density bounded by a given constant. Under some assumptions on the jump measure (e.g. the completely monotone Lévy density assumption as in the current paper), analogously to the barrier strategy that is optimal in the classical case, the threshold strategy—that pays dividends at the maximal rate as long as the surplus is above a certain fixed level—is optimal in this case. For a spectrally negative Lévy surplus process, Kyprianou *et al.* [12] showed the optimality of the threshold strategy under a completely monotone assumption on the Lévy measure. The spectrally positive Lévy model has been solved by Yin *et al.* [22]. In both cases, the optimally controlled process becomes the *refracted Lévy process* of Kyprianou and Loeffen [11], and the fluctuation identities for this process can be used efficiently to solve de Finetti's problem under the absolutely continuous condition.

Following the recent work of Hernández *et al.* [8], we study the case in which the longevity feature is added to the problem by considering a constraint on the time of ruin. The longevity aspect of the firm remained as a separate problem; see [20] for a survey on this matter. Despite efforts to integrate both features [6], [17], [21], it was not until very recently a successful solution to a model that actually accounts for the trade-off between performance and longevity was presented. Hernández and Junca [7] considered de Finetti's problem in the setting of Cramér–Lundberg reserves with i.i.d. exponentially distributed jumps adding a constraint to the expected time of ruin of the firm.

The contribution of this paper is twofold.

1. We first solve the optimal dividend problem *with a terminal value at ruin* under the absolutely continuous assumption. We solve this problem for the spectrally negative Lévy case under the assumption that the Lévy measure has a completely monotone density and also for the general spectrally positive Lévy case. In both models we show that a threshold strategy is optimal (see Theorems 4.1 and 6.1). The optimal refraction level as well as the value function are concisely expressed in terms of the scale function. Its optimality is confirmed by a verification lemma.
2. We then use these results to solve the constrained dividend maximization problem over the set of strategies that satisfy a particular constraint on the ruin time. This is an extension of [8] under the absolutely continuous assumption. Theorem 5.1 shows the result when the reserves are modeled by a spectrally negative Lévy process with a completely monotone Lévy density and Theorem 6.2 for the general dual model.

The rest of the paper is organized as follows. In Section 2, we formulate the problem. In Section 3, we present an overview of scale functions and some fluctuation identities related to spectrally negative Lévy processes and their respective refracted processes. In Section 4,

we solve the optimal dividend problem with a terminal cost and the absolutely continuous assumption for the case of a spectrally negative Lévy process with a completely monotone Lévy density. In Section 5, we extend the results to solve the constrained dividends problem. In Section 6, we solve the same problems for the spectrally positive case. Finally, in Section 7, we give some numerical results.

2. Formulation of the problem

In this section, we formulate the constrained de Finetti problem driven by a spectrally negative Lévy process. The spectrally positive Lévy process is its dual and a slight modification is only needed to formulate the spectrally positive case (see Section 6).

2.1. Spectrally negative Lévy processes

Recall that a spectrally negative Lévy process is a stochastic process, which has càdlàg paths and stationary and independent increments such that there are no positive discontinuities. To avoid degenerate cases in the forthcoming discussion, we shall additionally exclude from this definition the case of monotone paths. This means that we are not interested in the case of a deterministic increasing linear drift or the negative of a subordinator. Henceforth, we assume that $X = \{X_t : t \geq 0\}$ is a spectrally negative Lévy process defined on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ with Lévy triplet given by (γ, σ, Π) , where $\gamma \in \mathbb{R}$, $\sigma \geq 0$, and Π is a measure concentrated on $(0, \infty)$ satisfying

$$\int_{(0, \infty)} (1 \wedge z^2) \Pi(dz) < \infty.$$

The Laplace exponent of X is given by

$$\psi(\lambda) = \log \mathbb{E}[e^{\lambda X_1}] = \gamma\lambda + \frac{1}{2}\sigma^2\lambda^2 - \int_{(0, \infty)} (1 - e^{-\lambda z} - \lambda z \mathbf{1}_{\{0 < z \leq 1\}}) \Pi(dz), \tag{2.1}$$

which is well defined for $\lambda \geq 0$. Here \mathbb{E} denotes the expectation with respect to \mathbb{P} . In addition, the reader should note that, for convenience, we have arranged the representation of the Laplace exponent in such a way that the support of the Lévy measure is positive even though the process experiences only negative jumps. As a strong Markov process, we shall endow X with probabilities $\{\mathbb{P}_x : x \in \mathbb{R}\}$ such that, under \mathbb{P}_x , we have $X_0 = x$ with probability 1. Note that $\mathbb{P}_0 = \mathbb{P}$ and $\mathbb{E}_0 = \mathbb{E}$.

The spectrally negative Lévy process X has paths of bounded variation if and only if $\sigma = 0$ and $\int_{(0, 1]} z \Pi(dz) < \infty$, in which case, X can be written as

$$X_t = ct - S_t, \quad t \geq 0,$$

where $c := \gamma + \int_{(0, 1]} z \Pi(dz)$ and $\{S_t : t \geq 0\}$ is a drift-less subordinator. Note that we must have $c > 0$, since it is assumed that X does not have monotone paths.

In the classical model of the wealth of an insurance company, premium is received at a constant rate c whereas F -distributed claims arrive at the jump times of a Poisson process with rate μ . The corresponding surplus process is called the Cramér–Lundberg risk process and is a special case of the spectrally negative Lévy process with $\gamma = c - \mu \int_{(0, 1]} z F(dz)$, $\sigma = 0$, and $\Pi = \mu F$.

2.2. Admissible strategies

Let $D = \{D_t : t \geq 0\}$ be a dividend strategy, meaning that it is a nonnegative and nondecreasing process adapted to the completed and right continuous filtration $\mathbb{F} := \{\mathcal{F}_t : t \geq 0\}$ of X . Here, for each fixed $t \geq 0$, the quantity D_t represents the cumulative dividends paid

out up to time t by the insurance company whose risk process is modeled by X . The controlled Lévy process becomes $U^D = \{U_t^D = X_t - D_t : t \geq 0\}$ and we write

$$\tau^D := \inf\{t > 0 : U_t^D < 0\}$$

for the time at which ruin occurs when the dividend payments are taken into account.

In this work, we are interested in adding a constraint to the dividend processes. Specifically, we will only work with absolutely continuous strategies of bounded rate, i.e.

$$D_t = \int_0^t d(s) \, ds, \quad t \geq 0,$$

such that the *dividend rate* d satisfies $0 \leq d(t) \leq \delta$ for $t \geq 0$, where $\delta > 0$ is a ceiling rate. We will denote by Θ the family of admissible strategies satisfying the conditions mentioned above.

2.3. Constrained de Finetti's problem and its dual

The expected net present value under the dividend policy $D \in \Theta$ with discounting at rate $q > 0$ and initial capital $x \geq 0$ is given by

$$v^D(x) = \mathbb{E}_x \left[\int_0^{\tau^D} e^{-qt} \, dD_t \right].$$

The dividend problem, originally considered by de Finetti, asks to maximize the expected net present value of dividend payments over the set of strategies Θ .

Now, as studied in [8], we are interested in addressing a modification of this problem by adding a restriction to the dividend process D , which is given by the constraint

$$\mathbb{E}_x[e^{-q\tau^D}] \leq K, \quad 0 \leq K \leq 1 \text{ fixed.}$$

Strategies in Θ satisfying this constraint are called *feasible*, and are called *infeasible* otherwise.

Remark 2.1. By this constraint, we consider only strategies that avoid early ruin of the company. The classical problem corresponds to the case with $K = 1$. By selecting the constraint written in terms of the expectation of $e^{-q\tau^D}$, the Lagrangian subproblem (considered in Section 5) can be solved explicitly using the scale function and hence the constrained case as well. It is also of interest to consider other constraints, e.g. $\mathbb{E}_x[e^{-p\tau^D}] \leq K$ for $p \neq q$ and $\mathbb{E}_x[\tau^D] \leq K$. However, in these cases, the techniques in this paper cannot be directly applied, and the optimality of a threshold strategy may not hold.

We want to maximize the expected net present value of dividend payments over the set of feasible strategies. That is, we aim to solve the optimization problem, for $x \geq 0$ and $0 \leq K \leq 1$,

$$V(x; K) := \sup_{D \in \Theta} v^D(x) \quad \text{such that} \quad \mathbb{E}_x[e^{-q\tau^D}] \leq K, \quad (2.2)$$

where, in the case $\mathbb{E}_x[e^{-q\tau^D}] > K$ for all $D \in \Theta$, we set $V(x; K) = -\infty$ and call problem (2.2) infeasible.

Proceeding as in [8], we use Lagrange multipliers to reformulate the problem. For $\Lambda \geq 0$, we define the function

$$v_\Lambda^D(x; K) := v^D(x) + \Lambda(K - \mathbb{E}_x[e^{-q\tau^D}]). \quad (2.3)$$

Note that we can write problem (2.2) as $V(x; K) = \sup_{D \in \Theta} \inf_{\Lambda \geq 0} v_{\Lambda}^D(x; K)$ since any infeasible strategy D will make $\inf_{\Lambda \geq 0} v_{\Lambda}^D(x; K) = -\infty$, and any feasible strategy D will make $\inf_{\Lambda \geq 0} v_{\Lambda}^D(x; K) = v_0^D(x; K) = v^D(x)$.

The dual problem of (2.2) is obtained by interchanging the sup with the inf in the expression above, yielding an upper bound,

$$V(x; K) = \sup_{D \in \Theta} \inf_{\Lambda \geq 0} v_{\Lambda}^D(x; K) \leq \inf_{\Lambda \geq 0} V_{\Lambda}(x; K), \tag{2.4}$$

where

$$V_{\Lambda}(x; K) := \sup_{D \in \Theta} v_{\Lambda}^D(x; K). \tag{2.5}$$

Therefore, the main goal is to prove that $V(x; K) = \inf_{\Lambda \geq 0} V_{\Lambda}(x; K)$ and to find, if it exists, an optimal Λ (Lagrange multiplier) with which the infimum is attained. In order to do this, we will first solve (2.5).

We remark that if we set

$$V_{\Lambda}(x) := V_{\Lambda}(x; 0) \quad \text{and} \quad v_{\Lambda}^D(x) := v_{\Lambda}^D(x; 0), \quad D \in \Theta, \tag{2.6}$$

then $v_{\Lambda}^D(x; K) = v_{\Lambda}^D(x) + \Lambda K$ and hence, $V_{\Lambda}(x; K) = V_{\Lambda}(x) + \Lambda K$. Therefore, solving (2.5) is equivalent to solving

$$V_{\Lambda}(x) := \sup_{D \in \Theta} v_{\Lambda}^D(x). \tag{2.7}$$

3. Review of scale functions

In this section, we review the scale function of spectrally negative Lévy processes. First, we define the process $Y = \{Y_t = X_t - \delta t : t \geq 0\}$ with its Laplace exponent

$$\psi_Y(\theta) := \psi(\theta) - \delta\theta, \quad \theta \geq 0. \tag{3.1}$$

We assume here that Y is a spectral negative Lévy process and not the negative of a subordinator (see Assumption 4.3).

Fix $q > 0$. Following the same notation as in [11], we use $W^{(q)}$ and $\mathbb{W}^{(q)}$ for the scale functions of X and Y , respectively. These are the mappings from \mathbb{R} to $[0, \infty)$ that are 0 on the negative half-line, while on the positive half-line they are strictly increasing functions that are defined by their Laplace transforms,

$$\begin{aligned} \int_0^{\infty} e^{-\theta x} W^{(q)}(x) dx &= \frac{1}{\psi(\theta) - q}, & \theta > \Phi(q), \\ \int_0^{\infty} e^{-\theta x} \mathbb{W}^{(q)}(x) dx &= \frac{1}{\psi_Y(\theta) - q}, & \theta > \varphi(q), \end{aligned} \tag{3.2}$$

where

$$\Phi(q) := \sup\{\lambda \geq 0 : \psi(\lambda) = q\} \quad \text{and} \quad \varphi(q) := \sup\{\lambda \geq 0 : \psi_Y(\lambda) = q\}. \tag{3.3}$$

By the strict convexity of ψ on $[0, \infty)$, we derive the strict inequality $\varphi(q) > \Phi(q) > 0$.

We also define, for $x \in \mathbb{R}$,

$$\begin{aligned}\bar{W}^{(q)}(x) &:= \int_0^x W^{(q)}(y) \, dy, \\ Z^{(q)}(x) &:= 1 + q\bar{W}^{(q)}(x), \\ \bar{Z}^{(q)}(x) &:= \int_0^x Z^{(q)}(z) \, dz = x + q \int_0^x \int_0^z W^{(q)}(w) \, dw \, dz.\end{aligned}$$

Noting that $W^{(q)}(x) = 0$ for $-\infty < x < 0$, we have

$$\bar{W}^{(q)}(x) = 0, \quad Z^{(q)}(x) = 1, \quad \text{and} \quad \bar{Z}^{(q)}(x) = x, \quad x \leq 0.$$

In addition, we define $\bar{\mathbb{W}}^{(q)}$, $\mathbb{Z}^{(q)}$, and $\bar{\mathbb{Z}}^{(q)}$ analogously for Y . The scale functions of X and Y are related for $x \in \mathbb{R}$ by the equality

$$\delta \int_0^x \mathbb{W}^{(q)}(x-y)W^{(q)}(y) \, dy = \bar{\mathbb{W}}^{(q)}(x) - \bar{W}^{(q)}(x). \quad (3.4)$$

This can be confirmed by showing that the Laplace transforms on both sides are equal.

Regarding their asymptotic behaviors, we have, as in Lemmas 3.1 and 3.2 of [9],

$$\begin{aligned}W^{(q)}(0) &= \begin{cases} 0 & \text{if } X \text{ is of unbounded variation,} \\ c^{-1} & \text{if } X \text{ is of bounded variation,} \end{cases} \\ \mathbb{W}^{(q)}(0) &= \begin{cases} 0 & \text{if } Y \text{ is of unbounded variation,} \\ (c-\delta)^{-1} & \text{if } Y \text{ is of bounded variation,} \end{cases}\end{aligned} \quad (3.5)$$

and

$$\begin{aligned}W^{(q)'}(0+) &:= \lim_{x \downarrow 0} W^{(q)'}(x) = \begin{cases} \frac{2}{\sigma^2} & \text{if } \sigma > 0, \\ \infty & \text{if } \sigma = 0 \text{ and } \Pi(0, \infty) = \infty, \\ \frac{q + \Pi(0, \infty)}{c^2} & \text{if } \sigma = 0 \text{ and } \Pi(0, \infty) < \infty, \end{cases} \\ \mathbb{W}^{(q)'}(0+) &:= \lim_{x \downarrow 0} \mathbb{W}^{(q)'}(x) = \begin{cases} \frac{2}{\sigma^2} & \text{if } \sigma > 0, \\ \infty & \text{if } \sigma = 0 \text{ and } \Pi(0, \infty) = \infty, \\ \frac{q + \Pi(0, \infty)}{(c-\delta)^2} & \text{if } \sigma = 0 \text{ and } \Pi(0, \infty) < \infty. \end{cases}\end{aligned} \quad (3.6)$$

On the other hand, as in Lemma 3.3 of [13], as $x \rightarrow \infty$,

$$e^{-\Phi(q)x}W^{(q)}(x) \nearrow \psi'(\Phi(q))^{-1} \quad \text{and} \quad e^{-\varphi(q)x}\mathbb{W}^{(q)}(x) \nearrow \psi'_Y(\varphi(q))^{-1}. \quad (3.7)$$

4. Optimal dividend problem with a terminal value

In this section, we solve problem (2.7). The obtained results will be applied to the constrained case in the next section. In this and next sections where we deal with the spectrally negative case, we assume the following.

