# Existence results of solitons in discrete non-linear Schrödinger equations<sup>†</sup>

HAIPING SHI<sup>1</sup> and YUANBIAO ZHANG<sup>2</sup>

<sup>1</sup>Modern Business and Management Department, Guangdong Construction Polytechnic, Guangzhou 510440, China email: shp7971@163.com

<sup>2</sup>Packaging Engineering Institute, Jinan University, Zhuhai 519070, China email: abiaoa@163.com

(Received 1 July 2015; revised 17 January 2016; accepted 19 January 2016; first published online 15 February 2016)

The discrete non-linear Schrödinger equation is one of the most important inherently discrete models, having a crucial role in the modelling of a great variety of phenomena, ranging from solid-state and condensed-matter physics to biology. In this paper, a class of discrete non-linear Schrödinger equations are considered. Using critical point theory in combination with periodic approximations, we establish some new sufficient conditions on the existence results for solitons of the equation. The classical Ambrosetti–Rabinowitz superlinear condition is improved.

Key words: existence, solitons, discrete non-linear Schrödinger equations, critical point theory

# 1 Introduction

Below N, Z and R denote the sets of all natural numbers, integers and real numbers respectively.  $l^2$  denotes the space of all real functions whose second powers are summable on Z. Also, \* denotes the transpose of a vector.

In this paper, we consider the following discrete non-linear Schrödinger (DNLS) equation:

$$i\dot{\psi}_n = -\Delta\psi_n + \varepsilon_n\psi_n - f_n(\psi_n), \ n \in \mathbb{Z},$$
(1.1)

where  $\Delta \psi_n = \psi_{n+1} + \psi_{n-1} - 2\psi_n$  is discrete Laplacian operator,  $\varepsilon_n$  is real valued for each  $n \in \mathbb{Z}$ ,  $\varepsilon_{n+T} = \varepsilon_n$ ,  $f_n \in C(\mathbb{R}, \mathbb{R})$ ,  $f_{n+T}(\cdot) = f_n(\cdot)$ . Here, T is a positive integer. We assume that  $f_n(0) = 0$  and the non-linearity  $f_n(u)$  is gauge invariant, that is,

$$f_n\left(e^{i\theta}u\right) = e^{i\theta}f_n(u), \ \theta \in \mathbf{R}.$$
(1.2)

Since solitons are spatially localized time-periodic solutions and decay to zero at infinity. Thus,  $\psi_n$  has the form

$$\psi_n = u_n e^{-i\omega t},$$

<sup>†</sup> This project is supported by the National Natural Science Foundation of China (No. 11401121) and Natural Science Foundation of Guangdong Province (No. S2013010014460).

and

$$\lim_{|n|\to\infty}\psi_n=0,$$

where  $\psi_n$  is real valued for each  $n \in \mathbb{Z}$  and  $\omega \in \mathbb{R}$  is the temporal frequency. Then, (1.1) becomes

$$-\Delta u_n + \varepsilon_n u_n - \omega u_n = f_n(u_n), \ n \in \mathbb{Z},$$
(1.3)

and

$$\lim_{|n| \to \infty} u_n = 0 \tag{1.4}$$

holds.

The DNLS equation is a non-linear lattice system that appears in many areas of physics such as non-linear optics [7], biomolecular chains [17] and Bose–Einstein condensates [18]. Fundamental states supported by the DNLS equations are discrete solitons. For example, experimental observations of two-dimensional discrete solitons have been reported in [8].

In the past decade, the existence of solitons of the DNLS equations has drawn a great deal of interest [15, 16, 21–24, 31–35]. The existence for the periodic DNLS equations with superlinear non-linearity [21–24] and with saturable non-linearity [34, 35] has been studied. And the existence results of solitons of the DNLS equations without periodicity assumptions were established in [15, 16, 31, 32]. As for the existence of the homoclinic orbits of non-linear Schrödinger equations, we refer to [5, 26–29].

Actually, we consider a more general equation:

$$Lu_n - \omega u_n = f_n(u_n), \ n \in \mathbb{Z},$$
(1.5)

with the same boundary condition (1.4). Here, L is the Jacobi operator (see [30]) given by

$$Lu_n = a_n u_{n+1} + a_{n-1} u_{n-1} + b_n u_n,$$

where  $a_n$  and  $b_n$  are real valued for each  $n \in \mathbb{Z}$ ,  $a_{n+T} = a_n$ ,  $b_{n+T} = b_n$ . When  $a_n \equiv -1$  and  $b_n \equiv 2 + \varepsilon_n$ , we obtain (1.3). As usual, we say that a solution  $u = \{u_n\}$  of (1.5) is homoclinic (to 0) if (1.4) holds. Naturally, if we look for solitons of (1.1), we just need to get the homoclinic solutions of (1.5).

