

## STACKING FAULTS WITH 180° LAYER ROTATION IN CELADONITE, AN Fe- AND Mg-RICH DIOCTAHEDRAL MICA

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**Abstract**—Stacking disorder in celadonite, a dioctahedral mica with Fe and Mg as major octahedral cations and which generally adopts the 1M stacking sequence, was investigated mainly by using transmission electron microscopy (TEM). The selected-area electron diffraction patterns with *0kl* reflections along the [100] beam direction correspond to the 1M stacking but those along the [110], [1 $\bar{1}$ 0], [010], [310], and [3 $\bar{1}$ 0] directions are frequently streaked along the [001]\* direction or contain extra spots from twinned domains. Three-dimensional stacking analyses using sets of two high-resolution TEM images along different directions of the same area of one crystal indicate that all stacking faults involve only 180° layer rotations. These stacking faults produce greater peaks of *0kl* reflections than expected in powder X-ray diffraction (XRD) patterns. Simulation of the XRD patterns indicated that 180° layer rotations occur at >10% of total interlayer regions in one celadonite specimen. The interlayer region of celadonite is characterized by a near-zero ditrigonal rotation angle, a small surface corrugation of the basal oxygen plane, and a small amount of Al substitution in the tetrahedral sheets. These features suggest that there is no preference for any of the six stacking angles around the interlayer region. The abundance of 180° layer rotation rather than  $\pm 60^\circ$  and  $\pm 120^\circ$  in the present specimens may be related to their ribbon-like morphologies elongated along the *a* axis.

**Key Words**—Celadonite, Electron Microscopy, Mica, Polytype, Stacking Faults, X-ray Diffraction.

### INTRODUCTION

Polytypism is a common phenomenon in many phyllosilicates. Polytypism of micas, in particular, has been investigated extensively, both experimentally and theoretically, but is not yet fully understood. Polytypism of micas may be described as the rotation between adjacent layers, the angle of which is restricted to a multiple (*n*) of 60° about the *c*\* axis. Polytypes with layer rotations of  $2n60^\circ$  (subfamily A) occur more frequently than those of  $(2n+1)60^\circ$  (subfamily B). This is primarily explained by the atomic configuration of the interlayer region of micas (e.g. Ferraris and Ivaldi, 2002). However, the occurrence of subfamily B polytypes, although uncommon, is not rare. Lepidolite often occurs as the  $2M_2$  polytype with  $\pm 60^\circ$  rotations (e.g. Levinson, 1953), and several illite specimens adopt the  $2M_2$  configuration (Threadgold, 1959; Drits *et al.*, 1966; Zhukhlistov *et al.*, 1973; Slonimskaya *et al.*, 1978). The origins of the  $2M_2$  polytype in lepidolite were discussed in previous works (e.g. Takeda *et al.*, 1971; Brigatti *et al.*, 2005). Furthermore, the 2O polytype with 180° rotation occurs in anandite from Sri Lanka (Giuseppetti and Tadini, 1972) and in phlogopite (Ferraris *et al.*, 2001).

Stacking faults in micas have origins similar to the structural variables that produce polytypes, but they are not periodic and occur at a unit-layer scale. An effective technique to investigate the stacking faults or long-period polytypes is high-resolution transmission electron microscopy (HRTEM) (e.g. Iijima and Buseck, 1978; Baronnet, 1992; Kogure and Nespolo, 1999a, 1999b). Kogure and his coworkers (Kogure and Nespolo, 1999a; Kogure and Bunno, 2004; Kameda *et al.*, 2007) reported the occurrence of stacking faults with  $(2n+1)60^\circ$  rotations in various micas. Kogure and Nespolo (1999a), in particular, found that stacking faults that involve various rotation angles exist in a  $2M_1$ -dominant Mg-rich annite in granite. They employed a ‘three dimensional stacking analysis’ using a set of two HRTEM images along the  $X_i$  direction and the  $Y_i$  direction where  $X_i$  is, for example, [100], [110], or [1 $\bar{1}$ 0] and  $Y_i$  is [010], [310], or [3 $\bar{1}$ 0] in the 1M and  $2M_1$  polytypes (Bailey, 1984). Interstratification of  $2n60^\circ$  and  $(2n+1)60^\circ$  rotations was also reported in lepidolite (Kogure and Bunno, 2004) and in magnesian illite (gümbelite) (Kameda *et al.*, 2007); both sets of authors used HRTEM.