Assumption 4.1. The Lévy measure Π of the process X has a completely monotone density. That is, Π admits a density π whose n th derivative $\pi^{(n)}$ exists for all $n \geq 1$ and satisfies

$$(-1)^n \pi^{(n)}(x) \geq 0, \quad x > 0.$$

Remark 4.1. Under Assumption 4.1, the scale functions $W^{(q)}$ and $\mathbb{W}^{(q)}$ defined in Section 3 are infinitely continuously differentiable on $(0, \infty)$. For more details, see Lemma A.1.

This assumption is known to be a sufficient optimality condition for threshold strategies in the classical spectrally negative case by Loeffen [15], for the absolutely continuous case (with $\Lambda = 0$) by Kyprianou *et al.* [12], and for the periodic case by Noba *et al.* [16] (with $\Lambda = 0$).

In this section, we allow Λ to be negative (in which case a positive payoff is collected at ruin time), but need to assume the following in order to avoid the trivial case (see Remark 4.2).

Assumption 4.2. We assume that $q\Lambda + \delta > 0$.

Remark 4.2. Suppose that Assumption 4.2 does not hold (i.e. $q\Lambda + \delta \leq 0$). Because the dividend rate is bounded by δ , for any dividend policy $D \in \Theta$ and $x \geq 0$,

$$\begin{aligned} v_\Lambda^D(x) &\leq \delta \mathbb{E}_x \left[\int_0^{\tau^D} e^{-qt} dt \right] - \Lambda \mathbb{E}_x [e^{-q\tau^D}] \\ &\leq \frac{\delta}{q} - \frac{q\Lambda + \delta}{q} \mathbb{E}_x [e^{-q\tau^D}] \\ &\leq \frac{\delta}{q} - \frac{q\Lambda + \delta}{q} \sup_{D' \in \Theta} \mathbb{E}_x [e^{-q\tau^{D'}}]. \end{aligned}$$

This implies that v_Λ^D is maximized by taking the strategy that pays dividends at the ceiling rate δ for all $t \geq 0$ because it maximizes $\mathbb{E}_x [e^{-q\tau^{D'}}]$ over Θ .

Finally, we make the following assumption.

Assumption 4.3. If X has paths of bounded variation then $\delta < c$.

This assumption is commonly used in the literature (see, for instance, [12] and [18])—without this assumption, the techniques used in this paper are not capable of giving a complete solution.

For the case when Assumption 4.3 is not satisfied, as long as the starting value is below the barrier, a reflection strategy—which is optimal when the absolutely continuous condition is relaxed—is feasible in the considered problem and is therefore optimal as well. On the other hand, when the starting value is above the barrier, the optimal (reflection) strategy in the classical case pushes the process instantaneously to the barrier, whereas this is not feasible in our problem setting. The optimal strategy in this case cannot be directly obtained by the methods used in this paper, and it is out of scope of this paper. The reader is referred to Azcue and Muler [3] for complete results for the Cramér–Lundberg case when Assumption 4.3 is violated in a related problem.

4.1. Threshold strategies

The objective of this section is to show that the optimal strategies for (2.7) are of the threshold type. Under the threshold strategy D^b for $b \geq 0$, the resulting controlled process U^b is known as a refracted Lévy process of [11], which is the unique strong solution to

$$U_t^b = X_t - D_t^b, \quad \text{where} \quad D_t^b := \delta \int_0^t \mathbf{1}_{\{U_s^b > b\}} ds, \quad t \geq 0.$$

Let its ruin time be

$$\tau_b := \inf\{t > 0: U_t^b < 0\}.$$

The next identities are lifted from Theorems 5(ii) and 6(ii) of [11]. For $x \in \mathbb{R}$ and $b \geq 0$, we have

$$\mathbb{E}_x \left[\int_0^{\tau_b} e^{-qt} dD_t^b \right] = -\delta \overline{\mathbb{W}}^{(q)}(x-b) + \frac{1}{h(b)} \left(W^{(q)}(x) + \delta \int_b^x \mathbb{W}^{(q)}(x-y) W^{(q)'}(y) dy \right) \quad (4.1)$$

and

$$\begin{aligned} \Psi_x(b) &:= \mathbb{E}_x[e^{-q\tau_b}] \\ &= Z^{(q)}(x) + \delta q \int_b^x \mathbb{W}^{(q)}(x-y) W^{(q)}(y) dy \\ &\quad - \frac{q\varphi(q)e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)}(y) dy}{h(b)} \left(W^{(q)}(x) + \delta \int_b^x \mathbb{W}^{(q)}(x-y) W^{(q)'}(y) dy \right), \end{aligned} \quad (4.2)$$

where

$$h(b) := \varphi(q)e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)'}(y) dy. \quad (4.3)$$

Under the threshold strategy D^b , the expected net present value is denoted by

$$v_\Lambda^b(x) := \mathbb{E}_x \left[\int_0^{\tau_b} e^{-qt} dD_t^b \right] - \Lambda \Psi_x(b) \quad \text{for } x \geq 0. \quad (4.4)$$

Using (4.1) and (4.2), we have the following result.

Proposition 4.1. *The function v_Λ^b , with $b \geq 0$, is given by*

$$\begin{aligned} v_\Lambda^b(x) &= \xi_\Lambda(b) \left(W^{(q)}(x) + \delta \int_b^x \mathbb{W}^{(q)}(x-y) W^{(q)'}(y) dy \right) \\ &\quad - \Lambda \left(Z^{(q)}(x) + \delta q \int_b^x \mathbb{W}^{(q)}(x-y) W^{(q)}(y) dy \right) - \delta \overline{\mathbb{W}}^{(q)}(x-b) \quad \text{for } x \geq 0, \end{aligned} \quad (4.5)$$

where

$$\xi_\Lambda(b) := \frac{1}{h(b)} \left(1 + \varphi(q)q\Lambda e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)}(y) dy \right). \quad (4.6)$$

In particular, for $x \leq b$, we have

$$v_\Lambda^b(x) = \xi_\Lambda(b) W^{(q)}(x) - \Lambda Z^{(q)}(x). \quad (4.7)$$

Remark 4.3. From (4.3) and integration by parts,

$$\varphi(q)e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)}(y) dy = W^{(q)}(b) + \frac{h(b)}{\varphi(q)}, \quad b \geq 0. \quad (4.8)$$

Hence, the function ξ_Λ , given in (4.6), can be rewritten as

$$\xi_\Lambda(b) = \frac{1}{h(b)} \left(1 + q\Lambda \left(W^{(q)}(b) + \frac{h(b)}{\varphi(q)} \right) \right), \quad b \geq 0. \tag{4.9}$$

In particular, for the case $b = 0$, these expressions can be simplified as follows; the proof is deferred to Appendix A.1.

Lemma 4.1. *We have*

$$\begin{aligned} h(0) &= \varphi(q) \int_0^\infty e^{-\varphi(q)y} W^{(q)'}(y) dy = \varphi(q)(-W^{(q)}(0) + \delta^{-1}), \\ \xi_\Lambda(0) &= \frac{1}{h(0)} \left(1 + \frac{q\Lambda}{\delta} \right) = \frac{\delta + q\Lambda}{\varphi(q)(1 - \delta W^{(q)}(0))}, \end{aligned} \tag{4.10}$$

and, for $x \geq 0$,

$$v_\Lambda^0(x) = \xi_\Lambda(0)\mathbb{W}^{(q)}(x)(1 - \delta W^{(q)}(0)) - \Lambda \mathbb{Z}^{(q)}(x) - \delta \overline{\mathbb{W}}^{(q)}(x). \tag{4.11}$$

4.2. Selection of the optimal threshold b_Λ

Focusing on the set of threshold strategies, we now select our candidate optimal threshold, which we call b_Λ . In view of (4.7), such b_Λ must maximize ξ_Λ . Motivated by this fact, we pursue b_Λ such that $\xi'_\Lambda(b_\Lambda)$ vanishes if such a value exists.

First, we rewrite the form of $\xi'_\Lambda(b)$ as follows. Fix $b > 0$. Taking a derivative in (4.9) and using the fact that $h'(b) = \varphi(q)(h(b) - W^{(q)'}(b))$ (by (4.3)),

$$\begin{aligned} \xi'_\Lambda(b) &= q\Lambda - \frac{h'(b)}{h(b)} \xi_\Lambda(b) \\ &= q\Lambda - \frac{\varphi(q)}{h(b)} \left(1 + q\Lambda \left(W^{(q)}(b) + \frac{h(b)}{\varphi(q)} \right) \right) + \frac{\varphi(q)W^{(q)'}(b)}{h(b)} \xi_\Lambda(b) \\ &= \frac{\varphi(q)W^{(q)'}(b)}{h(b)} (\xi_\Lambda(b) - g_\Lambda(b)), \end{aligned} \tag{4.12}$$

where

$$g_\Lambda(b) := \frac{1 + q\Lambda W^{(q)}(b)}{W^{(q)'}(b)}. \tag{4.13}$$

In view of (4.12), we now define the (candidate) optimal threshold level for (2.7) by

$$b_\Lambda := \inf\{b > 0 : \xi'_\Lambda(b) \leq 0\} = \inf\{b > 0 : \xi_\Lambda(b) - g_\Lambda(b) \leq 0\}. \tag{4.14}$$

Here, we set $\inf \emptyset = \infty$ for convenience, but we will see in Proposition 4.2 that b_Λ is necessarily finite.

Remark 4.4. Following the proof of Lemma 3 of [12], the function h as in (4.3) has the following properties.

- (i) As a special case with $\Lambda = 0$, b_0 is the point where $b \mapsto h(b)$ attains its global minimum. Hence, $h'(b) < 0$ for $b < b_0$ and $h'(b) > 0$ for $b > b_0$.
- (ii) We have $\lim_{b \rightarrow \infty} h(b) = \infty$.

Remark 4.5. Note that the function g_Λ plays a key role in [14] and satisfies the following assertions.

- (i) $g_\Lambda(0+) = (1 + q\Lambda W^{(q)}(0))/W^{(q)'(0+)}$ and $g_\Lambda(b) \rightarrow q\Lambda/\Phi(q)$ as $b \rightarrow \infty$ by (3.7).
- (ii) If we define

$$a_\Lambda := \sup\{b \geq 0: g_\Lambda(b) \geq g_\Lambda(x) \text{ for all } x \geq 0\},$$

we know that a_Λ is finite (see [14, Proposition 3]) and is the unique point where g_Λ has a global maximum; see [14, proof of Theorem 1]. Moreover if $a_\Lambda = 0$ then $g'_\Lambda(b) < 0$ for $b \in (0, \infty)$, and if $a_\Lambda > 0$ then $g'_\Lambda(a_\Lambda) = 0$, $g'_\Lambda(b) > 0$ for $b < a_\Lambda$ and $g'_\Lambda(b) < 0$ for $b > a_\Lambda$.

We will now prove an auxiliary result which describes the asymptotic behavior of the function ξ_Λ .

Lemma 4.2. *We have*

$$\lim_{b \rightarrow \infty} \xi_\Lambda(b) = \lim_{b \rightarrow \infty} g_\Lambda(b) = \frac{q\Lambda}{\Phi(q)}. \tag{4.15}$$

Proof. Recall from Remark 4.5 the convergence of g_Λ . Now, letting $b \rightarrow \infty$ in (4.9), we observe that

$$\lim_{b \rightarrow \infty} \xi_\Lambda(b) = q\Lambda \left(\frac{1}{\varphi(q)} + \lim_{b \rightarrow \infty} \frac{W^{(q)}(b)}{h(b)} \right), \tag{4.16}$$

since $h(b) \rightarrow \infty$ as $b \rightarrow \infty$. On the other hand, by the dominated convergence theorem and using (3.7), it follows that

$$\begin{aligned} \frac{W^{(q)}(b)}{h(b)} &= \left(\varphi(q) \int_0^\infty e^{-(\varphi(q)-\Phi(q))y} \frac{W^{(q)'(y+b)} e^{-\Phi(q)[y+b]} W^{(q)}(y+b)}{W^{(q)}(y+b) e^{-\Phi(q)b} W^{(q)}(b)} dy \right)^{-1} \\ &\xrightarrow{b \uparrow \infty} \left(\varphi(q)\Phi(q) \int_0^\infty e^{-(\varphi(q)-\Phi(q))y} dy \right)^{-1} \\ &= \frac{\varphi(q) - \Phi(q)}{\varphi(q)\Phi(q)}, \end{aligned} \tag{4.17}$$

where we recall that $\varphi(q) > \Phi(q)$. Now, applying (4.17) in (4.16), we get (4.15). □

Proposition 4.2. (i) *Under Assumption 4.1, we have $b_\Lambda \in [0, a_\Lambda]$.*

(ii) *Moreover, $b_\Lambda = 0$ if and only if one of the following two cases holds:*

- (1) $\sigma = 0$, $\Pi(0, \infty) < \infty$, and $\varphi(q) \geq (\delta + q\Lambda)(q + \Pi(0, \infty))/(c + q\Lambda)(c - \delta) =: \phi_1(\Lambda)$,
or
- (2) $\sigma > 0$ and $\varphi(q) \geq 2(\delta + q\Lambda)/\sigma^2 =: \phi_2(\Lambda)$.

Proof. (i) By the definition of b_Λ as in (4.14) and the continuity of ξ_Λ and g_Λ , in order to show that $b_\Lambda \leq a_\Lambda$, it is sufficient to show that $\xi'_\Lambda(b) \leq 0$ (or, equivalently, $\xi_\Lambda(b) - g_\Lambda(b) \leq 0$) on $[a_\Lambda, \infty)$. To show this, suppose that there exists $\bar{b} > a_\Lambda$ such that $\xi_\Lambda(\bar{b}) - g_\Lambda(\bar{b}) > 0$. Then, since g_Λ is decreasing on (a_Λ, ∞) as in Remark 4.5(ii), it follows by (4.12) that $\xi_\Lambda(b) - g_\Lambda(b)$ is increasing on (\bar{b}, ∞) . However, this contradicts (4.15). Hence, $\xi'_\Lambda(b) \leq 0$ on $[a_\Lambda, \infty)$.

(ii) Using (4.12) and the definition of b_Λ given in (4.14), we obtain $b_\Lambda = 0$ if and only if $g_\Lambda(0+) \geq \xi_\Lambda(0+)$. This is equivalent, by Lemma 4.1 and Remark 4.5(i), to

$$\varphi(q) \geq \frac{(\delta + q\Lambda)W^{(q)'}(0+)}{(1 + q\Lambda W^{(q)}(0))(1 - \delta W^{(q)}(0))}.$$

From here, using (3.5) and (3.6), we obtain the two cases announced in the proposition. \square

Remark 4.6. For the case $\sigma = 0$ and $\Pi(0, \infty) < \infty$ and the case $\sigma > 0$, the functions ϕ_1 and ϕ_2 , respectively, are both strictly increasing (since $c > \delta$ in the bounded variation case by Assumption 4.3), with $\phi_1(-\delta/q) = \phi_2(-\delta/q) = 0$. Hence, there exists $\bar{\Lambda} \in (-\delta/q, \infty]$ such that $b_\Lambda = 0$ for all $-\delta/q < \Lambda \leq \bar{\Lambda}$ and $b_\Lambda > 0$ for all $\Lambda > \bar{\Lambda}$.