Let  $F_n(u) = \int_0^u f_n(t) dt, t \in \mathbf{R}$  and

$$\underline{\lambda} = \min_{n \in \mathbf{Z}} \left( b_n - |a_{n-1}| - |a_n| \right) > \omega, \ \overline{\lambda} = \max_{n \in \mathbf{Z}} \left( b_n + |a_{n-1}| + |a_n| \right).$$

Our main results are the following theorems.

**Theorem 1.1** Suppose that the following hypotheses are satisfied: (L)  $b_n - |a_{n-1}| - |a_n| > 0$ , for all  $n \in \mathbb{Z}$ ; (F<sub>1</sub>) there exist positive constants  $\varrho$  and  $a < \frac{\lambda - \omega}{2}$  such that  $|F_n(u)| \leq au^2$  for all  $n \in \mathbb{Z}$  and  $|u| \leq \varrho$ ; (F<sub>2</sub>) there exist constants  $\rho, c > \frac{\lambda - \omega}{2}$  and b such that  $F_n(u) \geq cu^2 + b$  for all  $n \in \mathbb{Z}$  and  $|u| \geq \rho$ ; (F<sub>3</sub>)  $f_n(u)u - 2F_n(u) > 0$ , for all  $n \in \mathbb{Z}$  and  $u \in \mathbb{R} \setminus \{0\}$ ; (F<sub>4</sub>)  $f_n(u)u - 2F_n(u) \rightarrow +\infty$  as  $|u| \rightarrow +\infty$ . Then, (1.5) has a non-trivial homoclinic solution.

**Remark 1.1** Since  $f_{n+T}(u) = f_n(u)$ , from  $(F_1)$ , it is easy to see that  $(f) |f_n(u)| \leq 2a|u|$ , for all  $n \in \mathbb{Z}$  and  $|u| \leq \varrho$ .

**Remark 1.2** By  $(F_2)$ , it is easy to see that there exists a constant  $\zeta > 0$  such that  $(F'_2) F_n(u) \ge cu^2 + b - \zeta, \ \forall (n, u) \in \mathbb{Z} \times \mathbb{R}.$ 

As a matter of fact, let  $\zeta = \max \{ |F_n(u) - cu^2 - b| : n \in \mathbb{Z}, |u| \leq \rho \}$ , we can easily get the desired result.

**Remark 1.3** A crucial role that the classical Ambrosetti–Rabinowitz condition plays is to ensure the boundedness of Palais–Smale sequences. This is very crucial in applying the critical-point theory. In many studies (see e.g. [23,24]), the following classical Ambrosetti– Rabinowitz condition is assumed.

(AR) there exists a constant  $\beta > 2$  such that

$$0 < \beta F_n(u) \leq u f_n(u)$$
 for all  $n \in \mathbb{Z}$  and  $u \in \mathbb{R} \setminus \{0\}$ .

It is easily checked that (AR) satisfies  $(F_2) - (F_4)$ . Thus,  $(F_2) - (F_4)$  improve (AR).

Example 1.1 Let

$$f_n(u)=\frac{\gamma u^3}{1+u^2},$$

and

$$F_n(u) = \frac{1}{2}\gamma \left[ u^2 - \ln \left( 1 + u^2 \right) \right],$$

where  $\gamma > \overline{\lambda}$ . If (L) is satisfied, then it is easy to verify all the assumptions of Theorem 1.1 are satisfied. Consequently, a non-trivial homoclinic solution is obtained.

**Theorem 1.2** Suppose that (L),  $(F_1) - (F_4)$  and the following hypothesis are satisfied: ( $F_5$ )  $a_{-n} = a_n$ ,  $b_{-n} = b_n$ ,  $f_{-n}(\cdot) = f_n(\cdot)$ . Then, (1.5) has a non-trivial even homoclinic solution.

The main idea in this paper is an application of Mountain Pass Lemma combined with an approximation technique. This idea has been employed in [20]. We mention that critical-point theory is a powerful tool to deal with the homoclinic solutions of differential equations [9–14] and is used to study homoclinic solutions of discrete systems in recent years [1–4, 6, 19, 20, 33, 34]. We should emphasize that the results are obtained without the classical Ambrosetti–Rabinowitz condition [23, 24].

#### 2 Preliminaries

In this section, we shall present some definitions and lemmas that will be used in the proof of our results.

Let S be the set of sequences  $u = (..., u_{-n}, ..., u_{-1}, u_0, u_1, ..., u_n, ...) = \{u_n\}_{n=-\infty}^{+\infty}$ , that is

$$S = \{\{u_n\} | u_n \in \mathbf{R}, n \in \mathbf{Z}\}.$$

For any  $u, v \in S$ ,  $a, b \in \mathbf{R}$ , au + bv is defined by

$$au + bv = \{au_n + bv_n\}_{n = -\infty}^{+\infty}$$

Then, S is a vector space.