The analyses of stacking faults in layer minerals by XRD methods is difficult because stacking errors only broaden reflections or change several reflections to a ‘band.’ However, XRD studies are enhanced when computer simulation of the experimental XRD patterns is used. Comparison of calculated positions, intensities, and profiles with observed reflections provides an

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estimate of the content and distribution of stacking errors (*e.g.* Drits and Tchoubar, 1990). For example, Sakharov *et al.* (1990) found that in glauconite mica layers involving  $2n60^\circ$  and  $(2n+1)60^\circ$ , layer rotations occur with equal probability, but layers with different orientations are generally segregated. Studies by McCarty and Reynolds (1995), Cuadros and Altaner (1998a, 1998b), and Ylagan *et al.* (2000) on illite-smectite showed that stacking faults involving layer rotations of  $2n60^\circ$  and  $(2n+1)60^\circ$  were found to occur.

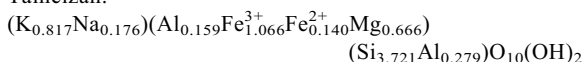
Celadonite is an  $\text{Fe}^{3+}$ -rich dioctahedral mica containing similar amounts of  $\text{Fe}^{2+}$  and/or  $\text{Mg}^{2+}$  at the octahedral sites (Rieder *et al.*, 1999), and it commonly adopts the *1M* polytype (Buckley *et al.*, 1978). Although celadonite is commonly impure and fine-grained, and therefore difficult to study, Zhukhlistov *et al.* (1977), Zhukhlistov (2005), and Tshipursky and Drits (1986) refined the structure of celadonite using oblique-texture electron diffraction.

Kogure *et al.* (2007) examined unheated and heat-treated celadonite and found that celadonite-*1M* generally contains frequent stacking faults. Surprisingly, they found that all faults involve the layer rotation of  $180^\circ$  only. The present paper describes the observation of the faults by HRTEM, the powder XRD patterns and their simulation, and possible origins for these stacking faults.

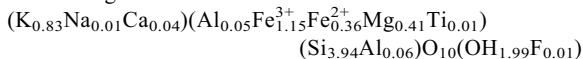
## MATERIALS AND METHODS

Three celadonite specimens were investigated: from Taiheizan, Akita, Japan (Kimbara and Shimoda, 1973); from Krivoi Rog, Ukraine (Zhukhlistov *et al.*, 1977); and from Zaval'e, Pobuj'e, Ukraine (Malkova, 1956; Tshipursky and Drits, 1986). The reported formulae are:

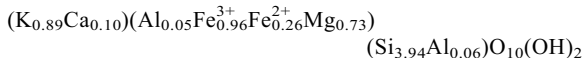
Taiheizan:



Krivoi Rog:



Zaval'e:



The Taiheizan specimen contains more Al at the octahedral and tetrahedral sites than the other two specimens.

Each specimen for TEM examination was prepared following the method proposed by Kogure (2002). In short, powder was embedded in epoxy resin between two glass slides. After hardening, the glass slides were cut using a diamond wheel to laths of  $\sim 1$  mm thick. The laths were thinned to  $\sim 50$   $\mu\text{m}$  by mechanical grinding and then argon ion milled. Examination by HRTEM was performed at 200 kV using a JEOL JEM-2010 UHR with a nominal point resolution of 2.0 Å. The HRTEM images

were captured on film or using a Gatan MSC 794 bottom-mounted CCD camera; they were taken at sufficiently thin regions of the specimen, and by adjusting the defocus value, to record the contrast that corresponds to the projected potential of the crystal structure (Kogure, 2002). Noisy contrast from amorphous materials in HRTEM images was removed using a Wiener-filter (Marks, 1996; Kilaas, 1998) developed by K. Ishizuka (HREM Research, Inc.) and implemented using a Gatan DigitalMicrograph version 3.10.0.

Powder XRD patterns were collected using a Rigaku RINT-Ultima<sup>+</sup> diffractometer with  $\text{CuK}\alpha$  radiation, a graphite monochromator, a 0.3 mm receiving slit, and  $0.5^\circ$  divergence and anti-scatter slits. A continuous scan rate of  $1^\circ 2\theta \text{ min}^{-1}$  was adopted. Simulation of powder XRD patterns was performed using *DIFFaX* (Treacy *et al.*, 1991) which calculates diffraction patterns from layered materials with various stacking sequences.

