

Lepton-hadronic processes and high-energy neutrinos in NGC 1275

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Abstract. The nearby active galaxy NGC 1275, has been widely detected from radio to gamma rays. Its spectral energy distribution (SED) shows a double-peak feature, which is well explained by synchrotron self-Compton (SSC) model. However, recent TeV detections might suggest that very-high-energy γ -rays ($E \geq 100$ GeV) may not have a leptonic origin. We test a lepto-hadronic model to describe the whole SED through SSC emission and neutral pion decay resulting from $p\gamma$ interactions. Also, we estimate the neutrino events expected in a km^3 Cherenkov telescope.

Keywords. galaxies: individual (NGC1275), radiation mechanisms: non-thermal

1. Introduction

NGC 1275, also known as Perseus A and 3C 84, is the nearby active galaxy located at the center of the Perseus cluster at redshift of $z = 0.0179$ (Veron (1978)). This object has been well studied in radio, optical/UV, X-ray and MeV - GeV emission (Abdo *et al.* 2009, Aleksić *et al.* 2012, Acciari *et al.* 2009). Recently, this source has been detected by MAGIC telescopes with a statistical significance of 6.6σ above 100 GeV in 46 hr of stereo observations carried out between August 2010 and February 2011. The measured differential energy spectrum between 70 GeV and 500 GeV can be described by a power law with a steep spectral index of $\Gamma = -4.1 \pm 0.7_{stat} \pm 0.3_{syst}$, and the average flux above 100 GeV, $F_\gamma = 1.3 \pm 0.2_{stat} \pm 0.3_{syst} \times 10^{-11} \text{ cm}^{-2} \text{ s}^{-1}$ (Aleksić *et al.* (2012)). We propose a lepto-hadronic model to describe the whole SED as the superposition of synchrotron self-Compton (SSC) emission and the neutral pion decay products.

2. Theoretical model

Fermi-accelerated electrons in the emitting region are described by a broken power-law written as: $N_e(\gamma_e) \propto \gamma_e^{-\alpha_e}$ for $\gamma_{e,m} < \gamma_e < \gamma_{e,b}$, and $N_e(\gamma_e) \propto \gamma_{e,b}\gamma_e^{-(\alpha_e+1)}$ for $\gamma_{e,b} \leq \gamma_e < \gamma_{e,max}$ and co-accelerated protons at the same place for a simple power law as: $N_p(\gamma_p) \propto \gamma_p^{-\alpha_p}$, where α_e (α_p) is the power index for electrons (protons) and $\gamma_{e,i}$ (γ_p) are the electron (proton) Lorentz factors. The indices i , c and max are for minimum, break, and maximum, respectively. Assuming that the emitting region is endowed by a magnetic field, electrons and protons are cooled down by electromagnetic (synchrotron radiation and Compton scattering) and hadronic (proton-photon interaction) channels, respectively. Taking into account the cooling processes for electron and proton distributions, the observed spectra for electron synchrotron radiation, Compton scattering and

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| Parameter | Value | Parameter | Value |
|---|-------------------------------------|---|------------------------------------|
| $A_{syn,\gamma}$ (MeVcm ⁻² s ⁻¹) | $(6.248 \pm 0.899) \times 10^{-5}$ | $A_{ssc,\gamma}$ (MeVcm ⁻² s ⁻¹) | $(1.806 \pm 0.754) \times 10^{-5}$ |
| α_e | 2.809 ± 0.0520 | α_e | 2.809 ± 0.0520 |
| $\epsilon_{\gamma,c}^{syn}$ (eV) | 0.100 ± 0.001 | $\epsilon_{\gamma,c}^{ssc}$ (MeV) | 7.254 ± 2.574 |
| $\epsilon_{\gamma,m}^{syn}$ (eV) | $(0.001 \pm 9.308) \times 10^{-05}$ | $\epsilon_{\gamma,m}^{ssc}$ (keV) | (11.374 ± 3.696) |

Table 1. Best fit electron synchrotron radiation (left) and Compton scattering (right) parameters obtained after fitting the SED.