(i) Suppose that $\sigma = 0$ and $\Pi(0, \infty) < \infty$. Define

$$\phi_1(\infty) := \lim_{\Lambda \rightarrow \infty} \phi_1(\Lambda) = \frac{q + \Pi(0, \infty)}{c - \delta}.$$

If $\varphi(q) \geq \phi_1(\infty)$ then, by Proposition 4.2(ii), $\bar{\Lambda} = \infty$. Otherwise, we must have $\bar{\Lambda} < \infty$ and $\phi_1(\bar{\Lambda}) = \varphi(q)$.

(ii) Suppose that $\sigma > 0$. Because $\lim_{\Lambda \rightarrow \infty} \phi_2(\Lambda) = \infty$, we must have $\bar{\Lambda} < \infty$. This also implies that $\phi_2(\bar{\Lambda}) = \varphi(q)$.

(iii) Suppose that $\sigma = 0$ and $\Pi(0, \infty) = \infty$. Then, we can set $\bar{\Lambda} = -\delta/q$.

4.3. Verification

For the case $b_\Lambda > 0$, by how b_Λ is selected as in (4.14), together with (4.5) and (4.12), we can write

$$\begin{aligned} v_\Lambda^{b_\Lambda}(x) &= g_\Lambda(b_\Lambda) \left(W^{(q)}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy \right) \\ &\quad - \Lambda \left(Z^{(q)}(x) + \delta q \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)}(y) dy \right) - \delta \bar{\mathbb{W}}^{(q)}(x - b_\Lambda) \quad \text{for } x \geq 0. \end{aligned} \tag{4.18}$$

For the case $b_\Lambda = 0$, the function $v_\Lambda^{b_\Lambda} \equiv v_\Lambda^0$ is given in (4.11).

Given the spectrally negative Lévy process X , we call a function $F: \mathbb{R} \rightarrow \mathbb{R}$ sufficiently smooth if F is continuously differentiable on $(0, \infty)$ when X has paths of bounded variation and is twice continuously differentiable on $(0, \infty)$ when X has paths of unbounded variation. We let Γ be the operator acting on a sufficiently smooth function F , defined by

$$\Gamma F(x) := \gamma F'(x) + \frac{\sigma^2}{2} F''(x) + \int_{(0, \infty)} (F(x-z) - F(x) + F'(x)z \mathbf{1}_{\{0 < z \leq 1\}}) \Pi(dz), \quad x > 0.$$

The following lemma constitutes standard technology as far as optimal control is concerned. For its proof we refer the reader to that of Lemma 1 in [14].

Lemma 4.3. Suppose that $\hat{D} \in \Theta$ is an admissible dividend strategy such that $v_\Lambda^{\hat{D}}$ is sufficiently smooth on $(0, \infty)$, $v_\Lambda^{\hat{D}}(0) \geq -\Lambda$, and, for all $x > 0$,

$$(\Gamma - q)v_\Lambda^{\hat{D}}(x) + \sup_{0 \leq r \leq \delta} (r - rv_\Lambda^{\hat{D}}(x)) \leq 0. \tag{4.19}$$

Then $v_\Lambda^{\hat{D}}(x) = V_\Lambda(x)$ for all $x \geq 0$ and hence, \hat{D} is an optimal strategy.

We first show that our candidate value function $v_{\Lambda}^{b_{\Lambda}}$ is indeed sufficiently smooth on $(0, \infty)$. The proof is given in Appendix A.2.

Lemma 4.4. Consider $b_{\Lambda} \geq 0$ given by (4.14). Then $v_{\Lambda}^{b_{\Lambda}}$ is sufficiently smooth on $(0, \infty)$.

In order to prove the HJB inequality (4.19), we use a more friendly sufficient condition. For the proof of the following result, we refer the reader to the proof of Lemma 7 in [12].

Lemma 4.5. The value function $v_{\Lambda}^{b_{\Lambda}}$ satisfies (4.19) if and only if

$$v_{\Lambda}^{b_{\Lambda}'}(x) \geq 1 \quad \text{if } 0 < x \leq b_{\Lambda}, \quad v_{\Lambda}^{b_{\Lambda}'}(x) \leq 1 \quad \text{if } x > b_{\Lambda}. \tag{4.20}$$

We now state the main theorem of this section. Its proof is provided in Appendix A.3.1.

Theorem 4.1. The optimal strategy for (2.7) consists of a refraction strategy at level b_{Λ} , given by (4.14), and the corresponding value function is given by (4.18).

5. Solution of the constrained de Finetti problem

In this section, we study the constrained de Finetti problem given in (2.2) under Assumptions 4.1 and 4.3. In order to solve this problem, we use the results in Section 4, noting that the optimal strategy for (2.5) for any $K \in [0, 1]$ is the same as the case $K = 0$, i.e. $D^{b_{\Lambda}}$, with b_{Λ} as in (4.14), is the optimal strategy for (2.5). See the discussion at the end of Section 2.3.

Throughout this section, we assume the following (see Remark 5.2 for the case it does not hold).

Assumption 5.1. We assume that $\bar{\Lambda} < \infty$, which, by Remark 4.6, is equivalent to

$$\varphi(q) < \phi_1(\infty) = \frac{q + \Pi(0, \infty)}{c - \delta} \quad \text{if } \sigma = 0 \text{ and } \Pi(0, \infty) < \infty. \tag{5.1}$$

First we need to study the relationship between Λ and its corresponding optimal threshold level b_{Λ} , which will give us enough tools to see if problem (2.2) is feasible or not.

Recall Remark 4.6 and fix $\Lambda > \bar{\Lambda}$ (then necessarily $b_{\Lambda} > 0$). Since $\xi'_{\Lambda}(b_{\Lambda}) = 0$ and by the first equality of (4.12), we observe that Λ and b_{Λ} satisfy the relation $\Lambda = \lambda(b_{\Lambda})$, where

$$\lambda(b) := (qH(b))^{-1} \tag{5.2}$$

with

$$H(b) := \frac{[h(b)]^2}{h'(b)} - \left(W^{(q)}(b) + \frac{h(b)}{\varphi(q)} \right). \tag{5.3}$$

The following results proved in Appendices A.4 and A.5 give properties of the functions defined above.

Lemma 5.1. The function $H(b)$ given in (5.3) is positive and strictly decreasing for $b \in (b_0, \infty)$, where we recall that b_0 is as in (4.14) when $\Lambda = 0$.

Note that Lemma 5.1 implies that $\lambda(b)$ is finite, positive, and strictly increasing for $b \in (b_0, \infty)$.

Proposition 5.1. Assume that (5.1) holds. Then, the function $\lambda(b)$, given in (5.2), satisfies

(i) $\lim_{b \downarrow b_0} \lambda(b) = \bar{\Lambda} \vee 0$,

- (ii) $\lim_{b \rightarrow \infty} \lambda(b) = \infty$, and
- (iii) $b_{\lambda(b)} = b$ for all $b > b_0$.

Now, by (4.2) and (4.8), we see that

$$\Psi_x(b) = Z^{(q)}(x) + \delta q \int_b^x \mathbb{W}^{(q)}(x-y)W^{(q)}(y) dy - \frac{q(W^{(q)}(b) + h(b)/\varphi(q))}{h(b)} \left(W^{(q)}(x) + \delta \int_b^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy \right), \tag{5.4}$$

for all $b \in [0, \infty)$.

Remark 5.1. Note that if $x = 0$ and X is of unbounded variation, we immediately obtain $\Psi_0(b) = 1$ for all $b \geq 0$. Hence, we obtain $V(0; K) = 0$ if $K = 1$, and the problem is infeasible otherwise.

The proof of the following result is deferred to Appendix A.6.

Lemma 5.2. Assume that $x \geq 0$ and in the case X is of unbounded variation that $x > 0$. Then the function $b \mapsto \Psi_x(b)$ defined in (5.4) is strictly decreasing on $[0, \infty)$ and

$$K_x := \lim_{b \rightarrow \infty} \Psi_x(b) = Z^{(q)}(x) - \frac{q}{\Phi(q)} W^{(q)}(x). \tag{5.5}$$

By Remark 4.4(ii), the limit of (4.1) becomes

$$\lim_{b \rightarrow \infty} \mathbb{E}_x \left[\int_0^{\tau_b} e^{-qt} dD_t^b \right] = 0, \quad x \geq 0. \tag{5.6}$$

We will define, for $K \geq 0$,

$$v_\Lambda^b(x; K) := v_\Lambda^b(x) + \Lambda K.$$

Using (5.6) and Lemma 5.2 in (4.4), we obtain

$$\lim_{b \rightarrow \infty} v_\Lambda^b(x; K) = \Lambda(K - K_x) \quad \text{for all } x > 0. \tag{5.7}$$

We are now ready to characterize the solution of (2.2). We define the *do-nothing* strategy as $D^\infty = 0$ uniformly in time and hence, $U_t^{D^\infty} = X$. By Equation (8.9) of [10], (5.7), and Lemma 5.2, we confirm the convergence results

$$\mathbb{E}_x[e^{-q\tau^{D^\infty}}] = K_x = \lim_{b \rightarrow \infty} \Psi_x(b) \quad \text{and} \quad v_\Lambda^{D^\infty}(x; K) = \Lambda(K - K_x) = \lim_{b \rightarrow \infty} v_\Lambda^b(x; K), \tag{5.8}$$

where $\tau^{D^\infty} := \inf\{t > 0 : X_t < 0\}$. From (2.4) and (5.8), we observe that if $K \geq K_x$ then D^∞ is feasible for problem (2.2) and hence,

$$V(x; K) = \sup_{D \in \Theta} \inf_{\Lambda \geq 0} v_\Lambda^D(x; K) \geq \inf_{\Lambda \geq 0} v_\Lambda^{D^\infty}(x; K) = 0 \quad \text{if } K \in [K_x, 1], \tag{5.9}$$

with v_Λ^D as in (2.3).

Theorem 5.1. *Let $x \geq 0$ be fixed. Assume that (5.1) and one of the following cases hold:*

- (i) $x > 0$ and X is of unbounded variation;
- (ii) $x \geq 0$ and X is of bounded variation.

Then

$$V(x; K) = \begin{cases} v_0^{b_0}(x; K) & \text{if } K \in [\Psi_x(b_0), 1], \\ \inf_{\Lambda \geq 0} V_\Lambda(x; K) & \text{if } K \in (K_x, \Psi_x(b_0)), \\ 0 & \text{if } K = K_x, \\ -\infty & \text{if } K \in [0, K_x). \end{cases}$$

Proof. We will only prove (i) since the other case is similar. Recall inequality (2.4) and that $V_\Lambda(x; K)$ is defined as in (2.5).

(a) If $K \in [\Psi_x(b_0), 1]$ then the threshold strategy at level b_0 is feasible for problem (2.2) (see Section 2), and therefore,

$$v_0^{b_0}(x; K) \leq V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) \leq V_0(x; K) = v_0^{b_0}(x; K).$$

Here the second inequality holds by (2.4) and the last equality holds because the case $\Lambda = 0$ is solved by the threshold strategy with b_0 in problem (2.5) (which is equivalent to (2.7)).

(b) If $K \in (K_x, \Psi_x(b_0))$, since $b \mapsto \Psi_x(b)$ is continuous and strictly decreasing, by Lemma 5.2, there exists a unique $b^* > b_0$ such that $K - \Psi_x(b^*) = 0$. Therefore,

$$V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) \leq V_{\lambda(b^*)}(x; K) = \mathbb{E}_x \left[\int_0^{\tau_{b^*}} e^{-qt} dD_t^{b^*} \right] \leq V(x; K).$$

Here the first inequality holds by (2.4). The equality follows from Proposition 5.1(iii) since D^{b^*} is the optimal strategy for (2.5) when $\Lambda = \lambda(b^*)$. The last inequality follows since the threshold strategy at level b^* is feasible for problem (2.2).

(c) If $K = K_x$, by Lemma 5.2, we have $\lambda(b)(K - \Psi_x(b)) \leq 0$ for all $b > b_0$. Hence,

$$0 \leq V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) \leq \inf_{b > b_0} \left(\mathbb{E}_x \left[\int_0^{\tau_b} e^{-qt} dD_t^b \right] + \lambda(b)(K - \Psi_x(b)) \right) \leq 0.$$

Here the first inequality follows from (5.9), the second by (2.4), and the last by (5.6).

(d) Finally, if $K \in [0, K_x)$ then Lemma 5.2 gives $\lim_{b \rightarrow \infty} [K - \Psi_x(b)] = K - K_x < 0$. Hence, we obtain

$$V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) \leq \inf_{b > b_0} \left(\mathbb{E}_x \left[\int_0^{\tau_b} e^{-qt} dD_t^b \right] + \lambda(b)(K - \Psi_x(b)) \right) = -\infty,$$

where the second inequality holds by equation (2.4) and the last equality follows from Proposition 5.1. □

Remark 5.2. Consider the case when Assumption 5.1 is violated. By Remark 4.6, for all $\Lambda \geq 0$, we must have $b_\Lambda = 0$ and hence, $V_\Lambda(x; K) = \mathbb{E}_x[\int_0^{\tau_0} e^{-qt} dD_t^0] + \Lambda(K - \Psi_x(0))$. If $K \in [\Psi_x(b_0), 1]$ then the threshold strategy at level 0 is feasible for problem (2.2) (see Section 2), and therefore,

$$v_0^0(x; K) \leq V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) \leq V_0(x; K) = v_0^0(x; K).$$

On the other hand, if $K \in [0, \Psi_x(0))$, we obtain

$$V(x; K) \leq \inf_{\Lambda \geq 0} V_\Lambda(x; K) = \inf_{\Lambda \geq 0} \left(\mathbb{E}_x \left[\int_0^{\tau_0} e^{-qt} dD_t^0 \right] + \Lambda(K - \Psi_x(0)) \right) = -\infty.$$

In sum, we have

$$V(x; K) = \begin{cases} v_0^0(x, K) & \text{if } K \in [\Psi_x(0), 1], \\ -\infty & \text{if } K \in [0, \Psi_x(0)). \end{cases}$$

6. Spectrally positive case

In this section, we solve analogous problems driven by a spectrally positive Lévy process \bar{Y} . We assume that its dual process $Y = -\bar{Y}$ has its Laplace exponent ψ_Y as in (3.1) so that its right inverse and scale function are given by $\varphi(q)$ and $\mathbb{W}^{(q)}$, respectively. We also define the drift-changed process $\bar{X} = \{\bar{X}_t = \bar{Y}_t - \delta t; t \geq 0\}$ whose dual $X = -\bar{X}$ has its Laplace exponent ψ as in (2.1), right inverse $\Phi(q)$, and scale function $W^{(q)}$ described in Section 3. We denote by $\bar{\mathbb{E}}_x$ the expectation with respect to the law of the process \bar{Y} when it starts at x .

In addition, for $x \geq 0$ and $0 < K \leq 1$, we define $\bar{v}(x; K)$, $\bar{v}_\Lambda^D(x; K)$, $\bar{v}_\Lambda(x; K)$ and $\bar{v}_\Lambda(x)$ analogously to (2.2), (2.3), (2.5), and (2.6), respectively.