For any fixed positive integers m and T,  $E_m$  is defined as a subspace of S by

$$E_m = \{ u \in S | u_{n+2mT} = u_n, n \in \mathbf{Z} \}.$$

Clearly,  $E_m$  is isomorphic to  $\mathbf{R}^{2mT}$ .  $E_m$  can be equipped with the inner product

$$\langle u, v \rangle = \sum_{j=-mT}^{mT-1} u_j v_j, \ u, v \in E_m,$$
(2.1)

by which the norm  $\|\cdot\|$  can be induced by

$$||u|| = \left(\sum_{j=-mT}^{mT-1} u_j^2\right)^{\frac{1}{2}}, u \in E_m,$$
(2.2)

respectively. We also define a norm  $\|\cdot\|_{\infty}$  in  $E_m$  by

$$\|u\|_{\infty}=\max_{j\in\mathbf{Z}}|u_j|\,,\ u\in E_m.$$

Consider the functional J on  $E_m$  defined by

$$J(u) = \frac{1}{2} \sum_{n=-mT}^{mT-1} Lu_n \cdot u_n - \frac{\omega}{2} \sum_{n=-mT}^{mT-1} u_n^2 - \sum_{n=-mT}^{mT-1} F_n(u_n).$$
(2.3)

Then,

$$\langle J'(u), v \rangle = \sum_{n=-mT}^{mT-1} \left[ Lu_n \cdot v_n - \omega u_n v_n - f_n(u_n) v_n \right], \ u, v \in E_m.$$
(2.4)

Since  $\{a_n\}$  and  $\{b_n\}$  are *T*-periodic, it is easy to see that the critical points of *J* in  $E_m$  are exactly 2mT-periodic solutions of equation (1.5).

Let *E* be a real Banach space,  $J \in C^1(E, \mathbf{R})$ , i.e., *J* is a continuously Fréchet-differentiable functional defined on *E*. *J* is said to satisfy the Palais–Smale condition (P.S. condition for short) if any sequence  $\{u_n\} \subset E$  for which  $\{J(u_n)\}$  is bounded and  $J'(u_n) \to 0$   $(n \to \infty)$ possesses a convergent subsequence in *E*. Let  $B_{\rho}$  denote the open ball in *E* about 0 of radius  $\rho$  and let  $\partial B_{\rho}$  denote its boundary.

**Lemma 2.1** (Mountain Pass Lemma [25]). Let E be a real Banach space and  $J \in C^1(E, \mathbb{R})$  satisfy the P.S. condition. If J(0) = 0 and

 $(J_1)$  there exist constants  $\rho$ ,  $\alpha > 0$  such that  $J|_{\partial B_{\rho}} \ge \alpha$ , and  $(J_1)$  there exists  $\alpha \in E \setminus B$ , such that  $J(\alpha) \le 0$ .

 $(J_2)$  there exists  $e \in E \setminus B_\rho$  such that  $J(e) \leq 0$ . Then, J possesses a critical value  $c \ge \alpha$  given by

$$c = \inf_{g \in \Gamma} \max_{s \in [0,1]} J(g(s)),$$
(2.5)

where

$$\Gamma = \{ g \in C([0,1], E) | g(0) = 0, \ g(1) = e \}.$$
(2.6)

Lemma 2.2 The following inequality is true:

$$\underline{\lambda} \|u\|^2 \leqslant \sum_{n=-mT}^{mT-1} Lu_n \cdot u_n \leqslant \overline{\lambda} \|u\|^2.$$
(2.7)

Proof Let

$$\sum_{m=-mT}^{mT-1} Lu_n \cdot u_n = \langle P_m u, u \rangle, \qquad (2.8)$$

where  $u = (u_{-mT}, \dots, u_{-1}, u_0, u_1, \dots, u_{mT-1})^*$ ,

$$P_m = \begin{pmatrix} b_{-mT} & a_{-mT} & 0 & \cdots & 0 & a_{-mT-1} \\ a_{-mT} & b_{-mT+1} & a_{-mT+1} & \cdots & 0 & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & 0 & 0 & \cdots & b_{mT-2} & a_{mT-2} \\ a_{mT-1} & 0 & 0 & \cdots & a_{mT-2} & b_{mT-1} \end{pmatrix}_{2mT \times 2mT}.$$

By (*L*),  $P_m$  is positive definite. Suppose that the eigenvalues of  $P_m$  are  $\lambda_{-mT}$ ,  $\lambda_{-mT+1}$ , ...,  $\lambda_{-1}$ ,  $\lambda_0$ ,  $\lambda_1$ , ...,  $\lambda_{mT-2}$ ,  $\lambda_{mT-1}$ , then they are all greater than zero. So, by (2.2) and (2.8), we get

$$\underline{\lambda}\|u\|^2 \leqslant \sum_{n=-mT}^{mT-1} Lu_n \cdot u_n \leqslant \overline{\lambda}\|u\|^2.$$

The proof of Lemma 2.2 is complete.