## RESULTS

### HRTEM analyses

Typical images (Figure 1a) of the Zaval'e celadonite crystals, dispersed on a microgrid with holey carbon film, revealed that the crystals have a lath- or ribbon-like morphology that is common for the three specimens investigated. Cross sections of the crystals (Figure 1b) showed them to be 50 to  $\sim 150$  nm thick. Selected-area electron diffraction (SAED) patterns from the same area in a celadonite crystal, along the [100] (Figure 1c) and  $[3\bar{1}0]$  (or [310]) (Figure 1d) directions, indicate that the stacking along [100] is *1M* and well ordered because any streaking parallel to the  $c^*$  axis is below detection. In contrast, the pattern along the  $[3\bar{1}0]$  direction (Figure 1d) contains extra spots on the 13 $l$  row, indicating the existence of twinned domains formed by mutual layer rotation with angles of  $(2n+1)60^\circ$  (Kogure and Nespolo, 1999a). This apparent inconsistency between the two patterns suggests that stacking disorder involves  $180^\circ$  layer rotations in the *1M* stacking sequence. The intralayer shift of the 2:1 layer along the  $[\bar{1}00]$  direction for the *1M* polytype changes to the [100] direction by a  $180^\circ$  layer rotation, which does not affect the diffraction pattern along the [100] direction.

The HRTEM image of the Krivoi Rog celadonite recorded along the [100] direction (Figure 2a) revealed that the stacking is uniform and no stacking faults are observed. In contrast, the HRTEM image of the same area, but along the  $[3\bar{1}0]$  (or [310]) direction (Figure 2b), shows several stacking faults. The white bar in each 2:1 layer connects the closest dark spots in the lower and upper tetrahedral sheets. These bars are slanted by  $\pm a/6$ . The symbol '+' and '-' at the bottom indicates the direction of the slant of the white bar in each layer; '+' for the bar diagonally down ( $+a/6$ ) and '-' diagonally up ( $-a/6$ ). This slant corresponds to the direction of the projection of the intralayer shift in a 2:1 layer on the

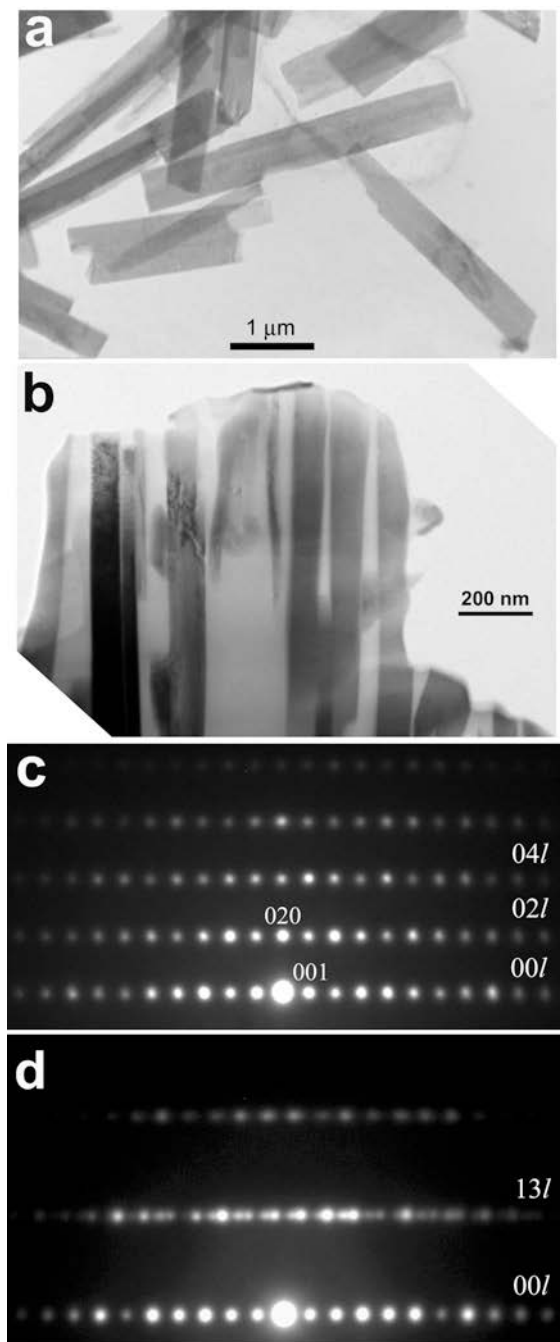


Figure 1. (a) Bright-field image of Zaval'e celadonite, showing a lath- or ribbon-like morphology with the (001) surface and elongation along the  $a$  axis. (b) Bright-field image of the cross-sectional specimen used for selected-area electron diffraction (SAED) and HRTEM. The material between the crystals is epoxy resin embedded during preparation of the sample. (c) SAED pattern from a crystal in (b), with the beam direction parallel to the  $a$  axis of celadonite-1M. Note that the reflection rows are discrete with no streak. (d) SAED pattern from the crystal in (c) but the crystal is rotated by  $30^\circ$  about the  $c^*$  axis, showing extra spots which indicate twinned domains and mutual layer rotation at the twin boundaries with angles of  $(2n+1)60^\circ$ .