We first solve the optimal dividend problem with terminal payoff/penalty (2.7) with X replaced with \bar{Y} . Similarly to the spectrally negative Lévy case, we define for $b \geq 0$ the threshold strategy \bar{D}^b and the resulting controlled surplus process, which is a refracted spectrally positive Lévy process defined as the unique strong solution to the stochastic differential equation

$$\bar{U}_t^b := \bar{Y}_t - \bar{D}_t^b := \bar{Y}_t - \delta \int_0^t \mathbf{1}_{\{\bar{U}_s^b > b\}} ds, \quad t \geq 0.$$

Let its ruin time be denoted by

$$\bar{\tau}_b := \inf\{t > 0: \bar{U}_t^b < 0\}.$$

6.1. Scale functions under a change of measure

For each $\beta \geq 0$, we define the change of measure

$$\frac{d\tilde{\mathbb{P}}_x^\beta}{d\mathbb{P}_x} \Big|_{\mathcal{F}_t} = e^{\beta(Y_t - x) - \psi_Y(\beta)t}, \quad x \in \mathbb{R}, t \geq 0,$$

where $\tilde{\mathbb{P}}_x$ is the law of the process Y when it starts at x . It is known that Y is still a spectrally negative Lévy process on $(\Omega, \mathcal{F}, \tilde{\mathbb{P}}^\beta)$ and the scale function of Y on this probability space can be written

$$\begin{aligned} \mathbb{W}_\beta^{(u - \psi_Y(\beta))}(x) &= e^{-\beta x} \mathbb{W}^{(u)}(x), \\ \mathbb{Z}_\beta^{(u - \psi_Y(\beta))}(x) &= 1 + (u - \psi_Y(\beta)) \int_0^x e^{-\beta z} \mathbb{W}^{(u)}(z) dz, \end{aligned}$$

with $u - \psi_Y(\beta) \geq 0$; see [2, Remark 4]. In particular, (3.1) and (3.3) give $q - \psi_Y(\Phi(q)) = \delta\Phi(q)$ and hence,

$$\mathbb{W}_{\Phi(q)}^{(\delta\Phi(q))}(x) = e^{-\Phi(q)x} \mathbb{W}^{(q)}(x) \quad \text{and} \quad \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(x) = 1 + \delta\Phi(q) \int_0^x e^{-\Phi(q)z} \mathbb{W}^{(q)}(z) dz. \quad (6.1)$$

6.2. Optimal dividend problem with terminal value

As in the case of the spectrally negative Lévy process, we are first interested in solving problem (2.4) for the spectrally positive case. For this purpose, first we need to study the optimal dividend problem with a terminal value for the process \bar{Y} . Using Theorems 5(i) and 6(iii) of [11] we have the following result, whose proof is deferred to Appendix A.7.

Proposition 6.1. *For $x, b, q \geq 0$, we have*

$$\bar{\Psi}_x(b) := \bar{\mathbb{E}}_x[e^{-q\bar{\tau}_b}; \bar{\tau}_b < \infty] = e^{-\Phi(q)x} \frac{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b-x)}{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)}, \tag{6.2}$$

where $\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(x)$ is given in (6.1), $\bar{\tau}_b := \inf\{t > 0: \bar{U}_t^b = 0\}$, and

$$\bar{\mathbb{E}}_x \left[\int_0^{\bar{\tau}_b} e^{-qt} d\bar{D}_t^b \right] = \frac{\delta}{q} (\mathbb{Z}^{(q)}(b-x) - \mathbb{Z}^{(q)}(b) \bar{\Psi}_x(b)). \tag{6.3}$$

Using (6.2) and (6.3), we have the following result.

Proposition 6.2. *For $b \geq 0$, we have*

$$\bar{v}_\Lambda^b(x) := \bar{\mathbb{E}}_x \left[\int_0^{\bar{\tau}_b} e^{-qt} d\bar{D}_t^b \right] - \Lambda \bar{\Psi}_x(b) = \begin{cases} \frac{\delta}{q} \mathbb{Z}^{(q)}(b-x) - \bar{k}_x(b, \Lambda) & \text{if } 0 \leq x \leq b, \\ \frac{\delta}{q} - \bar{k}_b(b, \Lambda) e^{-\Phi(q)(x-b)} & \text{if } x > b, \end{cases} \tag{6.4}$$

where, for $x \geq 0$,

$$\bar{k}_x(b, \Lambda) := \bar{\Psi}_x(b) \left(\frac{\delta}{q} \mathbb{Z}^{(q)}(b) + \Lambda \right).$$

In order to select the optimal threshold, we apply smooth fit. Note that, by (6.4), \bar{v}_Λ^b is continuous on $[0, \infty)$ for any choice of b . Here, we will study the smoothness of \bar{v}_Λ^b at $x = b$ to propose a candidate threshold level \bar{b}_Λ such that \bar{v}_Λ^b is $C^1(0, \infty)$ and $C^2(0, \infty)$ when \bar{Y} is of bounded and unbounded variation, respectively. By differentiating (6.4), we see that

$$\bar{v}'_\Lambda(x) = \begin{cases} -\delta \mathbb{W}^{(q)}(b-x) + \Phi(q) \bar{k}_x(b, \Lambda) + \delta \bar{k}'_b(b, \Lambda) \mathbb{W}^{(q)}(b-x) & \text{if } 0 < x < b, \\ \Phi(q) \bar{k}'_b(b, \Lambda) e^{-\Phi(q)(x-b)} & \text{if } x > b, \end{cases} \tag{6.5}$$

and, in particular, for the unbounded variation case

$$\bar{v}''_\Lambda(x) = \begin{cases} \delta \mathbb{W}^{(q)'}(b-x) - [\Phi(q)]^2 \bar{k}_x(b, \Lambda) \\ \quad - \delta \Phi(q) \bar{k}'_b(b, \Lambda) [\Phi(q) \mathbb{W}^{(q)}(b-x) + \mathbb{W}^{(q)'}(b-x)] & \text{if } 0 < x < b, \\ -[\Phi(q)]^2 \bar{k}'_b(b, \Lambda) e^{-\Phi(q)(x-b)} & \text{if } x > b, \end{cases} \tag{6.6}$$

where we recall that, if Y is of unbounded variation, $\mathbb{W}^{(q)}$ is C^1 on $(0, \infty)$. From (6.5) and (6.6), together with (3.5) and (3.6), we have the following result.

Lemma 6.1. Suppose that $b > 0$ is such that

$$\bar{k}_b(b, \Lambda) = \frac{1}{\Phi(q)}, \quad \text{or, equivalently,} \quad \Lambda e^{-\Phi(q)b} = s(b), \tag{6.7}$$

where

$$s(b) := \frac{1}{\Phi(q)} \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b) - \frac{\delta e^{-\Phi(q)b}}{q} \mathbb{Z}^{(q)}(b), \quad b > 0. \tag{6.8}$$

Then, the function \bar{v}_Λ^b is $C^1(0, \infty)$ and $C^2(0, \infty)$ for the case of bounded and unbounded variation, respectively.

Lemma 6.2. If $\Lambda > 1/\Phi(q) - \delta/q$, then there exists a unique $b > 0$ that satisfies (6.7).

Proof. In order to prove the lemma, we will show that $s(b)$ as in (6.8) is strictly increasing and satisfies

$$\lim_{b \rightarrow 0} s(b) = \frac{1}{\Phi(q)} - \frac{\delta}{q} \quad \text{and} \quad \lim_{b \rightarrow \infty} s(b) = \infty. \tag{6.9}$$

(i) Since $s'(b) = \delta\Phi(q)e^{-\Phi(q)b}\mathbb{Z}^{(q)}(b)/q > 0$ for all $b > 0$, then $s(\cdot)$ is strictly increasing on $(0, \infty)$.

(ii) Letting $b \rightarrow 0$ in (6.8), it is easy to see that the first limit of (6.9) holds.

(iii) Note that

$$s(b) = \frac{\delta\mathbb{Z}^{(q)}(b)}{q\Phi(q)e^{\Phi(q)b}} \left(\frac{qe^{\Phi(q)b}\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)}{\delta\mathbb{Z}^{(q)}(b)} - \Phi(q) \right). \tag{6.10}$$

Using l'Hôpital's rule, (3.7), and the fact that $\varphi(q) > \Phi(q)$, the following limits can be verified:

$$\begin{aligned} \lim_{b \rightarrow \infty} \frac{\mathbb{Z}^{(q)}(b)}{e^{\Phi(q)b}} &= \lim_{b \rightarrow \infty} q[\Phi(q)]^{-1} e^{(\varphi(q)-\Phi(q))b} e^{-\varphi(q)b} \mathbb{W}^{(q)}(b) = \infty, \\ \lim_{b \rightarrow \infty} \frac{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)}{e^{(\varphi(q)-\Phi(q))b}} &= \lim_{b \rightarrow \infty} \delta\Phi(q) \frac{e^{-\varphi(q)b} \mathbb{W}^{(q)}(b)}{\varphi(q) - \Phi(q)} = \delta\Phi(q) \frac{\psi'_Y(\varphi(q))^{-1}}{\varphi(q) - \Phi(q)}, \\ \lim_{b \rightarrow \infty} \frac{qe^{\Phi(q)b}\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)}{\delta\mathbb{Z}^{(q)}(b)} &= \lim_{b \rightarrow \infty} \frac{\Phi(q)(\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)/\delta e^{(\varphi(q)-\Phi(q))b} + e^{-\varphi(q)b} \mathbb{W}^{(q)}(b))}{e^{-\varphi(q)b} \mathbb{W}^{(q)}(b)} \\ &= \Phi(q) \left(1 + \frac{\Phi(q)}{\varphi(q) - \Phi(q)} \right). \end{aligned}$$

Hence, it follows that $\lim_{b \rightarrow \infty} s(b) = \infty$. □

Now, we let \bar{b}_Λ be as in Lemma 6.2 for the case $\Lambda > 1/\Phi(q) - \delta/q$ and set it to 0 otherwise.

(i) When $\Lambda > 1/\Phi(q) - \delta/q$, applying (6.7) in (6.4), with $b = \bar{b}_\Lambda$, we see that $\bar{v}_\Lambda^{\bar{b}_\Lambda}$ is given by

$$\bar{v}_\Lambda^{\bar{b}_\Lambda}(x) = \begin{cases} \frac{\delta}{q} \mathbb{Z}^{(q)}(\bar{b}_\Lambda - x) - \frac{e^{-\Phi(q)(x-\bar{b}_\Lambda)}}{\Phi(q)} \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(\bar{b}_\Lambda - x) & \text{if } x \leq \bar{b}_\Lambda, \\ \frac{\delta}{q} - \frac{e^{-\Phi(q)(x-\bar{b}_\Lambda)}}{\Phi(q)} & \text{if } x > \bar{b}_\Lambda. \end{cases} \tag{6.11}$$

(ii) When $\Lambda \leq 1/\Phi(q) - \delta/q$, using (6.4) and because $\bar{k}_0(0, \Lambda) = \delta/q + \Lambda$, we have

$$\bar{v}_\Lambda^0(x) = \frac{\delta}{q} - e^{-\Phi(q)x} \left(\frac{\delta}{q} + \Lambda \right), \quad x \geq 0. \tag{6.12}$$

Theorem 6.1. *The optimal strategy for (2.7) consists of a threshold strategy at level \bar{b}_Λ .*

Proof. In view of (6.11) and (6.12), we confirm that $\bar{v}_\Lambda^{\bar{b}_\Lambda}$ is sufficiently smooth. Hence, as in the spectrally negative case, in order to verify that $\bar{D}^{\bar{b}_\Lambda}$ is the optimal strategy over all admissible strategies, it is sufficient to show that the cost function $\bar{v}_\Lambda^{\bar{b}_\Lambda}$, given by (6.11) and (6.12), satisfies (4.20) and that $\bar{v}_\Lambda^{\bar{b}_\Lambda}(0) \geq -\Lambda$.

(i) Suppose that $\bar{b}_\Lambda > 0$, and so the threshold level \bar{b}_Λ satisfies (6.7). From (6.11) we have

$$\bar{v}_\Lambda^{\bar{b}_\Lambda}(x) = \begin{cases} e^{-\Phi(q)(x-\bar{b}_\Lambda)} \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(\bar{b}_\Lambda - x) & \text{if } x \leq \bar{b}_\Lambda, \\ e^{-\Phi(q)(x-\bar{b}_\Lambda)} & \text{if } x > \bar{b}_\Lambda. \end{cases}$$

Clearly, $\bar{v}_\Lambda^{\bar{b}_\Lambda}(x) < 1$ if $x > \bar{b}_\Lambda$. On the other hand, $\bar{v}_\Lambda^{\bar{b}_\Lambda}(x)$ is strictly decreasing on $[0, \bar{b}_\Lambda]$ since $x \mapsto e^{-\Phi(q)(x-\bar{b}_\Lambda)}$ is strictly decreasing and $x \mapsto \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(\bar{b}_\Lambda - x)$ is nonincreasing in the interval. This together with $\bar{v}_\Lambda^{\bar{b}_\Lambda}(\bar{b}_\Lambda) = 1$ shows that $\bar{v}_\Lambda^{\bar{b}_\Lambda}(x) \geq 1$ if $x \leq \bar{b}_\Lambda$.

Finally, we note that, using (6.7), (6.8), and (6.11),

$$\bar{v}_\Lambda^{\bar{b}_\Lambda}(0) = \frac{\delta}{q} \mathbb{Z}^{(q)}(\bar{b}_\Lambda) - \frac{e^{\Phi(q)\bar{b}_\Lambda}}{\Phi(q)} \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(\bar{b}_\Lambda) = -e^{\Phi(q)\bar{b}_\Lambda} s(\bar{b}_\Lambda) = -\Lambda.$$

(ii) Suppose that $\bar{b}_\Lambda = 0$. Since $\Lambda \leq 1/\Phi(q) - \delta/q$, it follows that, for $x \geq 0$,

$$\bar{v}_\Lambda^0(x) = \Phi(q)e^{-\Phi(q)x} \left(\frac{\delta}{q} + \Lambda \right) \leq e^{-\Phi(q)x} \leq 1.$$

Finally, by (6.12) we obtain $\bar{v}_\Lambda^0(0) = -\Lambda$. □

6.3. Constrained de Finetti’s problem for spectrally positive Lévy processes

Now we consider problem (2.2) driven by the spectrally positive Lévy process \bar{Y} . Note that $\bar{D}^{\bar{b}_\Lambda}$ is the optimal strategy for (2.5) for any $K \in [0, 1]$.

Let us define $\tilde{\Lambda} := 1/\Phi(q) - \delta/q$. Following Lemma 6.2, if $\tilde{\Lambda} \geq 0$, we have $\bar{b}_\Lambda = 0$ for $\Lambda \in [0, \tilde{\Lambda}]$; on the other hand, if $\tilde{\Lambda} < 0$ then $\bar{b}_\Lambda > 0$ for all $\Lambda > 0$.