# **3** Proofs of theorems

In this section, we shall prove our main results by using the critical-point method.

**Lemma 3.1** Suppose that (L) and  $(F_1) - (F_4)$  are satisfied. Then, J satisfies the P.S. condition.

**Proof** Assume that  $\{u^{(i)}\}_{i\in\mathbb{N}}$  in  $E_m$  is a sequence such that  $\{J(u^{(i)})\}_{i\in\mathbb{N}}$  is bounded. Then, there is a positive constant K such that  $-K \leq J(u^{(i)})$ . By (2.7) and  $(F'_2)$ , we have

$$\begin{split} -K &\leq J\left(u^{(i)}\right) \leq \frac{\bar{\lambda} - \omega}{2} \left\|u^{(i)}\right\|^2 - \sum_{n=-mT}^{mT-1} \left[c\left(u^{(i)}_n\right)^2 + b - \zeta\right] \\ &\leq \left(\frac{\bar{\lambda} - \omega}{2} - c\right) \left\|u^{(i)}\right\|^2 + 2mT\left(\zeta - b\right). \end{split}$$

Therefore,

$$\left(c - \frac{\bar{\lambda} - \omega}{2}\right) \left\| u^{(i)} \right\|^2 \leq 2mT \left(\zeta - b\right) + K.$$
(3.1)

Since  $c > \frac{\bar{\lambda} - \omega}{2}$ , (3.1) implies that  $\{u^{(i)}\}_{i \in \mathbb{N}}$  is bounded in  $E_m$ . Thus,  $\{u^{(i)}\}_{i \in \mathbb{N}}$  possesses a convergence subsequence in  $E_m$ . The desired result follows.

**Lemma 3.2** Suppose that (L) and  $(F_1)-(F_4)$  are satisfied. Then, for which u, (1.5) possesses a 2mT-periodic solution  $u^{(m)} \in E_m$ .

**Proof** In our case, it is clear that J(0) = 0. By Lemma 3.1, J satisfies the P.S. condition. By  $(F_1)$ , we have

$$J(u) \ge \frac{\lambda}{2} \sum_{n=-mT}^{mT-1} u_n^2 - \frac{\omega}{2} \sum_{n=-mT}^{mT-1} u_n^2 - a \sum_{n=-mT}^{mT-1} u_n^2$$
$$= \frac{\lambda - \omega - 2a}{2} \|u\|^2.$$

Taking  $\alpha = \frac{\lambda - \omega - 2a}{2} \varrho^2 > 0$ , we obtain

$$J(u)|_{\partial B_{\rho}} \geqslant \alpha > 0,$$

which implies that J satisfies the condition  $(J_1)$  of the Mountain Pass Lemma.

Next, we shall verify the condition  $(J_2)$ .

There exists a sufficiently large number  $\varepsilon > \max\{\varrho, \rho\}$  such that

$$\left(c - \frac{\bar{\lambda} - \omega}{2}\right)\varepsilon^2 \ge |b|. \tag{3.2}$$

Let  $e \in E_m$  and

$$e_n = \begin{cases} \varepsilon, \text{ if } n = 0, \\ 0, \text{ if } n \in \{ j \in \mathbf{Z} : -mT \leq j \leq mT - 1 \text{ and } j \neq 0 \}. \end{cases}$$

Then,

$$F_n(e_n) = \begin{cases} F_n(\varepsilon), \text{ if } n = 0, \\ 0, \quad \text{ if } n \in \{j \in \mathbf{Z} : -mT \le j \le mT - 1 \text{ and } j \neq 0\}. \end{cases}$$

With (3.2), we have

$$J(e) = \frac{1}{2} \sum_{n=-mT}^{mT-1} Le_n \cdot e_n - \frac{\omega}{2} \sum_{n=-mT}^{mT-1} e_n^2 - \sum_{n=-mT}^{mT-1} F_n(e_n)$$
  
$$\leq \frac{\bar{\lambda}}{2} \|e\|^2 - \frac{\omega}{2} \|e\|^2 - c \|e\|^2 - b$$
  
$$= -\left(c - \frac{\bar{\lambda} - \omega}{2}\right) \varepsilon^2 - b \leq 0.$$