imaging plane (Baronnet and Kang, 1989). As described by Kogure and Bunno (2004), the sequence of different symbols ('+' or '-' or '-+') in the image along  $Y_i$ , *i.e.* along the [010], [310], and  $[3\bar{1}0]$  directions for the 1M polytype, indicates stacking with a rotation angle of  $(2n+1)60^\circ$ .

In Figure 2a, the contrast at the octahedral sheet in the 2:1 layer does not correspond to the expected contrast for the *trans*-vacant dioctahedral sheet that was identified by previous structure refinements (Zhukhlistov *et al.*, 1977; Zhukhlistov, 2005; Tsipursky and Drits, 1986). This result occurs by cation migration induced by dehydroxylation. Kogure (2007) suggested that dioctahedral 2:1 phyllosilicates (muscovite, paragonite, pyrophyllite, *etc.*) are dehydroxylated by the electron beam radiation in a TEM. On the other hand, Tsipursky *et al.* (1985), Muller *et al.* (2000a, 2000b), and Drits *et al.* (1995) reported that cation migration from the *trans*-vacant dioctahedral arrangement to the *cis*-vacant arrangement occurs in celadonite by dehydroxylation with heating. From these results, the recorded structure in the HRTEM images is probably not *trans*-vacant dioctahedral as in the unheated celadonite. This phenomenon will be discussed further in a future study. However, this cation migration would not be expected to change the stacking sequence of the layers. In addition, the number of 2:1 layers recorded in Figure 2a is smaller than that in Figure 2b, although the same area was recorded. This result is related to radiation damage. The images in Figure 2b were recorded prior to the specimen being rotated by  $30^\circ$  about the  $c^*$  axis to record Figure 2a. Beam damage before recording Figure 2a destroyed a few layers at the right and left edges of the packet, resulting in an amorphous region. Thus, making the correspondence of layers between the two images in Figure 2 is difficult. Nevertheless, stacking faults in Figure 2b must be contained in the area shown in Figure 2a, although they do not appear in the image. This result requires a rotation angle of  $180^\circ$  at the faults.

Another example of the three-dimensional stacking analysis is shown in Figure 3. The beam direction in Figure 3a is [110] or  $[1\bar{1}0]$  of the 1M polytype. Stacking faults are observed in this image and in Figure 3b. As for Figure 2b, the orientation of each layer was determined from image contrast and denoted as '+' or '-' symbols at the bottom of Figures 3a and 3b (the degree of slant in Figure 3a is  $\pm b/6$ ). The 2:1 layers indicated by the black arrow (top right) in Figures 3a and 3b show a variation in contrast, the reason for which is not clear; the arrows are assumed to indicate identical layers. With this assumption, all layers recorded in the two images can be related to each other. By comparing the sequences of the '+' and '-' symbols in Figures 3a and 3b, the sign of each corresponding layer is reversed. Thus, for a symbol change of '+ -' (or '- +') at an interlayer in Figure 3b (this indicates  $(2n+1)60^\circ$  as described above), the symbols at the corresponding interlayer are '- +' (or