Similarly to Section 5, we need to establish the relationship between Λ and its corresponding threshold level \bar{b}_Λ given by Lemma 6.2. From (6.7) we obtain $\Lambda = \tilde{\lambda}(\bar{b}_\Lambda)$ for $\Lambda > \tilde{\Lambda}$, where

$$\tilde{\lambda}(b) = e^{\Phi(q)b} s(b),$$

with s defined in (6.10). Since s is strictly increasing (see the proof of Lemma 6.2) and satisfies (6.9), it immediately follows that $\tilde{\lambda}$ is also strictly increasing, $\lim_{b \rightarrow \infty} \tilde{\lambda}(b) = \infty$, and

$$\lim_{b \rightarrow \bar{b}_0} \tilde{\lambda}(b) = \begin{cases} \frac{1}{\Phi(q)} - \frac{\delta}{q} & \text{if } \bar{b}_0 = 0, \\ 0 & \text{if } \bar{b}_0 > 0. \end{cases}$$

Here, the convergence for the case $\bar{b}_0 > 0$ holds by the fact that

$$\lim_{b \rightarrow \bar{b}_0} \tilde{\lambda}(b) = \lim_{b \rightarrow \bar{b}_0} e^{\Phi(q)b} s(b) = e^{\Phi(q)\bar{b}_0} s(\bar{b}_0) = 0,$$

where the last equality follows because (6.7) and Lemma 6.2 imply that $s(\bar{b}_0) = 0$. Note that $\bar{b}_{\tilde{\lambda}(b)} = b$ for all $b > \bar{b}_0$.

Next, we need to show that the function $b \mapsto \bar{\Psi}_x(b)$, given in (6.2), is strictly decreasing with $x > 0$ fixed. In the case that $x = 0$, we see that $\bar{\Psi}_0(b) = 1$ by (6.2). The proof of the following lemma is given in Appendix A.8.

Lemma 6.3. *Let $x > 0$ be fixed. Then the function $b \mapsto \bar{\Psi}_x(b)$ defined in (6.2) is strictly decreasing and satisfies*

$$\lim_{b \rightarrow 0} \bar{\Psi}_x(b) = e^{-\Phi(q)x} \quad \text{and} \quad \lim_{b \rightarrow \infty} \bar{\Psi}_x(b) = e^{-\varphi(q)x}.$$

Finally, using similar arguments as in Theorem 5.1 (noting that we have results analogous to Lemmas 5.1(iv) and 5.2), we obtain the following theorem.

Theorem 6.2. *Let $x > 0$ be fixed. Then*

$$\bar{v}(x; K) = \begin{cases} \bar{v}_0^{\bar{b}_0}(x; K) & \text{if } K \in [\bar{\Psi}_x(\bar{b}_0), 1], \\ \inf_{\Lambda \geq 0} \bar{v}_\Lambda(x; K) & \text{if } K \in (e^{-\varphi(q)x}, \bar{\Psi}_x(\bar{b}_0)), \\ 0 & \text{if } K = e^{-\varphi(q)x}, \\ -\infty & \text{if } K \in [0, e^{-\varphi(q)x}). \end{cases}$$

7. Numerical examples

In this section, we confirm the obtained results through a sequence of numerical examples for both spectrally negative and positive cases. Throughout this section, we set $q = 0.05$.

7.1. Spectrally negative case

We first consider the spectrally negative case as studied in Sections 4 and 5. Here we assume that X is of the form

$$X_t - X_0 = ct + 0.2B_t - \sum_{n=1}^{N_t} Z_n, \quad 0 \leq t < \infty, \tag{7.1}$$

where $B = \{B_t : t \geq 0\}$ is a standard Brownian motion, $N = \{N_t : t \geq 0\}$ is a Poisson process with arrival rate κ , and $Z = \{Z_n; n = 1, 2, \dots\}$ is an i.i.d. sequence of exponential variables with parameter 1 (so that Assumption 4.1 is satisfied). Here, the processes B , N , and Z are assumed mutually independent. We refer the reader to [5] and [9] for the forms of the corresponding scale functions.

We consider the following two parameter sets.

Case 1: $\kappa = 1, \sigma = 0.2, c = 1.5, \delta = 1$.

Case 2: $\kappa = 0.01, \sigma = 0, c = 5, \delta = 0.1$.

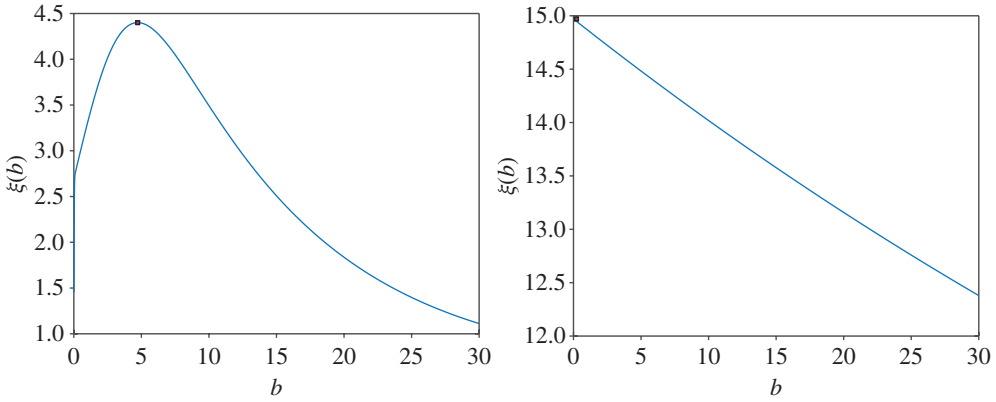


FIGURE 1: Plots $b \mapsto \xi_\Lambda(b)$ for case 1 (left) and case 2 (right). The points at b_Λ are indicated by squares.

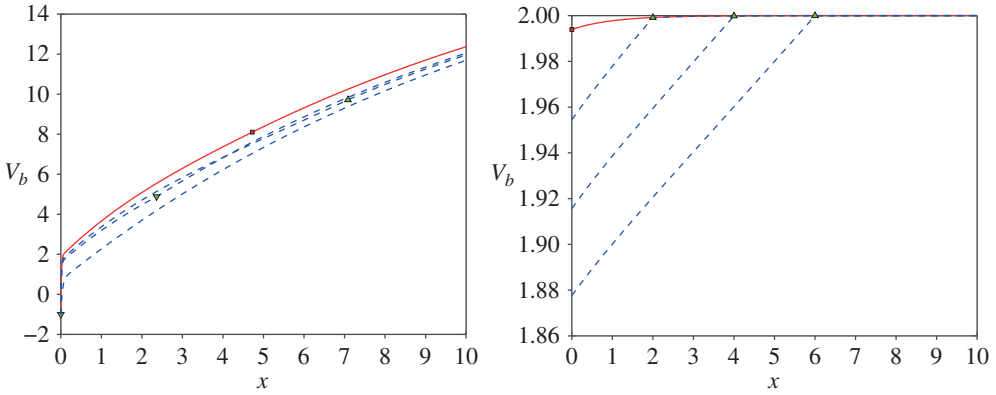


FIGURE 2: Plots of $x \mapsto V_\Lambda(x)$ (solid lines) for case 1 (left) and case 2 (right). Suboptimal value functions v_Λ^b (dotted lines) are also plotted for the choice of $b = 0, \bar{b}_\Lambda/2, 3\bar{b}_\Lambda/2$ for case 1 and $b = 2, 4, 6$ for case 2. The points at b_Λ are indicated by squares and those at b in the suboptimal cases are indicated by up- (respectively down-) pointing triangles when $b > b_\Lambda$ (respectively $b < b_\Lambda$).

Here, case 2 corresponds to the case $\bar{\Lambda} = \infty$ where we have $b_\Lambda = 0$ for any choice of Λ as in Remark 4.6.

We first show the optimal solutions for the problem considered in Section 4 focusing on the case $\Lambda = 1$. In Figure 1 we plot the function $b \mapsto \xi_\Lambda(b)$ as in (4.6) and (4.9). Here, in case 1, it attains a global maximum and the maximizer becomes b_Λ by (4.14). In contrast, in case 2, it is monotonically decreasing and, by (4.14), we have $b_\Lambda = 0$. In Figure 2, we plot the optimal value function $x \mapsto V_\Lambda(x) = v_\Lambda^{b_\Lambda}(x)$ along with the suboptimal value functions v_Λ^b for the choice of $b = 0, \bar{b}_\Lambda/2, 3\bar{b}_\Lambda/2$ for case 1 and $b = 2, 4, 6$ for case 2. In both cases, we confirm that V_Λ dominates the suboptimal ones uniformly in x .

We now move onto the constrained problem (2.2) studied in Section 5, focusing on case 1 with $K = 0.1$. Recall that the optimal solutions are given in Theorem 5.1. In the left-hand panel of Figure 3, we plot the function $x \mapsto V_\Lambda(x; K) = V_\Lambda(x) + \Lambda K$ for various values of Λ ranging from 0 to 20 000. For $x \in (\underline{x}, \bar{x})$, where \underline{x} and \bar{x} are such that $K_{\underline{x}} = K$ and $\Psi_{\bar{x}}(b_0) = K$, respectively, its minimum over the considered Λ gives (an approximation of) $V(x; K)$, indicated by the solid red line in the plot. On the other hand, $V(x; K)$ equals

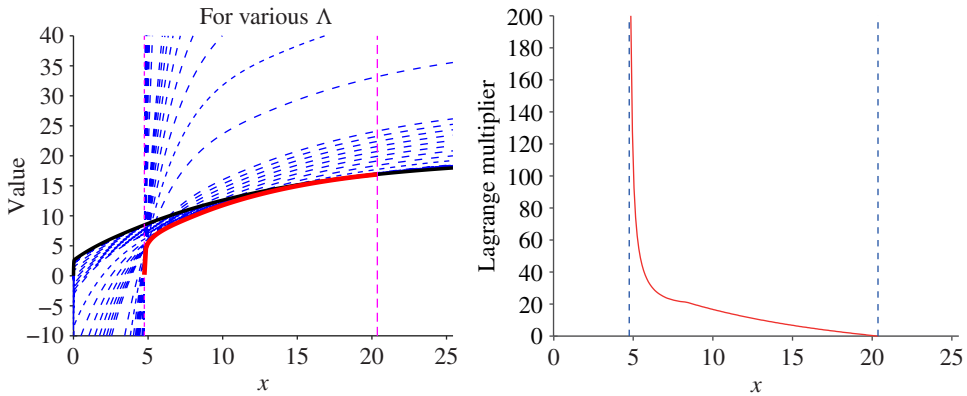


FIGURE 3: (Left) Plots of $x \mapsto V_\Lambda(x; K)$ for $\Lambda = 0.1, \dots, 1, 2, \dots, 10, 20, \dots, 100, 200, \dots, 1000, 2000, \dots, 10\,000, 20\,000$ (dotted lines) and for $\Lambda = 0$ (solid, boldface line) for the case $K = 0.1$. The two vertical dotted lines indicate the values of \underline{x} and \bar{x} such that $K_{\underline{x}} = K$ and $\Psi_{\bar{x}}(b_0) = K$. On $[\underline{x}, \bar{x}]$, the minimum of $V_\Lambda(x; K)$ over Λ is shown by a solid boldface red line. (Right) Plots of the Lagrange multiplier Λ^* on $[\underline{x}, \bar{x}]$ with the same two vertical lines as in the left plot.

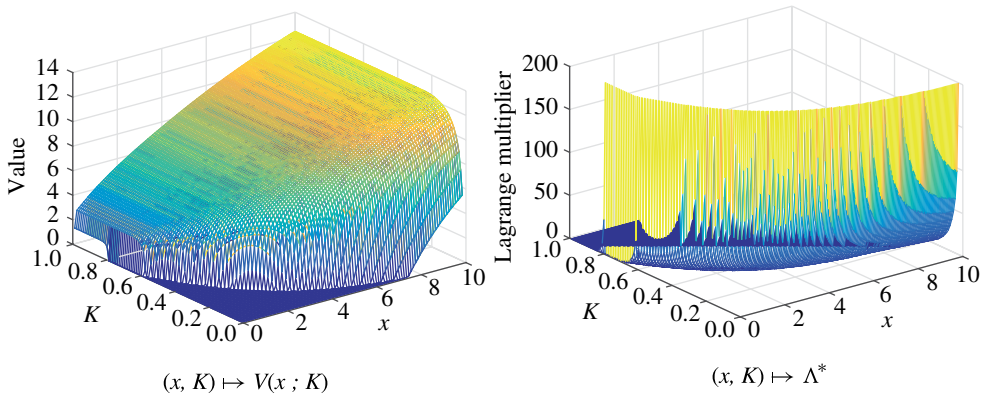


FIGURE 4: Plots of $V(x; K)$ (left) and the Lagrange multiplier Λ^* (right) as functions of x and K .

$V_0(x; K) = v_0^{b_0}(x; K)$ for $x \in [\bar{x}, \infty)$ and it is infeasible for $x \in [0, \underline{x}]$. In the right-hand panel of Figure 3, we plot, for $x \in (\underline{x}, \bar{x})$, the Lagrange multiplier $\Lambda^* := \arg \min_{\Lambda \geq 0} V_\Lambda(x; K)$. We observe that Λ^* goes to ∞ as $x \downarrow \underline{x}$ and to 0 as $x \uparrow \bar{x}$.

In Figure 4, we show the values of $V(x; K)$ and the Lagrange multiplier Λ^* as functions of (x, K) . Here, those (x, K) at which the problem is infeasible are indicated by dark shades on the $z = 0$ plane. It is confirmed that $V(x; K)$ increases as x and K increase, while Λ^* increases as (x, K) decrease.

7.2. Spectrally positive case

Similarly, we confirm the results in Section 6 focusing on the case \bar{Y} is of the form

$$\bar{Y}_t - \bar{Y}_0 = -t + 0.2B_t + \sum_{n=1}^{N_t} Z_n \quad \text{for } t \geq 0.$$

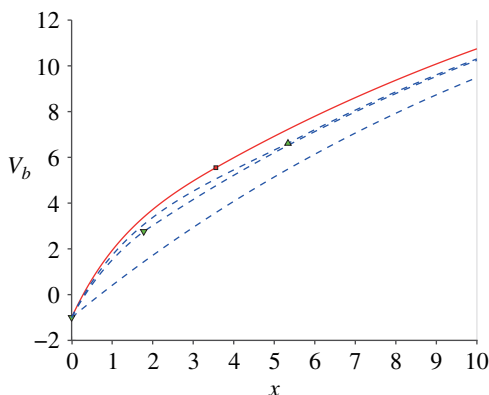


FIGURE 5: Plots of $x \mapsto \bar{v}_\Lambda(x)$ along with suboptimal value functions \bar{v}_Λ^b (dotted lines) for the choice of $b = 0, \bar{b}_\Lambda/2, 3\bar{b}_\Lambda/2$. The point at \bar{b}_Λ is indicated by a square and the points at b in the suboptimal cases are indicated by up- (respectively down-) pointing triangles when $b > \bar{b}_\Lambda$ (respectively $b < \bar{b}_\Lambda$).