All the assumptions of the Mountain Pass Lemma have been verified. Consequently by Lemmas 2.1 and 3.1, J possesses a critical value  $c_m$  given by (2.5) and (2.6) with  $E = E_m$  and  $\Gamma = \Gamma_m$ , where  $\Gamma_m = \{g_m \in C([0, 1], E_m) | g_m(0) = 0, g_m(1) = e, e \in E_m \setminus B_{\varepsilon}\}$ . Let  $u^{(m)}$  denote the corresponding critical point of J in  $E_m$ . Note that  $||u^{(m)}|| \neq 0$  since  $c_m > 0$ .  $\Box$ 

**Lemma 3.3** Suppose that (L) and  $(F_1) - (F_4)$  are satisfied. Then, there exist positive constants  $\rho$  and  $\eta$  independent of m such that

$$\varrho \leqslant \left\| u^{(m)} \right\|_{\infty} \leqslant \eta. \tag{3.3}$$

**Proof** The continuity of  $F_n(u)$  with respect to the variable u implies that there exists a constant  $\tau > 0$  such that  $|F_n(u)| \leq \tau$  for  $|u| \leq \varrho$ . It is clear that

$$J\left(u^{(m)}\right) \leq \max_{0 \leq s \leq 1} \left\{ \frac{1}{2} \sum_{n=-mT}^{mT-1} \left| L(se)_n \cdot (se)_n - \omega(se)_n^2 \right| - \sum_{n=-mT}^{mT-1} F_n\left((se)_n\right) \right\}$$
$$\leq \frac{\bar{\lambda} + \omega}{2} \|e\|^2 + \tau$$
$$= \frac{(\bar{\lambda} + \omega)\varepsilon^2}{2} + \tau.$$

Let  $\xi = \frac{(\bar{\lambda}+\omega)\varepsilon^2}{2} + \tau$ , we have that  $J(u^{(m)}) \leq \xi$ , which is independent of *m*. From (2.5) and (2.6), we have

$$J(u^{(m)}) = \frac{1}{2} \sum_{n=-mT}^{mT-1} f_n(u_n^{(m)}) u_n^{(m)} - \sum_{n=-mT}^{mT-1} F_n(u_n^{(m)}) \leqslant \xi.$$

By  $(F_3)$  and  $(F_4)$ , there exists a constant  $\eta > 0$  such that

$$\frac{1}{2}f_n(v)v - F_n(v) > \xi, \text{ for all } |v| \ge \eta,$$

which implies that  $|u_n^{(m)}| \leq \eta$  for all  $n \in \mathbb{Z}$ , that is,  $||u^{(m)}||_{\infty} \leq \eta$ .

From the definition of J, we have

$$0 = \langle J'(u^{(m)}), u^{(m)} \rangle \ge (\underline{\lambda} - \omega) \sum_{n=-mT}^{mT-1} |u_n^{(m)}|^2 - \sum_{n=-mT}^{mT-1} f_n(u_n^{(m)}) u_n^{(m)}.$$

Therefore, combined with  $(F_2)$ , we get

$$(\underline{\lambda} - \omega) \| u^{(m)} \|^{2} \leq \sum_{n=-mT}^{mT-1} f_{n} (u_{n}^{(m)}) u_{n}^{(m)} \leq \left\{ \sum_{n=-mT}^{mT-1} \left[ f_{n} (u_{n}^{(m)}) \right]^{2} \right\}^{\frac{1}{2}} \| u^{(m)} \|.$$

That is,

$$(\underline{\lambda} - \omega) \left\| u^{(m)} \right\| \leq \left\{ \sum_{n=-mT}^{mT-1} \left[ f_n \left( u_n^{(m)} \right) \right]^2 \right\}^{\frac{1}{2}}$$

Thus,

$$(\underline{\lambda} - \omega)^2 \left\| u^{(m)} \right\|^2 \leqslant \sum_{n=-mT}^{mT-1} \left[ f_n \left( u_n^{(m)} \right) \right]^2.$$
(3.4)

Combined with  $(F_1)$ , we get

$$(\underline{\lambda} - \omega)^2 \|u^{(m)}\|^2 \leq \sum_{n=-mT}^{mT-1} [2a |u_n^{(m)}|]^2 = 4a^2 \|u^{(m)}\|^2.$$

Thus, we have  $u^{(m)} = 0$ . But this contradicts  $||u^{(m)}|| \neq 0$ , which shows that

$$\left\| u^{(m)} \right\|_{\infty} \ge \varrho,$$

and the proof of Lemma 3.3 is finished.

**Proof of Theorem 1.1.** Consider the sequence  $\{u_n^{(m)}\}_{n \in \mathbb{Z}}$  of 2mT-periodic solutions found in Lemma 3.2. First, by (3.3), for any  $m \in \mathbb{N}$ , there exists a constant  $n_m \in \mathbb{Z}$  independent of m such that

$$\left|u_{n_m}^{(m)}\right| \geqslant \varrho. \tag{3.5}$$

Since  $a_n$ ,  $b_n$  are *T*-periodic in n,  $\{u_{n+jT}^{(m)}\}$  ( $\forall j \in \mathbb{N}$ ) is also 2mT-periodic solution of (1.3). Hence, making such shifts, we can assume that  $0 \leq n_m \leq T - 1$  in (3.5). Moreover, passing to a subsequence of ms, we can even assume that  $n_m = n_0$  is independent of m.