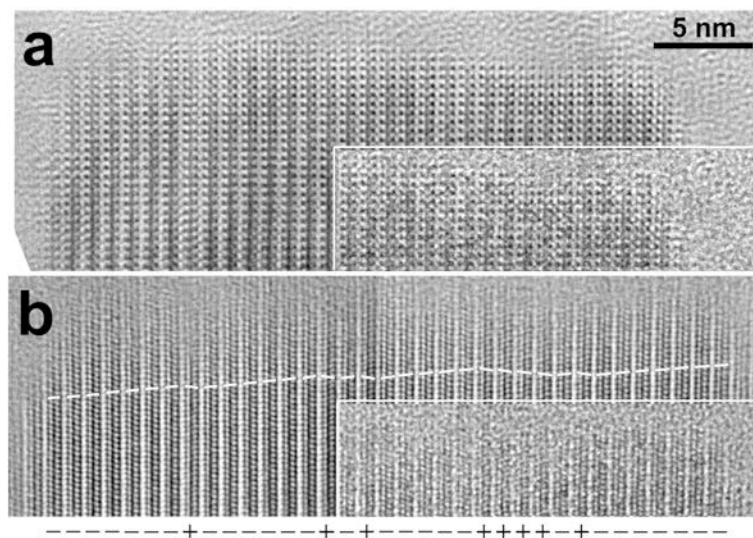


Figure 2. (a) Processed HRTEM image of a crystallite from Krivoi Rog celadonite recorded along the [100] direction. No stacking faults are observed. (b) Processed HRTEM image of the same region but the crystallite is rotated by 30° about the  $c^*$  axis. The white bar in each layer connects the two closest dark spots within the upper and lower tetrahedral sheets in a 2:1 layer, and the sign at the bottom of the figure indicates the slant direction of the bar. The inset at the bottom right of each figure shows the original, unprocessed HRTEM image.

‘+ –’) in Figure 3a. Figure 3c shows the sequence of signs across the interlayer with +60° and 180° rotations, if observed along various  $X_i$  and  $Y_i$  directions (the result with –60° rotation is just a flip of that with +60° rotation horizontally). Two arrows numbered 1 and 2 indicate the intralayer shift with the length of  $a/3$ , for a layer and the overlying layer. In this case, arrow 1 is set parallel to the  $-X_1$  direction. The two symbols, e.g. ‘+ –’, indicate how the two intralayer shifts are imaged in the HRTEM image along the  $X_i$  and  $Y_i$  directions. ‘0’ indicates no slant for the bar, with the intralayer shift that is parallel or anti-parallel to the beam direction. In Figure 3c, the derivation of the symbols for  $Y_i$  directions requires an explanation. If the arrow (the intralayer shift) is perpendicular to the beam direction, the shift is imaged by the amount  $a/3$ . On the other hand, the periodicity of the projection for the 2:1 layer is  $a/2$ , owing to the C-centered unit cell. Consequently, the shift of  $+a/3$  is apparently observed as  $-a/6 (+a/3 - a/2)$  in the HRTEM image recorded along the  $Y_i$  direction. Hence, the shift of  $+a/3$  has the symbol ‘–’. The sequences of ‘+ –’ and ‘– +’ never appear along any  $X_i$  directions for +60° rotation but these sequences do appear for 180°. In conclusion, the stacking sequence in Figure 3 involves only 0° (1M stacking) and 180° layer rotations.

#### XRD analysis

Electron diffraction and three-dimensional stacking analyses using HRTEM for the three celadonite specimens showed identical results. However, the areas investigated by TEM are too limited to conclude that stacking faults with 180° rotation are common throughout the entire specimen. To overcome this limitation, the

effect of stacking faults on the powder XRD pattern was simulated and compared to experimental patterns.

Simulations using *DIFFaX* are similar to that described by Kogure *et al.* (2006). According to the HRTEM results, each celadonite crystal was assumed to consist of two kinds of layers rotated by 180° from one another (this is equivalent to reversal of sign of the  $x$  coordinate of each atom in the orthogonal coordinate system). These two kinds of layers are expressed as layer A and layer B. The ratio of the numbers of the two kinds of layers in a crystal is defined as  $W_A$  and  $W_B$  and it is assumed that layer B is equal to or less than layer A ( $W_B \leq W_A$ ). Next, the probability that layer B succeeds layer A is expressed as  $P_{AB}$ . Similarly  $P_{AA}$ ,  $P_{BB}$ , and  $P_{BA}$  are defined to express the stacking. Of these six parameters, only two are independent because they have the following relations to each other:

$$\begin{aligned} W_A + W_B &= 1, P_{AA} + P_{AB} = 1 \\ P_{BA} + P_{BB} &= 1, P_{AB} = (W_B/W_A)P_{BA} = (W_B/W_A)(1 - P_{BB}) \end{aligned}$$

Among the six parameters, use of  $W_B$  and  $P_{BB}$  to express the stacking is sensible. With the assumption of  $W_B \leq W_A$ ,  $P_{BB}$  can be changed from 0 to 1. If  $P_{BB} = 1$ , the two kinds of layers are completely segregated. If  $P_{BB} = 0$ , the stacking contains a maximum amount of ‘ABA’ sequence, *i.e.* without BB pairs. In the simulation, the atomic parameters for the Zaval’e celadonite (by Tshipursky and Drits, 1986) were used, and a pseudo-Voigt function (the same ratio of Gauss and Lorentz functions) with a full-width at the half maximum of 0.22° was assumed for the peak profile.