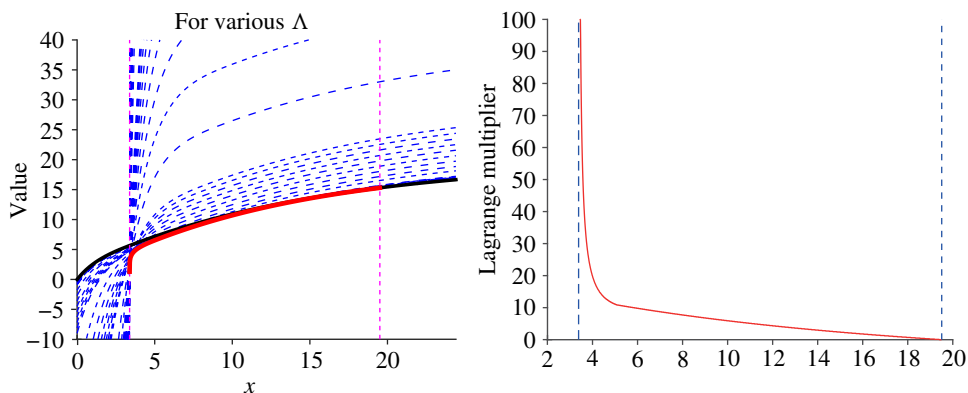


FIGURE 6: (Left) Plots of $x \mapsto \bar{v}_\Lambda(x; K)$ for $\Lambda = 0.1, \dots, 1, 2, \dots, 10, 20, \dots, 100, 200, \dots, 1000, 2000, \dots, 10\,000, 20\,000$ (dotted lines) and for $\Lambda = 0$ (solid, boldface line) for the case $K = 0.1$. The two vertical dotted lines indicate the values of \underline{x} and \bar{x} such that $\exp(-\varphi(q)\underline{x}) = K$ and $\bar{\Psi}_{\bar{x}}(\bar{b}_0) = K$. On $[\underline{x}, \bar{x}]$, the minimum of $\bar{v}_\Lambda(x; K)$ over Λ is shown by the solid boldface red line. (Right) Plots of the Lagrange multiplier Λ^* on $(\underline{x}, \bar{x}]$ with the same two vertical lines as in the left plot.

Here B and N (with $\kappa = 1.5$) are the same as in the case of (7.1), and Z is a phase-type random variable that approximates the Weibull distribution with shape parameter 2 and scale parameter 1 (see [1] for the parameters of the phase-type distribution). Throughout, we set $\delta = 1$.

For the (Lagrangian) problem considered in Section 6.2, the optimal threshold \bar{b}_Λ is such that (6.7) holds and the value function $\bar{v}_\Lambda(x) = \bar{v}_\Lambda^{b_\Lambda}(x)$ is given in (6.11). In Figure 5 we plot the optimal value function $x \mapsto \bar{v}_\Lambda(x)$ along with suboptimal value functions \bar{v}_Λ^b for the choice of $b = 0, \bar{b}_\Lambda/2, 3\bar{b}_\Lambda/2$ when $\Lambda = 1$. For the constrained case considered in Section 6.3, in Figures 6 and 7, we plot analogous results to those shown in Figures 3 and 4, where we assume that $K = 0.1$ for Figure 6. It is confirmed that similar behaviors of the value function and the Lagrange multiplier can be observed as in the spectrally negative case.

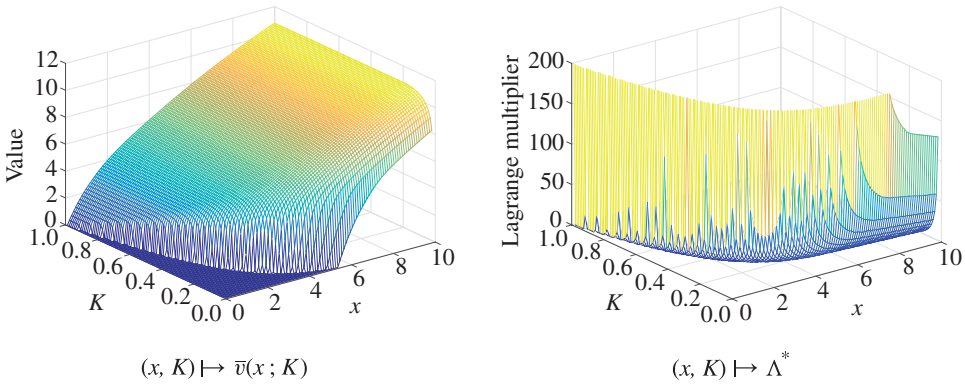


FIGURE 7: Plots of $\bar{v}(x; K)$ (left) and the Lagrange multiplier Λ^* (right) as functions of x and K .

8. Conclusions

In this paper, we studied versions of de Finetti’s problem with a constraint on the time of ruin over the set of absolutely continuous strategies. We solved it for a spectrally negative Lévy process with a completely monotone Lévy density and for a general spectrally positive Lévy process. Thanks to our analysis that the optimal solution to the Lagrangian subproblem can be characterized by a single threshold, the strong duality can be verified and hence the constrained problem can be solved efficiently.

A natural and interesting problem will be to consider the case of a general spectrally negative Lévy process, without the completely monotone Lévy density assumption. In this case, threshold strategies are in general not optimal, but as is observed in Azcue and Muler [3] in a related de Finetti’s problem, a band strategy is expected to be optimal. Here a big challenge is to show the strong duality to solve the constrained problem via Lagrangian subproblems. Our methods in this paper take advantage of the explicit expressions via the scale function as well as the relation between the (single) optimal threshold and Λ ; these fail to hold when bands are required to characterize the optimal strategy for the Lagrangian subproblem. We expect that novel approaches are required to solve the problem, but the results of this paper can potentially be generalized. These are important and challenging problems and we leave them for future work.

Appendix A. Proofs

A.1. Proof Lemma 4.1

By (3.1)–(3.3), we have

$$\int_0^\infty e^{-\varphi(q)y} W^{(q)}(y) dy = (\delta\varphi(q))^{-1}. \tag{A.1}$$

Using this and integration by parts,

$$\int_0^\infty W^{(q)'}(y) e^{-\varphi(q)y} dy = -W^{(q)}(0) + \varphi(q) \int_0^\infty W^{(q)}(y) e^{-\varphi(q)y} dy = -W^{(q)}(0) + \delta^{-1}. \tag{A.2}$$

From here, (4.10) is immediate.

Now, by differentiating (3.4) and changing variables,

$$\mathbb{W}^{(q)}(x) - W^{(q)}(x) = \delta \left(\int_0^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy + \mathbb{W}^{(q)}(x)W^{(q)}(0) \right).$$

This implies that

$$\delta \int_0^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy = \mathbb{W}^{(q)}(x) - W^{(q)}(x) - \delta \mathbb{W}^{(q)}(x)W^{(q)}(0). \quad (\text{A.3})$$

Substituting (3.4) and (A.3) into (4.5) and after simplification, we obtain (4.11).

A.2. Proof Lemma 4.4

(i) Let us consider the case $b_\Lambda > 0$. For $x \neq b_\Lambda$, by differentiating (4.18),

$$\begin{aligned} v_\Lambda^{b_\Lambda'}(x) &= g_\Lambda(b_\Lambda) \left(W^{(q)'}(x) + \delta \left[\mathbb{W}^{(q)}(0)W^{(q)'}(x) \right. \right. \\ &\quad \left. \left. + \int_{b_\Lambda}^x \mathbb{W}^{(q)'}(x-y)W^{(q)'}(y) dy \right] \mathbf{1}_{\{x > b_\Lambda\}} \right) \\ &\quad - q\Lambda \left(W^{(q)}(x) + \delta \left[\mathbb{W}^{(q)}(0)W^{(q)}(x) + \int_{b_\Lambda}^x \mathbb{W}^{(q)'}(x-y)W^{(q)}(y) dy \right] \mathbf{1}_{\{x > b_\Lambda\}} \right) \\ &\quad - \delta \mathbb{W}^{(q)}(x - b_\Lambda) \\ &= g_\Lambda(b_\Lambda) \left(W^{(q)'}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)''}(y) dy \right) \\ &\quad - q\Lambda \left(W^{(q)}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy \right) \\ &\quad + \delta \mathbb{W}^{(q)}(x - b_\Lambda)(g_\Lambda(b_\Lambda)W^{(q)'}(b_\Lambda) - \Lambda qW^{(q)}(b_\Lambda) - 1), \end{aligned} \quad (\text{A.4})$$

where the last equality holds by integration by parts. Now, by the definition of g_Λ as in (4.13), we have

$$\begin{aligned} v_\Lambda^{b_\Lambda'}(x) &= g_\Lambda(b_\Lambda) \left(W^{(q)'}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)''}(y) dy \right) \\ &\quad - q\Lambda \left(W^{(q)}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy \right). \end{aligned} \quad (\text{A.5})$$

Differentiating this further, we obtain, for $x \neq b_\Lambda$,

$$\begin{aligned} v_\Lambda^{b_\Lambda''}(x) &= g_\Lambda(b_\Lambda) \left((1 + \delta \mathbb{W}^{(q)}(0) \mathbf{1}_{\{x > b_\Lambda\}}) W^{(q)''}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)'}(x-y)W^{(q)''}(y) dy \right) \\ &\quad - q\Lambda \left((1 + \delta \mathbb{W}^{(q)}(0) \mathbf{1}_{\{x > b_\Lambda\}}) W^{(q)'}(x) + \delta \int_{b_\Lambda}^x \mathbb{W}^{(q)'}(x-y)W^{(q)'}(y) dy \right). \end{aligned} \quad (\text{A.6})$$

By Remark 4.1, (A.5), and (A.6), the functions $v_\Lambda^{b_\Lambda'}$ and $v_\Lambda^{b_\Lambda''}$ are continuous on $\mathbb{R} \setminus \{b_\Lambda\}$. Regarding the continuity at b_Λ , from (A.5) we have $v_\Lambda^{b_\Lambda'}(b_\Lambda +) = v_\Lambda^{b_\Lambda'}(b_\Lambda -)$. In particular, for the case that X is of unbounded variation (where $\mathbb{W}^{(q)}(0) = 0$ as in (3.5)), we have, using (A.6),

$$v_\Lambda^{b_\Lambda''}(b_\Lambda +) - v_\Lambda^{b_\Lambda''}(b_\Lambda -) = 0.$$

(ii) For the case $b_\Lambda = 0$, the result follows by a direct application of Lemma 4.1.

A.3. Proof of Theorem 4.1

In order to verify inequality (4.20), we will need the following preliminary results.

Lemma A.1. ([14, Theorem 2].) *Under Assumption 4.1, the q -scale function $\mathbb{W}^{(q)}$ can be written as*

$$\mathbb{W}^{(q)}(x) = \varphi'(q)e^{\varphi(q)x} - \hat{f}(x), \tag{A.7}$$

where \hat{f} is a nonnegative, completely monotone function given by $\hat{f}(x) = \int_{0+}^{\infty} e^{-xt} \hat{\mu}(dt)$, where $\hat{\mu}$ is a finite measure on $(0, \infty)$. Moreover, $\mathbb{W}^{(q) \prime}$ is strictly log-convex (and hence convex) on $(0, \infty)$.

Remark A.1. We note that a result analogous to Lemma A.1 holds for $W^{(q)}$, with f and μ playing the role of \hat{f} and $\hat{\mu}$.

The following result will be crucial for the proof of Theorem 4.1.

Lemma A.2. *For $x > b_{\Lambda}$, we have*

$$v_{\Lambda}^{b_{\Lambda} \prime}(x) = \int_{0+}^{\infty} e^{-tx} l(t) \hat{\mu}(dt), \tag{A.8}$$

where

$$\begin{aligned} l(t) := & g_{\Lambda}(b_{\Lambda}) \left((1 - \delta W^{(q)}(0))t - \delta t \int_0^{b_{\Lambda}} e^{ty} W^{(q) \prime}(y) dy \right) \\ & + q\Lambda \left(1 + \delta t \int_0^{b_{\Lambda}} e^{ty} W^{(q)}(y) dy \right) + \delta e^{tb_{\Lambda}}. \end{aligned} \tag{A.9}$$

Proof. By integration by parts applied to (A.3),

$$\delta \int_0^x \mathbb{W}^{(q) \prime}(x-y) W^{(q)}(y) dy = \mathbb{W}^{(q)}(x) - (1 + \delta \mathbb{W}^{(q)}(0)) W^{(q)}(x). \tag{A.10}$$

On the other hand, by differentiating (A.3), we obtain

$$\delta \int_0^x \mathbb{W}^{(q) \prime \prime}(x-y) W^{(q) \prime}(y) dy = (1 - \delta W^{(q)}(0)) \mathbb{W}^{(q) \prime \prime}(x) - (1 + \delta \mathbb{W}^{(q)}(0)) W^{(q) \prime \prime}(x). \tag{A.11}$$

Applying (A.10) and (A.11) in the first equality in (A.4) it follows that, for $x > b_{\Lambda}$,

$$\begin{aligned} v_{\Lambda}^{b_{\Lambda} \prime}(x) = & g_{\Lambda}(b_{\Lambda}) \left((1 - \delta W^{(q)}(0)) \mathbb{W}^{(q) \prime}(x) - \delta \int_0^{b_{\Lambda}} \mathbb{W}^{(q) \prime}(x-y) W^{(q) \prime}(y) dy \right) \\ & - q\Lambda \left(\mathbb{W}^{(q)}(x) - \delta \int_0^{b_{\Lambda}} \mathbb{W}^{(q) \prime}(x-y) W^{(q)}(y) dy \right) - \delta \mathbb{W}^{(q)}(x - b_{\Lambda}). \end{aligned}$$

By (A.7), we can write

$$v_{\Lambda}^{b_{\Lambda} \prime}(x) = G_1(x) + G_2(x),$$

where

$$G_1(x) := g_\Lambda(b_\Lambda)\varphi'(q)\varphi(q)e^{\varphi(q)x} \left((1 - \delta W^{(q)}(0)) - \delta \int_0^{b_\Lambda} e^{-\varphi(q)y} W^{(q)'}(y) dy \right) \\ - q\Lambda\varphi'(q)e^{\varphi(q)x} \left(1 - \delta\varphi(q) \int_0^{b_\Lambda} e^{-\varphi(q)y} W^{(q)}(y) dy \right) - \varphi'(q)\delta e^{\varphi(q)(x-b_\Lambda)}, \\ G_2(x) := -g_\Lambda(b_\Lambda) \left((1 - \delta W^{(q)}(0))\hat{f}'(x) - \delta \int_0^{b_\Lambda} \hat{f}'(x-y)W^{(q)'}(y) dy \right) \\ + q\Lambda \left(\hat{f}(x) - \delta \int_0^{b_\Lambda} \hat{f}'(x-y)W^{(q)}(y) dy \right) + \delta\hat{f}(x-b_\Lambda),$$

with \hat{f} as in Remark A.1. Now we note that, by (A.1) and (A.2),

$$\int_0^{b_\Lambda} e^{-\varphi(q)y} W^{(q)'}(y) dy = \frac{1}{\delta} - W^{(q)}(0) - \int_{b_\Lambda}^\infty e^{-\varphi(q)y} W^{(q)'}(y) dy, \\ \int_0^{b_\Lambda} e^{-\varphi(q)y} W^{(q)}(y) dy = \frac{1}{\delta\varphi(q)} - \int_{b_\Lambda}^\infty e^{-\varphi(q)y} W^{(q)}(y) dy.$$

Combining these, we obtain

$$G_1(x) = g_\Lambda(b_\Lambda)\varphi'(q)\varphi(q)e^{\varphi(q)x} \delta \int_{b_\Lambda}^\infty e^{-\varphi(q)y} W^{(q)'}(y) dy \\ - q\Lambda\varphi'(q)e^{\varphi(q)x} \delta\varphi(q) \int_{b_\Lambda}^\infty e^{-\varphi(q)y} W^{(q)}(y) dy - \varphi'(q)\delta e^{\varphi(q)(x-b_\Lambda)} \\ = \varphi'(q)e^{\varphi(q)(x-b_\Lambda)} \delta \left(\xi_\Lambda(b_\Lambda)h(b_\Lambda) - q\Lambda\varphi(q)e^{\varphi(q)b_\Lambda} \int_{b_\Lambda}^\infty e^{-\varphi(q)y} W^{(q)}(y) dy \right) \\ - \varphi'(q)\delta e^{\varphi(q)(x-b_\Lambda)} \\ = 0.$$

Now, using the fact that $\hat{f}'(x) = -\int_{0+}^\infty te^{-xt}\hat{\mu}(dt)$ and Tonelli's theorem, we have (A.8). \square

A.3.1. *Proof of Theorem 4.1.* By Lemmas 4.3–4.5, it is sufficient to verify (4.20), and the condition that $v_\Lambda^{b_\Lambda}(0) \geq -\Lambda$.