Next, we extract a subsequence, still denote by  $u^{(m)}$ , such that

$$u_n^{(m)} \to u_n, \ m \to \infty, \ \forall n \in \mathbb{Z}$$

Inequality (3.5) implies that  $|u_{n_0}| \ge \rho$  and, hence,  $u = \{u_n\}$  is a non-zero sequence. Moreover,

$$Lu_n - \omega u_n - f_n(u_n)$$
  
= 
$$\lim_{m \to \infty} \left[ Lu_n^{(m)} - \omega u_n^{(m)} - f_n(u_n^{(m)}) \right] = 0$$

So  $u = \{u_n\}$  is a solution of (1.5).

Finally, we show that  $u \in l^2$ . For  $u_m \in E_m$ , let

$$P_m = \left\{ n \in \mathbf{Z} : \left| u_n^{(m)} \right| < \varrho, -mT \leqslant n \leqslant mT - 1 \right\},\$$

 $\square$ 

 $Q_m = \left\{ n \in \mathbf{Z} : \left| u_n^{(m)} \right| \ge \varrho, -mT \le n \le mT - 1 \right\}.$ Since  $f_n(u) \in C(\mathbf{R}, \mathbf{R})$ , there exist constants  $\overline{\xi} > 0, \, \underline{\xi} > 0$  such that

 $\max\left\{|f_n(u)|: \varrho \leqslant |u| \leqslant \eta, \ n \in \mathbf{Z}\right\} \leqslant \overline{\xi},$ 

$$\min\left\{\frac{1}{2}f_n(u)u - F_n(u) : \varrho \leqslant |u| \leqslant \eta, \ n \in \mathbf{Z}\right\} \ge \underline{\xi}$$

For  $n \in Q_m$ ,

$$\left|f_{n}\left(u_{n}^{(m)}\right)\right| \leqslant \frac{\overline{\xi}}{\underline{\xi}} \left[\frac{1}{2}f_{n}\left(u_{n}^{(m)}\right)u_{n}^{(m)} - F_{n}\left(u_{n}^{(m)}\right)\right].$$
(3.6)

By  $(F_1)$ , (3.4) and (3.6), we have

$$\begin{aligned} (\underline{\lambda} - \omega)^2 \| u^{(m)} \|^2 &\leq \sum_{n \in P_m} \left[ f_n \left( u_n^{(m)} \right) \right]^2 + \sum_{n \in Q_m} \left[ f_n \left( u_n^{(m)} \right) \right]^2 \\ &\leq \sum_{n \in P_m} \left[ 2a \left| u_{n+1}^{(m)} \right| \right]^2 + \sum_{n \in Q_m} \left[ \frac{1}{2} f_n \left( u_n^{(m)} \right) u_n^{(m)} - F_n \left( u_n^{(m)} \right) \right] \\ &\leq 4a^2 \| u^{(m)} \|^2 + \frac{\overline{\zeta} \underline{\zeta}}{\underline{\zeta}}. \end{aligned}$$

Thus,

$$\left\|u^{(m)}\right\|^2 \leq \frac{\overline{\xi}\xi}{\underline{\xi}\left[(\underline{\lambda}-\omega)^2-4a^2\right]}.$$

For any fixed  $D \in \mathbb{Z}$  and *m* large enough, we have that

$$\sum_{n=-D}^{D} \left| u_n^{(m)} \right|^2 \leq \left\| u^{(m)} \right\|^2 \leq \frac{\overline{\xi}\xi}{\underline{\xi}\left[ (\underline{\lambda} - \omega)^2 - 4a^2 \right]}$$

Since  $\overline{\xi}$ ,  $\underline{\xi}$ ,  $\underline{\xi}$ ,  $\underline{\lambda}$ , a and  $\omega$  are constants independent of *m*, passing to the limit, we have that

$$\sum_{n=-D}^{D} |u_n|^2 \leq \frac{\overline{\xi}\xi}{\underline{\xi}\left[(\underline{\lambda}-\omega)^2 - 4a^2\right]}.$$

Due to the arbitrariness of  $D, u \in l^2$ . Therefore, u satisfies  $u_n \to 0$  as  $|n| \to \infty$ .