Figure 4a shows an example of the simulation where  $W_B$  was fixed at 1/3 and  $P_{BB}$  was varied from 0.0 to 1.0.

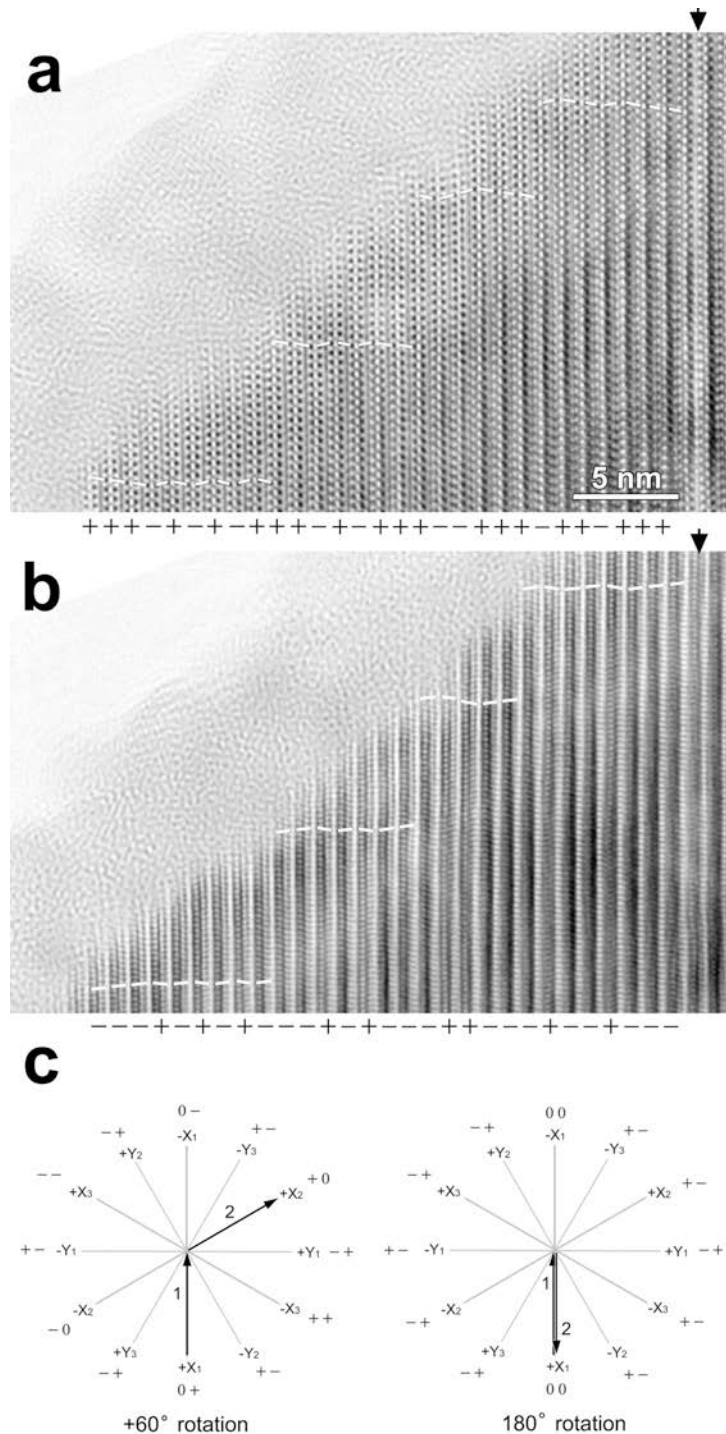


Figure 3. (a) Processed HRTEM image of a crystallite from Krivoi Rog celadonite recorded along the  $[110]$  or  $[\bar{1}\bar{1}0]$  direction. (b) Processed HRTEM image of the same crystal but with the crystallite rotated by  $30^\circ$  about the  $c^*$  axis. The layers indicated by the black arrows in both figures are regarded as identical, owing to their brighter contrast than others. (c) Diagrams showing how the  $60^\circ$  and  $180^\circ$  mutual rotation of the neighboring layers should be observed in the HRTEM images (see the text for details).