proton-photon interaction are (Fraija 2014a, Fraija *et al.* 2012):

$$(\epsilon^2 N(\epsilon))_{\gamma,syn} = A_{syn,\gamma} \begin{cases} \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}^{syn}}\right)^{4/3} & \epsilon_\gamma < \epsilon_{\gamma,m}^{syn} \\ \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}^{syn}}\right)^{-(\alpha-3)/2} & \epsilon_{\gamma,m}^{syn} < \epsilon_\gamma < \epsilon_{\gamma,c}^{syn} \\ \left(\frac{\epsilon_{\gamma,c}}{\epsilon_{\gamma,m}^{syn}}\right)^{-(\alpha-3)/2} \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,c}^{syn}}\right)^{-(\alpha-2)/2} & \epsilon_{\gamma,c}^{syn} < \epsilon_\gamma < \epsilon_{\gamma,max}^{syn} \end{cases}, \tag{2.1}$$

$$(\epsilon^2 N(\epsilon))_{\gamma,ssc} = A_{ssc,\gamma} \begin{cases} \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}^{ssc}}\right)^{4/3} & \epsilon_\gamma < \epsilon_{\gamma,m}^{ssc} \\ \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}^{ssc}}\right)^{-(\alpha-3)/2} & \epsilon_{\gamma,m}^{ssc} < \epsilon_\gamma < \epsilon_{\gamma,c}^{ssc} \\ \left(\frac{\epsilon_{\gamma,c}}{\epsilon_{\gamma,m}^{ssc}}\right)^{-(\alpha-3)/2} \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma,c}^{ssc}}\right)^{-(\alpha-2)/2} & \epsilon_{\gamma,c}^{ssc} < \epsilon_\gamma < \epsilon_{\gamma,max}^{ssc} \end{cases}, \tag{2.2}$$

$$(\epsilon^2 N(\epsilon))_{\pi^0,\gamma} = A_{p,\gamma} \begin{cases} \left(\frac{\epsilon_\gamma}{\epsilon_0}\right)^{-1} \left(\frac{\epsilon_{\gamma,c,\pi^0}}{\epsilon_0}\right)^{-\alpha_p+3} & \epsilon_\gamma < \epsilon_{\gamma,c,\pi^0} \\ \left(\frac{\epsilon_\gamma}{\epsilon_0}\right)^{-\alpha_p+2} & \epsilon_{\gamma,c,\pi^0} < \epsilon_\gamma \end{cases}, \tag{2.3}$$

respectively, where $A_{syn,\gamma}$ ($A_{ssc,\gamma}$), $\epsilon_{\gamma,m}^{syn}$ ($\epsilon_{\gamma,m}^{ssc}$), $\epsilon_{\gamma,c}^{syn}$ ($\epsilon_{\gamma,c}^{ssc}$) and $\epsilon_{\gamma,max}^{syn}$ ($\epsilon_{\gamma,max}^{ssc}$) are the proportionality constant and break energies for characteristic, cut-off and maximum of electron synchrotron (Compton scattering) spectrum and $A_{p,\gamma}$ and $\epsilon_{\gamma,c,\pi^0}$ are the proportionality constant and break photon energy for photo-pion spectrum (Fraija 2014b). It is important to highlight that the neutrino counterpart is calculated through $\int \frac{dN_\nu}{dE_\nu} E_\nu dE_\nu = \frac{1}{4} \int \left(\frac{dN}{d\epsilon}\right)_{\pi^0,\gamma} \epsilon_{\pi^0,\gamma} d\epsilon_{\pi^0,\gamma}$, where the neutrino flux is $dN_\nu/dE_\nu = A_\nu E_\nu^{-\alpha_\nu}$, with $\alpha_\nu \simeq \alpha_p$.

3. Results

We have presented a lepton-hadronic model through Eqns. 2.2, 2.2 and 2.3 to describe the broadband SED of NGC1275. For the hadronic interactions we invoke the $p\gamma$ interactions of Fermi-accelerated protons with the target photons at the second SSC peak. To find the best fit of our leptonic and hadronic models, we use the chi-square (χ^2) minimization method. The fit values for the leptonic model are shown in Table 1, while for hadronic model are $A_{p\gamma} = (2.932 \pm 0.6579) \times 10^{-6}$ MeV cm⁻² s⁻¹ and $\alpha_p = 3.467 \pm 0.090$. The leptonic model describes the photon spectrum up to a few GeVs while the hadronic model up to hundreds of GeVs. The neutrino event expected in a km³ Cherenkov telescope is 0.93×10^{-4} per year.

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