Proof of Theorem 1.2. Consider the following boundary problem:

$$\begin{cases} Lu_n - \omega u_n - f_n(u_n) = 0, & -mT \leq \mathbb{Z} \leq mT, \\ a_{-mT} = a_{mT} = 0, & b_{-mT} = b_{mT} = 0, \\ a_{-n} = a_n, & b_{-n} = b_n, & -mT \leq \mathbb{Z} \leq mT. \end{cases}$$

Let S be the set of sequences  $u = (..., u_{-n}, ..., u_{-1}, u_0, u_1, ..., u_n, ...) = \{u_n\}_{n=-\infty}^{+\infty}$ , that is

$$S = \{\{u_n\} | u_n \in \mathbf{R}, n \in \mathbf{Z}\}.$$

For any  $u, v \in S$ ,  $a, b \in \mathbf{R}$ , au + bv is defined by

$$au + bv = \{au_n + bv_n\}_{n = -\infty}^{+\infty}.$$

Then, S is a vector space.

For any given positive integers m and T,  $\tilde{E}_m$  is defined as a subspace of S by

$$\tilde{E}_m = \{ u \in S | u_{-n} = u_n, \ \forall n \in \mathbf{Z} \}.$$

Clearly,  $\tilde{E}_m$  is isomorphic to  $\mathbf{R}^{2mT+1}$ .  $\tilde{E}_m$  can be equipped with the inner product

$$\langle u, v \rangle = \sum_{j=-mT}^{mT} u_j v_j, \ \forall u, v \in \tilde{E}_m$$

by which the norm  $\|\cdot\|$  can be induced by

$$\|u\| = \left(\sum_{j=-mT}^{mT} u_j^2\right)^{\frac{1}{2}}, \ \forall u \in \tilde{E}_m.$$

It is obvious that  $\tilde{E}_m$  is Hilbert space with 2mT + 1-periodicity and linearly homeomorphic to  $\mathbf{R}^{2mT+1}$ .

Similarly to the proof of Theorem 1.1, we can also prove Theorem 1.2. For simplicity, we omit its proof.  $\Box$ 

# References

- CHEN, P. & TANG, X. H. (2011) Existence of infinitely many homoclinic orbits for fourth-order difference systems containing both advance and retardation. *Appl. Math. Comput.* 217(9), 4408–4415.
- [2] CHEN, P. & TANG, X. H. (2011) Existence and multiplicity of homoclinic orbits for 2nth-order nonlinear difference equations containing both many advances and retardations. J. Math. Anal. Appl. 381(2), 485–505.
- [3] CHEN, P. & TANG, X. H. (2013) Infinitely many homoclinic solutions for the second-order discrete p-Laplacian systems. Bull. Belg. Math. Soc. 20(2), 193-212.
- [4] CHEN, P. & TANG, X. H. (2013) Existence of homoclinic solutions for some second-order discrete Hamiltonian systems. J. Differ. Equ. Appl. 19(4), 633–648.
- [5] CHEN, P. & TIAN, C. (2014) Infinitely many solutions for Schrödinger-Maxwell equations with indefinite sign subquadratic potentials. *Appl. Math. Comput.* 226(1), 492–502.
- [6] CHEN, P. & WANG, Z. M. (2012) Infinitely many homoclinic solutions for a class of nonlinear difference equations. *Electron. J. Qual. Theory Differ. Equ.* 47(2), 1–18.
- [7] CHRISTODOULIDES, D. N., LEDERER, F. & SILBERBERG, Y. (2003) Discretizing light behaviour in linear and nonlinear waveguide lattices. *Nature* 424(3), 817–823.
- [8] FLEISCHER, J. W., SEGEV, M., EFREMIDIS, N. K. & CHRISTODOULIDES, D. N. (2003) Observation of two-dimensional discrete solitons in optically induced nonlinear photonic lattices. *Nature* 422(3), 147–150.