The value of  $1/3$  for  $W_B$  seems reasonable because the value is 0.23 and 0.31 in Figures 2 and 3, respectively, and it is 0.46 and 0.32 for the other two grains investigated using HRTEM. Figure 4b represents experi-

mental powder diffraction patterns from the Taiheizan and Zaval'e specimens. The cell dimensions refined from a number of the peak positions in these diffraction patterns are shown in Table 1 with a list of the observed

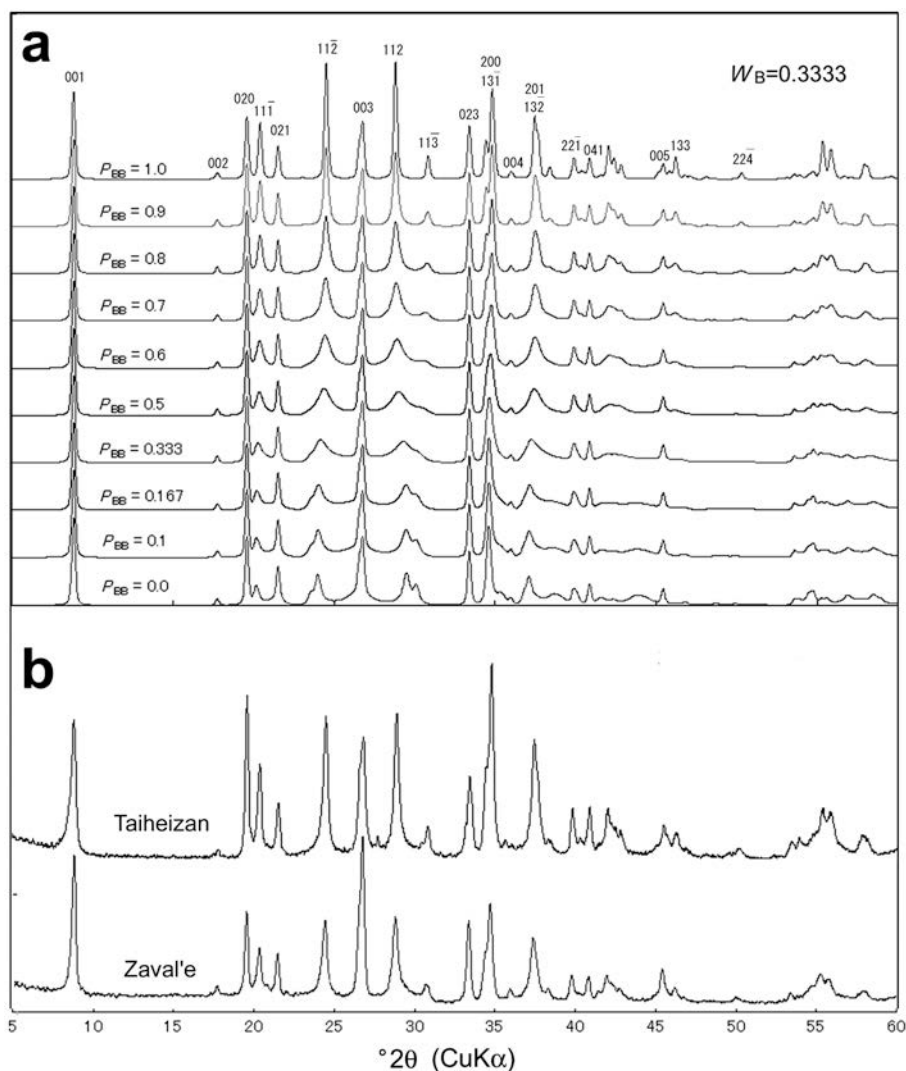


Figure 4. (a) Simulations of the powder XRD patterns from celadonite-1M containing the stacking faults with 180° stacking rotation, using the *DIFFaX* program. Parameters for the simulations,  $W_B$  and  $P_{BB}$  are explained in the text. (b) Experimental powder XRD patterns of celadonite from Taiheizan and Zaval'e.

and calculated  $d$  values for the typical peaks. The sample size of Krovoy Rog was too small to obtain sufficient diffraction intensity. Compared with the simulated pattern, where  $P_{BB} = 1.0$ , which is identical to the pattern for the fault-free 1M structure, the Zaval'e specimen shows a pattern characteristic of a structure which contains 180° stacking faults. Among peaks from 19 to 36°2 $\theta$ , the peak heights of the 11 $\bar{1}$ , 11 $\bar{2}$ , 112, 11 $\bar{3}$ , and (13 $\bar{1}$ , 200) reflections are less than expected when compared to 020, 021, (022, 003), and 023 (the indices in parenthesis indicate peak overlap). If compared with simulated patterns with various  $P_{BB}$  values, the pattern from Zaval'e is close to that with  $P_{BB} = 0.8$  (the largest peak for 003 may be due to preferred orientation). The ratio of the interlayer to form a 180° rotation is expressed as  $W_A P_{AB} + W_B P_{BA}$ . This value is 0.13 for  $W_B = 1/3$  and  $P_{BB} = 0.8$ , indicating that >10% of the

total interlayer regions are faulted with 180° rotation. On the other hand, the pattern of the Taiheizan specimen is close to the simulated pattern with  $P_{BB} = 0.9$ , suggesting more ordered stacking.