(i) First, consider the case $b_\Lambda > 0$. In this case, recall that $\xi_\Lambda(b_\Lambda) = g_\Lambda(b_\Lambda)$ (implied by $\xi'_\Lambda(b_\Lambda) = 0$ and the expression (4.12)).

1. Suppose that $x \leq b_\Lambda$. Since g_Λ is increasing on $(0, a_\Lambda)$ by Remark 4.5(ii) and $b_\Lambda \leq a_\Lambda$ by Proposition 4.2(i), we obtain

$$\xi_\Lambda(b_\Lambda) = g_\Lambda(b_\Lambda) \geq g_\Lambda(x) = \frac{1 + q\Lambda W^{(q)}(x)}{W^{(q)'}(x)} \quad \text{for all } 0 < x \leq b_\Lambda. \quad (\text{A.12})$$

Applying (A.12) in (A.5), it follows that $v_\Lambda^{b_\Lambda'}(x) \geq 1$.

2. Now suppose that $x > b_\Lambda$. Differentiating (A.9) twice, we have

$$l''(t) = -\delta g_\Lambda(b_\Lambda) \left(2 \int_0^{b_\Lambda} ye^{t,y} W^{(q)'}(y) dy + t \int_0^{b_\Lambda} y^2 e^{t,y} W^{(q)'}(y) dy \right) + \delta b_\Lambda^2 e^{tb_\Lambda} \\ + \delta q\Lambda \left(2 \int_0^{b_\Lambda} ye^{t,y} W^{(q)}(y) dy + t \int_0^{b_\Lambda} y^2 e^{t,y} W^{(q)}(y) dy \right).$$

On the other hand, using (A.12), we have $g_\Lambda(b_\Lambda)W^{(q)'}(y) \geq 1 + q\Lambda W^{(q)}(y)$ for all $y \in (0, b_\Lambda]$. Hence,

$$\begin{aligned} l''(t) &\leq -\delta \left(2 \int_0^{b_\Lambda} ye^{ty}(1 + q\Lambda W^{(q)}(y)) dy + t \int_0^{b_\Lambda} y^2e^{ty}(1 + q\Lambda W^{(q)}(y)) dy \right) + \delta b_\Lambda^2 e^{tb_\Lambda} \\ &\quad + \delta q\Lambda \left(2 \int_0^{b_\Lambda} ye^{ty}W^{(q)}(y) dy + t \int_0^{b_\Lambda} y^2e^{ty}W^{(q)}(y) dy \right) \\ &= -\delta \left(2 \int_0^{b_\Lambda} ye^{ty}dy + t \int_0^{b_\Lambda} y^2e^{ty}dy \right) + \delta b_\Lambda^2 e^{tb_\Lambda} \\ &= 0. \end{aligned}$$

Therefore, l is a concave function. In addition, since $l(0) = q\Lambda + \delta$, which is positive by Assumption 4.2, and recalling that $x > b_\Lambda$, it follows that there exists $0 < p \leq \infty$ such that l is positive on $(0, p)$ and negative on (p, ∞) . Consequently,

$$e^{-(x-b_\Lambda)t}l(t) \geq e^{-(x-b_\Lambda)p}l(t), \quad t > 0. \tag{A.13}$$

Now we note from (A.9) that there exists a constant $C(b_\Lambda)$ independent of t such that $|l(t)| \leq C(b_\Lambda)(1 + t)e^{tb_\Lambda}$. Therefore, using the fact that $x > b_\Lambda$ and the dominated convergence, we can take the derivative inside the integral in (A.8) and obtain

$$\begin{aligned} v_\Lambda^{b_\Lambda''}(x) &= - \int_{0+}^\infty e^{-(x-b_\Lambda)t}e^{-b_\Lambda t}tl(t)\hat{\mu}(dt) \\ &\leq -e^{-(x-b_\Lambda)p} \int_{0+}^\infty e^{-b_\Lambda t}tl(t)\hat{\mu}(dt) \\ &= e^{-(x-b_\Lambda)p}v_\Lambda^{b_\Lambda''}(b_\Lambda), \end{aligned} \tag{A.14}$$

where the inequality holds by (A.13). On the other hand, Proposition 4.2 implies that $b_\Lambda \leq a_\Lambda$, and hence, by Remark 4.5(ii),

$$0 \leq g'_\Lambda(b_\Lambda) = q\Lambda - \frac{W^{(q)''}(b_\Lambda)}{W^{(q)'}(b_\Lambda)}g_\Lambda(b_\Lambda).$$

Therefore, (A.6) gives

$$v_\Lambda^{b_\Lambda''}(b_\Lambda+) = (1 + \delta W^{(q)}(0))(g_\Lambda(b_\Lambda)W^{(q)''}(b_\Lambda) - q\Lambda W^{(q)'}(b_\Lambda)) \leq 0.$$

In combination with (A.14), it follows that $v_\Lambda^{b_\Lambda'}$ is nonincreasing on (b_Λ, ∞) . In addition, we note using (4.13) and (A.5) that

$$v_\Lambda^{b_\Lambda'}(b_\Lambda) = g_\Lambda(b_\Lambda)W^{(q)'}(b_\Lambda) - q\Lambda W^{(q)}(b_\Lambda) = 1.$$

Hence, we deduce that $v_\Lambda^{b_\Lambda'}(x) \leq 1$ for $x > b_\Lambda$.

3. Finally, using (4.18), and the fact that $g_\Lambda(b_\Lambda) \geq 0$,

$$v_\Lambda^{b_\Lambda}(0) = g_\Lambda(b_\Lambda)W^{(q)}(0) - \Lambda \geq -\Lambda,$$

as required.

(ii) Now, consider the case $b_\Lambda = 0$. By taking a derivative in (4.11), and using (4.10) and (A.7), we obtain

$$\begin{aligned} v_\Lambda^{0''}(x) &= (\delta + q\Lambda) \left(\frac{\mathbb{W}^{(q)'}(x)}{\varphi(q)} - \mathbb{W}^{(q)}(x) \right)' \\ &= (\delta + q\Lambda) \left(-\frac{\hat{f}'(x)}{\varphi(q)} + \hat{f}(x) \right)' \\ &= (\delta + q\Lambda) \left(-\frac{\hat{f}''(x)}{\varphi(q)} + \hat{f}'(x) \right), \end{aligned}$$

which is negative because \hat{f} is completely monotone. Therefore, $v_\Lambda^{0'}(x)$ is nonincreasing, and hence, it is enough to verify that $v_\Lambda^{0'}(0+) \leq 1$ or, equivalently,

$$\frac{(\delta + q\Lambda)\mathbb{W}^{(q)'}(0+)}{1 + (\delta + q\Lambda)\mathbb{W}^{(q)}(0)} \leq \varphi(q).$$

This inequality is automatically satisfied in cases 1 and 2 given in Proposition 4.2(ii). Therefore, we have (4.20) when $b_\Lambda = 0$. To complete the proof, using (4.10), (4.11), and Assumption 4.2, we have

$$v_\Lambda^0(0) = \frac{(\delta + q\Lambda)}{\varphi(q)} \mathbb{W}^{(q)}(0) - \Lambda \geq -\Lambda.$$

A.4. Proof of Lemma 5.1

First we note that, using (4.3),

$$h'(b) = \varphi(q)(h(b) - W^{(q)'}(b)). \tag{A.15}$$

Hence,

$$h'(b)H(b) = \varphi(q)W^{(q)}(b) \left(\frac{W^{(q)'}(b)}{W^{(q)}(b)} \left(W^{(q)}(b) + \frac{h(b)}{\varphi(q)} \right) - h(b) \right). \tag{A.16}$$

Now, since $h' > 0$ on (b_0, ∞) (see Remark 4.4(i)), it is enough to show that the right-hand side of (A.16) is positive. On the other hand, we know from [8] that $W^{(q)}$ is log-concave on $(0, a_0]$ and strictly log-concave on (a_0, ∞) , where a_0 is defined in Remark 4.5 for the case $\Lambda = 0$. Then

$$\frac{W^{(q)'(\eta)}}{W^{(q)}(\eta)} \geq \frac{W^{(q)'(\zeta)}}{W^{(q)}(\zeta)} \quad \text{for any } \eta \text{ and } \zeta \text{ with } 0 < \eta \leq \zeta.$$

Note that the previous inequality is strict when $a_0 < \eta < \zeta$. From here, it can be verified that

$$\frac{W^{(q)'(b)}}{W^{(q)}(b)} \int_b^\infty e^{-\varphi(q)y} W^{(q)}(y) dy > \int_b^\infty e^{-\varphi(q)y} W^{(q)'(y)} dy,$$

and using (4.3) and (4.8), it follows that

$$\frac{W^{(q)'(b)}}{W^{(q)}(b)} \left(W^{(q)}(b) + \frac{h(b)}{\varphi(q)} \right) > h(b) \quad \text{for all } b > 0. \tag{A.17}$$

From (A.16) and (A.17), we have $h'(b)H(b) > 0$ for $b \in (b_0, \infty)$, as desired.

Now, taking the first derivative in (5.3) and by (A.15), we have

$$\begin{aligned} H'(b) &= h(b) - \frac{[h(b)]^2 h''(b)}{[h'(b)]^2} \\ &= \frac{h(b)}{[h'(b)]^2} ([h'(b)]^2 - h(b)h''(b)) \\ &= \frac{[\varphi(q)h(b)]^2 W^{(q)'}(b)}{[h'(b)]^2} \left(\frac{W^{(q)'}(b)}{h(b)} + \frac{W^{(q)''}(b)}{\varphi(q)W^{(q)'}(b)} - 1 \right), \end{aligned}$$

where the last equality holds because

$$\begin{aligned} &[h'(b)]^2 - h(b)h''(b) \\ &= \varphi^2(q)(h(b) - W^{(q)'}(b))^2 - h(b)\varphi(q)[\varphi(q)(h(b) - W^{(q)'}(b)) - W^{(q)''}(b)] \\ &= \varphi(q)\{\varphi(q)(W^{(q)'}(b))^2 - h(b)W^{(q)'}(b) + h(b)W^{(q)''}(b)\}. \end{aligned}$$

Since, by Remark A.1, $W^{(q)'}$ is a strictly log-convex function, we have

$$\frac{W^{(q)''}(\eta)}{W^{(q)'}(\eta)} < \frac{W^{(q)''}(\zeta)}{W^{(q)'}(\zeta)} \quad \text{for any } \eta \text{ and } \zeta \text{ with } 0 < \eta < \zeta.$$

From the above and integration by parts we can show that

$$\frac{W^{(q)''}(b)}{W^{(q)'}(b)} h(b) < \varphi(q)e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)''}(y) dy = -\varphi(q)W^{(q)'}(b) + \varphi(q)h(b),$$

and hence,

$$\frac{W^{(q)'}(b)}{h(b)} + \frac{W^{(q)''}(b)}{\varphi(q)W^{(q)'}(b)} - 1 < \frac{W^{(q)'}(b)}{h(b)} + \frac{h(b) - W^{(q)'}(b)}{h(b)} - 1 = 0.$$

Hence, we conclude that the function H as in (5.3) is strictly decreasing.

A.5. Proof of Proposition 5.1

(i) For the case $\bar{\Lambda} \geq 0$, Remark 4.6 gives $b_0 = 0$. Then, from (4.10), (5.3), and (A.15),

$$\begin{aligned} \lim_{b \rightarrow 0} H(b) &= \frac{[h(0)]^2}{h'(0+)} - \left(W^{(q)}(0) + \frac{h(0)}{\varphi(q)} \right) \\ &= \frac{h(0)W^{(q)'}(0+) - \varphi(q)W^{(q)}(0)(h(0) - W^{(q)'}(0+))}{\varphi(q)(h(0) - W^{(q)'}(0+))} \\ &= \frac{\delta^{-1}W^{(q)'}(0+) - \varphi(q)W^{(q)}(0)(\delta^{-1} - W^{(q)}(0))}{\varphi(q)(\delta^{-1} - W^{(q)}(0)) - W^{(q)'}(0+)}. \end{aligned}$$

Using (3.5), (3.6), and the fact that $\phi_1(\bar{\Lambda}) = \varphi(q)$ or $\phi_2(\bar{\Lambda}) = \varphi(q)$ (see Remark 4.6), it can be verified that $\lim_{b \rightarrow 0} H(b) = 1/q\bar{\Lambda}$ (where in the case $\bar{\Lambda} = 0$ the right-hand side is understood to be ∞), and hence, $\lim_{b \downarrow b_0} \lambda(b) = \lim_{b \downarrow 0} \lambda(b) = \bar{\Lambda}$.

For the case $\bar{\Lambda} < 0$, Remark 4.6 gives $b_0 > 0$. By Lemma 3 of [12], we know that h attains its unique minimum at b_0 and, by the continuity of h' , $\lim_{b \rightarrow b_0} h'(b) = h'(b_0) = 0$.

In addition, by (A.15), $\lim_{b \rightarrow b_0} h(b) = W^{(q)'}(b_0) > 0$. Therefore, from (5.3), we obtain $\lim_{b \rightarrow b_0} H(b) = \infty$, and, hence, $\lim_{b \downarrow b_0} \lambda(b) = 0 = \bar{\Lambda} \vee 0$.

(ii) From Remark A.1, we can write

$$h(b) = \varphi(q) \left(\frac{\Phi(q)\Phi'(q)e^{\Phi(q)b}}{\varphi(q) - \Phi(q)} - \tilde{f}(b) \right),$$

where $\tilde{f}(b) := \int_0^\infty e^{-\varphi(q)y} f'(y + b) dy$, and hence we get the following expressions:

$$[h(b)]^2 = [\varphi(q)]^2 \left(\frac{[\Phi(q)\Phi'(q)]^2}{(\varphi(q) - \Phi(q))^2} e^{2\Phi(q)b} - \frac{2\Phi(q)\Phi'(q)}{\varphi(q) - \Phi(q)} e^{\Phi(q)b} \tilde{f}(b) + [\tilde{f}(b)]^2 \right),$$

$$h'(b) = \varphi(q) \left(\frac{[\Phi(q)]^2 \Phi'(q)}{\varphi(q) - \Phi(q)} e^{\Phi(q)b} - \varphi(q) \tilde{f}(b) + f'(b) \right),$$

and

$$W^{(q)}(b) + \frac{h(b)}{\varphi(q)} = \frac{\Phi'(q)\varphi(q)}{\varphi(q) - \Phi(q)} e^{\Phi(q)b} - \tilde{f}(b) - f(b).$$

Applying these identities in (5.3), it follows that

$$H(b) = \frac{\varphi(q)\Phi(q)\Phi'(q)}{\varphi(q) - \Phi(q)} H_1(b) + \frac{\varphi(q)^2 [\tilde{f}(b)]^2}{h'(b)} + \tilde{f}(b) + f(b),$$

where

$$H_1(b) := e^{\Phi(q)b} \left[\frac{\varphi(q)}{h'(b)} \left(\frac{\Phi(q)\Phi'(q)}{\varphi(q) - \Phi(q)} e^{\Phi(q)b} - 2\tilde{f}(b) \right) - \frac{1}{\Phi(q)} \right]$$

$$= \frac{\varphi(q)}{\Phi(q)} \frac{\varphi(q)\tilde{f}(b) - 2\Phi(q)\tilde{f}(b) - f'(b)}{h'(b)}.$$

By the dominated convergence theorem we have $f(b) \rightarrow 0, \tilde{f}(b) \rightarrow 0$, and $f'(b) \rightarrow 0$ as $b \rightarrow \infty$. Hence $\lim_{b \rightarrow \infty} H(b) = 0$ or, equivalently, $\lim_{b \rightarrow \infty} \lambda(b) = \infty$.