- [9] GUO, C. J., AGARWAL, R. P., WANG, C. J. & O'REGAN, D. (2014) The existence of homoclinic orbits for a class of first order superquadratic Hamiltonian systems. *Mem. Differ. Equ. Math. Phys.* 61(2), 83–102.
- [10] GUO, C. J., O'REGAN, D. & AGARWAL, R. P. (2010) Existence of homoclinic solutions for a class of the second-order neutral differential equations with multiple deviating arguments. *Adv. Dyn. Syst. Appl.* 5(1), 75–85.
- [11] GUO, C. J., O'REGAN, D., XU, Y. T. & AGARWAL, R. P. (2011) Existence of subharmonic solutions and homoclinic orbits for a class of high-order differential equations. *Appl. Anal.* 90(7), 1169–1183.
- [12] GUO, C. J., O'REGAN, D., XU, Y. T. & AGARWAL, R. P. (2010) Homoclinic orbits for a singular second-order neutral differential equation. J. Math. Anal. Appl. 366(2), 550–560.
- [13] GUO, C. J., O'REGAN, D., XU, Y. T. & AGARWAL, R. P. (2012) Existence and multiplicity of homoclinic orbits of a second-order differential difference equation via variational methods. *Appl. Math. Inform. Mech.* 4(1), 1–15.
- [14] GUO, C. J., O'REGAN, D., XU, Y. T. & AGARWAL, R. P. (2013) Existence of homoclinic orbits of a class of second order differential difference equations. *Dyn. Contin. Discrete Impuls. Syst.* Ser. B Appl. Algorithms 20(6), 675–690.
- [15] HUANG, M. H. & ZHOU, Z. (2013) Standing wave solutions for the discrete coupled nonlinear Schrödinger equations with unbounded potentials. *Abstr. Appl. Anal.* 2013(1), 1–6.
- [16] HUANG, M. H. & ZHOU, Z. (2013) On the existence of ground state solutions of the periodic discrete coupled nonlinear Schrödinger lattice. J. Appl. Math. 2013(2), 1–8.
- [17] KOPIDAKIS, G., AUBRY, S. & TSIRONIS, G. P. (2001) Targeted energy transfer through discrete breathers in nonlinear systems. *Phys. Rev. Lett.* 87(16), 165501.
- [18] LIVI, R., FRANZOSI, R. & OPPO, G. L. (2006) Self-localization of Bose–Einstein condensates in optical lattices via boundary dissipation. *Phys. Rev. Lett.* 97(6), 060401.
- [19] MA, M. J. & GUO, Z. M. (2006) Homoclinic orbits for second order self-adjoint difference equations. J. Math. Anal. Appl. 323(1), 513–521.
- [20] MA, M. J. & GUO, Z. M. (2007) Homoclinic orbits and subharmonics for nonlinear second order difference equations. *Nonlinear Anal.* 67(6), 1737–1745.
- [21] MAI, A. & ZHOU, Z. (2013) Discrete solitons for periodic discrete nonlinear Schrödinger equations. Appl. Math. Comput. 222(1), 34–41.
- [22] MAI, A. & ZHOU, Z. (2013) Ground state solutions for the periodic discrete nonlinear Schrödinger equations with superlinear nonlinearities. *Abstr. Appl. Anal.* 2013(3), 1–11.
- [23] PANKOV, A. (2006) Gap solitons in periodic discrete nonlinear Schrödinger equations. Nonlinearity 19(1), 27–41.
- [24] PANKOV, A. (2007) Gap solitons in periodic discrete nonlinear Schrödinger equations II: A generalized Nehari manifold approach. Discrete Contin. Dyn. Syst. 19(2), 419–430.
- [25] RABINOWITZ, P. H. (1986) Minimax Methods in Critical Point Theory with Applications to Differential Equations, Amer. Math. Soc., Providence, RI, New York, pp. 366–369.
- [26] TANG, X. H. (2015) Non-Nehari manifold method for asymptotically periodic Schrödinger equations. Sci. China Math. 58(4), 715–728.
- [27] TANG, X. H. (2014) New conditions on nonlinearity for a periodic Schrödinger equation having zero as spectrum. J. Math. Anal. Appl. 413(1), 392–410.
- [28] TANG, X. H. (2014) New super-quadratic conditions on ground state solutions for superlinear Schrödinger equation. Adv. Nonlinear Stud. 14(2), 361–374.
- [29] TANG, X. H. (2013) Infinitely many solutions for semilinear Schrödinger equations with signchanging potential and nonlinearity. J. Math. Anal. Appl. 401(1), 407–415.
- [30] TESCHL, G. (2000) Jacobi Operators and Completely Integrable Nonlinear Lattices, Amer. Math. Soc., Providence, RI, New York, pp. 232–239.
- [31] ZHANG, G. P. & LIU, F. S. (2009) Existence of breather solutions of the DNLS equations with unbounded potentials. *Nonlinear Anal.* 71(12), 786–792.
- [32] ZHOU, Z. & MA, D. F. (2015) Multiplicity results of breathers for the discrete nonlinear Schrödinger equations with unbounded potentials. *Sci. China Math.* 58(4), 781–790.

- [33] ZHOU, Z. & YU, J. S. (2013) Homoclinic solutions in periodic nonlinear difference equations with superlinear nonlinearity. *Acta Math. Sin. Engl. Ser.* **29**(9), 1809–1822.
- [34] ZHOU, Z. & YU, J. S. (2010) On the existence of homoclinic solutions of a class of discrete nonlinear periodic systems. J. Differ. Equ. 249(5), 1199–1212.
- [35] ZHOU, Z., YU., J. S. & CHEN, Y. M. (2010) On the existence of gap solitons in a periodic discrete nonlinear Schrödinger equation with saturable nonlinearity. *Nonlinearity* 23(7), 1727–1740.