(iii) Fix $b > b_0$. First let us assume that $b_0 > 0$ or that $b_0 = 0$ and $h'(0+) = 0$. Then, using (4.12) for $\Lambda = \lambda(b)$, we have

$$\frac{d\xi_{\lambda(b)}(\zeta)}{d\zeta} = q\lambda(b) - \frac{h'(\zeta)}{h(\zeta)} \xi_{\lambda(b)}(\zeta) \quad \text{for } \zeta > 0. \tag{A.18}$$

For $\zeta \in (0, b_0]$, by the fact that $h'(\zeta) \leq 0$ for $\zeta \in [0, b_0]$ (see Remark 4.4(i)) we have that $d\xi_{\lambda(b)}(\zeta)/d\zeta > 0$. On the other hand, applying (A.15) and (4.9) in (A.18),

$$\frac{d\xi_{\lambda(b)}(\zeta)}{d\zeta} = \frac{1}{[h(\zeta)]^2} \left(q\lambda(b) \left([h(\zeta)]^2 - h'(\zeta) \left(W^{(q)}(\zeta) + \frac{h(\zeta)}{\varphi(q)} \right) \right) - h'(\zeta) \right)$$

$$= \frac{h'(\zeta)}{[h(\zeta)]^2} \left(\frac{\lambda(b)}{\lambda(\zeta)} - 1 \right). \tag{A.19}$$

Then, using the fact that $\zeta \mapsto \lambda(\zeta)$ is strictly increasing as in (i), and that $h'(\zeta) > 0$ for $\zeta > b_0$ (see Remark 4.4(i)), we have $d\xi_{\lambda(b)}(\zeta)/d\zeta > 0$ for $b_0 < \zeta < b$ and vanishes at $\zeta = b$. Now let us assume that $b_0 = 0$ and that $h'(0+) > 0$. Then, by the proof of Lemma 3 of [12], we have $h'(\zeta) > 0$ for all $\zeta > 0$. Hence, (A.19) implies that $d\xi_{\lambda(b)}(\zeta)/d\zeta > 0$ for $0 < \zeta < b$ and vanishes at $\zeta = b$. The above implies that $b = \inf\{\zeta \geq 0 : d\xi_{\lambda(b)}(\zeta)/d\zeta \leq 0\}$. Therefore, by (4.14) we obtain $b_{\lambda(b)} = b$ for all $b > b_0$.

A.6. Proof of Lemma 5.2.

1. We have, by (A.15),

$$\frac{d}{db} \frac{W^{(q)}(b)}{h(b)} = \frac{\alpha(b)}{[h(b)]^2}, \quad b > 0, \tag{A.20}$$

with

$$\alpha(b) := W^{(q)'}(b)h(b) - W^{(q)}(b)\varphi(q)(h(b) - W^{(q)'}(b)) > 0, \quad b > 0,$$

where the positivity holds by (A.17).

If $x \leq b$, we have, by (5.4),

$$\Psi_x(b) = Z^{(q)}(x) - \frac{q(W^{(q)}(b) + h(b)/\varphi(q))}{h(b)} W^{(q)}(x). \tag{A.21}$$

Taking the derivative with respect to b , by (A.20) and the positivity of α ,

$$\frac{d\Psi_x(b)}{db} = -\frac{qW^{(q)}(x)}{[h(b)]^2} \alpha(b) < 0.$$

Therefore, Ψ_x is strictly decreasing on $[x, \infty)$.

Suppose that $b < x$. By (A.20) and (5.4),

$$\begin{aligned} \frac{d\Psi_x(b)}{db} &= -\delta q \mathbb{W}^{(q)}(x-b)W^{(q)}(b) \\ &\quad + \frac{q\delta \mathbb{W}^{(q)}(x-b)W^{(q)'}(b)(W^{(q)}(b) + h(b)/\varphi(q))}{h(b)} \\ &\quad - q \left(W^{(q)}(x) + \delta \int_b^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y)dy \right) \frac{d}{db} \left[\frac{W^{(q)}(b) + h(b)/\varphi(q)}{h(b)} \right] \\ &= \frac{q}{[h(b)]^2} r(b; x)\alpha(b), \end{aligned} \tag{A.22}$$

where, for $x, b \geq 0$,

$$r(b; x) := \frac{\delta \mathbb{W}^{(q)}(x-b)h(b)}{\varphi(q)} - W^{(q)}(x) - \delta \int_b^x \mathbb{W}(x-y)W^{(q)'}(y) dy.$$

To prove that $d\Psi_x/db < 0$ on $(0, x)$, by the positivity of α , we only need to verify that $r(b; x) < 0$ for all $b \in (0, x)$.

Note that, by (4.8),

$$\frac{\delta \mathbb{W}^{(q)}(x-b)h(b)}{\varphi(q)} = \delta \mathbb{W}^{(q)}(x-b) \left(\varphi(q)e^{\varphi(q)b} \int_b^\infty e^{-\varphi(q)y} W^{(q)}(y) dy - W^{(q)}(b) \right).$$

Using integration by parts and (A.10),

$$\begin{aligned} &-\delta \int_b^x \mathbb{W}^{(q)}(x-y)W^{(q)'}(y) dy \\ &= \delta \mathbb{W}^{(q)}(x-b)W^{(q)}(b) - \mathbb{W}^{(q)}(x) + W^{(q)}(x) + \delta \int_0^b \mathbb{W}^{(q)'}(x-y)W^{(q)}(y) dy. \end{aligned}$$

Substituting these,

$$r(b; x) = -\mathbb{W}^{(q)}(x) + \delta \int_0^b \mathbb{W}^{(q)'}(x - y)W^{(q)}(y) dy + \delta\varphi(q)e^{\varphi(q)b}\mathbb{W}^{(q)}(x - b) \int_b^\infty e^{-\varphi(q)y}W^{(q)}(y) dy.$$

Now, we rewrite this using Lemma A.1. By observing that the terms corresponding to $e^{\varphi(q)x}$ all cancel out (using the fact that $\int_0^\infty e^{-\varphi(q)y}W^{(q)}(y) dy = (\delta\varphi(q))^{-1}$), it follows that

$$\begin{aligned} r(b; x) &= \hat{f}(x) - \delta \int_0^b \hat{f}'(x - y)W^{(q)}(y) dy - \delta\varphi(q)e^{\varphi(q)b}\hat{f}(x - b) \int_b^\infty e^{-\varphi(q)y}W^{(q)}(y) dy \\ &= \int_0^\infty e^{-xt} \left(1 + \delta t \int_0^b e^{yt}W^{(q)}(y) dy - \delta\varphi(q)e^{b(t+\varphi(q))} \int_b^\infty e^{-\varphi(q)y}W^{(q)}(y) dy \right) \hat{\mu}(dt). \end{aligned} \tag{A.23}$$

Taking the derivative with respect to b in (A.23), it follows that

$$\begin{aligned} \frac{\partial r(b; x)}{\partial b} &= \int_0^\infty e^{-xt} \left(\delta(t + \varphi(q))e^{bt}W^{(q)}(b) - \delta\varphi(q)(t + \varphi(q))e^{b(t+\varphi(q))} \int_b^\infty e^{-\varphi(q)y}W^{(q)}(y) dy \right) \hat{\mu}(dt) \\ &< \int_0^\infty e^{-xt} \left(\delta(t + \varphi(q))e^{bt}W^{(q)}(b) - \delta\varphi(q)(t + \varphi(q))e^{b(t+\varphi(q))}W^{(q)}(b) \int_b^\infty e^{-\varphi(q)y} dy \right) \hat{\mu}(dt) \\ &= 0, \end{aligned}$$

where the inequality follows since $W^{(q)}$ is strictly increasing on $(0, \infty)$. From here we conclude that $r(b; x)$ is strictly decreasing on $(0, x)$.

On the other hand, by (A.23) we obtain $\lim_{b \rightarrow 0} r(b; x) = 0$, and hence, $r(b; x) < 0$ on $(0, x)$. Now we conclude by (A.22) that $b \mapsto \Psi_x(b)$ is also strictly decreasing on $(0, x)$.

2. By (4.17), we see that

$$\lim_{b \rightarrow \infty} \frac{W^{(q)}(b) + h(b)/\varphi(q)}{h(b)} = \frac{1}{\Phi(q)}. \tag{A.24}$$

Then, letting $b \rightarrow \infty$ in (A.21) and using (A.24), we obtain the expression in (5.5).

A.7. Proof of Proposition 6.1

Consider U^{-b} the (spectrally negative) refracted Lévy process with refraction level $-b$, driven by the process X (as in Section 2), and $\eta_{-b} := \inf\{t > 0 : U_t^{-b} > 0\}$. Then

$$\begin{aligned} \bar{\Psi}_x(b) &= \mathbb{E}_{-x}[e^{-q\eta_{-b}}; \eta_{-b} < \infty], \\ \mathbb{E}_x \left[\int_0^{\bar{\tau}_b} e^{-qt} d\bar{D}_t^b \right] &= \delta \left(\frac{1}{q} (1 - \bar{\Psi}_x(b)) - \mathbb{E}_{-x} \left[\int_0^{\eta_{-b}} e^{-qt} \mathbf{1}_{\{U_t^{-b} > -b\}} dt \right] \right). \end{aligned}$$

Now from Theorem 5(i) of [11], we have (6.2). On the other hand, by Theorem 6(iii) of [11], we obtain

$$\mathbb{E}_{-x} \left[\int_0^{\eta-b} e^{-qt} \mathbf{1}_{\{U_t^- > -b\}} dt \right] = \bar{\Psi}_x(b) \bar{\mathbb{W}}^{(q)}(b) - \bar{\mathbb{W}}^{(q)}(b-x).$$

Hence, putting the pieces together we get (6.3).

A.8. Proof of Lemma 6.3

For $b \in (0, x)$, we have $\bar{\Psi}_x(b) = e^{-\Phi(q)x} / \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)$, which is clearly strictly decreasing.

To see that it is strictly decreasing on $[x, \infty)$, we need to show that $d\bar{\Psi}_x(b)/db < 0$ on (x, ∞) . This is satisfied if we can show that

$$w(y_1, y_2) > 0, \quad y_1 < y_2,$$

where we define, with y_2 fixed,

$$w(y, y_2) := \frac{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(y)}{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(y_2)} - \frac{\mathbb{W}_{\Phi(q)}^{(\delta\Phi(q))}(y)}{\mathbb{W}_{\Phi(q)}^{(\delta\Phi(q))}(y_2)} \quad \text{for any } y \in \mathbb{R}.$$

Recall the change of measure addressed in Section 6.1. Because, for $y > 0$,

$$\{e^{-\delta\Phi(q)(t \wedge \tau_{0,y})} \mathbb{W}_{\Phi(q)}^{(\delta\Phi(q))}(Y_{t \wedge \tau_{0,y}}) : t \geq 0\} \quad \text{and} \quad \{e^{-\delta\Phi(q)(t \wedge \tau_{0,y})} \mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(Y_{t \wedge \tau_{0,y}}) : t \geq 0\}$$

are $\tilde{\mathbb{P}}_x^{\Phi(q)}$ -martingales (see Proposition 3 of [19]), where $\tau_{0,y} := \inf\{t > 0 : Y_t < 0 \text{ or } Y_t > y\}$, it follows that $\{e^{-\delta\Phi(q)(t \wedge \tau_{0,y_2})} w(Y_{t \wedge \tau_{0,y_2}}, y_2) : t \geq 0\}$ is a $\tilde{\mathbb{P}}_x^{\Phi(q)}$ -martingale. Now, taking $y_1 < y_2$ and using the optimal stopping theorem,

$$w(y_1, y_2) = \tilde{\mathbb{E}}_{y_1}^{\Phi(q)} [e^{-q(t \wedge \tau_{0,y_2})} w(Y_{t \wedge \tau_{0,y_2}}, y_2)],$$

where $\tilde{\mathbb{E}}_{y_1}^{\Phi(q)}$ is the expected value with respect to the probability measure $\tilde{\mathbb{P}}_{y_1}^{\Phi(q)}$. Noting that w is bounded (recalling that Y is spectrally negative) and taking $t \rightarrow \infty$, dominated convergence gives

$$w(y_1, y_2) = \tilde{\mathbb{E}}_{y_1}^{\Phi(q)} [e^{-q\tau_{0,y_2}} w(Y_{\tau_{0,y_2}}, y_2)].$$

Now we note that the following assertions hold.

- (i) If Y has paths of unbounded variation then $\mathbb{W}_{\Phi(q)}^{(\delta\Phi(q))}(0) = 0$, and hence, $w(y, y_2) > 0$ for $y \in (-\infty, 0]$. On the other hand, $\tilde{\mathbb{P}}_x^{\Phi(q)}(Y_{\tau_{0,y_2}} \leq 0) > 0$.
- (ii) If Y has paths of bounded variation then $w(y, y_2) > 0$ for $y \in (-\infty, 0)$, and

$$\tilde{\mathbb{P}}_x^{\Phi(q)}(Y_{\tau_{0,y_2}} \leq 0) = \tilde{\mathbb{P}}_x^{\Phi(q)}(Y_{\tau_{0,y_2}} < 0) > 0.$$

These facts imply that

$$w(y_1, y_2) = \tilde{\mathbb{E}}_{y_1}^{\Phi(q)} [e^{-q\tau_{0,y_2}} w(Y_{\tau_{0,y_2}}, y_2)] \geq \tilde{\mathbb{E}}_{y_1}^{\Phi(q)} [e^{-q\tau_{0,y_2}} w(Y_{\tau_{0,y_2}}, y_2) \mathbf{1}_{\{Y_{\tau_{0,y_2}} \leq 0\}}] > 0.$$

From here we conclude that $\bar{\Psi}_x$ is strictly decreasing on $(0, \infty)$. Finally, letting $b \rightarrow 0$ in (6.2), it is clear that $\lim_{b \rightarrow 0} \bar{\Psi}_x(b) = e^{-\Phi(q)x}$. On the other hand, using l'Hôpital's rule and (3.7), we have

$$\lim_{b \rightarrow \infty} \frac{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b-x)}{\mathbb{Z}_{\Phi(q)}^{(\delta\Phi(q))}(b)} = \frac{1}{e^{(\varphi(q)-\Phi(q))x}} \lim_{b \rightarrow \infty} \frac{e^{-\varphi(q)(b-x)} \mathbb{W}_{\Phi(q)}(b-x)}{e^{-\varphi(q)b} \mathbb{W}_{\Phi(q)}(b)} = \frac{1}{e^{(\varphi(q)-\Phi(q))x}}.$$

Hence, $\lim_{b \rightarrow \infty} \bar{\Psi}_x(b) = e^{-\varphi(q)x}$.

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