## DISCUSSION

### *Some factors responsible for the stacking angle between adjacent layers in micas*

The interlayer region in micas is partly determined by the ditrigonal symmetry of the tetrahedral sheets in the 2:1 layers. This symmetry is caused by the mutual rotation of adjacent tetrahedra in opposite directions around  $c^*$  by the angle  $\alpha$ , the ditrigonal rotation angle. The angle  $\alpha$  is determined by the misfit of the lateral dimensions between the octahedral and tetrahedral sheets. For instance, an increase in Fe and Mg in the

Table 1. Refined cell dimensions from the powder XRD pattern, and observed and calculated peak positions.

| <i>hkl</i>   | — Taiheizan —                |                               | — Zavalé —                   |                               |
|--------------|------------------------------|-------------------------------|------------------------------|-------------------------------|
|              | <i>d</i> <sub>obs.</sub> (Å) | <i>d</i> <sub>calc.</sub> (Å) | <i>d</i> <sub>obs.</sub> (Å) | <i>d</i> <sub>calc.</sub> (Å) |
| 001          | 9.894                        | 9.935                         | 10.121                       | 9.994                         |
| 002          | 4.946                        | 4.968                         | 5.007                        | 4.997                         |
| 020          | 4.510                        | 4.514                         | 4.551                        | 4.538                         |
| 11 $\bar{1}$ | 4.329                        | 4.338                         | 4.362                        | 4.365                         |
| 021          | 4.105                        | 4.110                         | 4.139                        | 4.132                         |
| 11 $\bar{2}$ | 3.613                        | 3.623                         | 3.645                        | 3.643                         |
| 022          | 3.333                        | 3.341                         | 3.351                        | 3.359                         |
| 003          | 3.306                        | 3.312                         | 3.332                        | 3.331                         |
| 112          | 3.074                        | 3.078                         | 3.108                        | 3.100                         |
| 11 $\bar{3}$ | 2.887                        | 2.893                         | 2.903                        | 2.908                         |
| 023          | 2.672                        | 2.670                         | 2.688                        | 2.685                         |
| 20 $\bar{1}$ | 2.593                        | 2.601                         | 2.607                        | 2.620                         |
| 130          |                              | 2.595                         |                              | 2.610                         |
| 13 $\bar{1}$ | 2.565                        | 2.571                         | 2.583                        | 2.585                         |
| 200          |                              | 2.562                         |                              | 2.582                         |
| 004          | 2.484                        | 2.484                         | 2.497                        | 2.499                         |
| 13 $\bar{2}$ | 2.394                        | 2.395                         | 2.408                        | 2.408                         |
| 201          |                              | 2.376                         |                              | 2.394                         |
| 11 $\bar{4}$ | 2.340                        | 2.337                         | 2.349                        | 2.350                         |
| 22 $\bar{1}$ | 2.254                        | 2.254                         | 2.268                        | 2.269                         |
| 041          | 2.200                        | 2.201                         | 2.212                        | 2.212                         |
|              | <i>a</i> (Å)                 | 5.216 (7)                     | <i>a</i> (Å)                 | 5.254 (7)                     |
|              | <i>b</i> (Å)                 | 9.028 (9)                     | <i>b</i> (Å)                 | 9.075 (9)                     |
|              | <i>c</i> (Å)                 | 10.114 (6)                    | <i>c</i> (Å)                 | 10.170 (6)                    |
|              | $\beta$ (°)                  | 100.78 (8)                    | $\beta$ (°)                  | 100.67 (9)                    |

octahedral sheet and a decrease of Al in the tetrahedral sheets of the 2:1 layers in muscovite would accompany a decrease in  $\alpha$  (e.g. Brigatti *et al.*, 1998). The rotation of adjacent layers by  $(2n+1)60^\circ$  should be less favorable because the pairs of the tetrahedral sheets across the interlayer region form a trigonal prism of oxygen atoms around K and create close superposition of the oxygen atoms vertically. However, the smaller the  $\alpha$  angle, the smaller this disadvantage is, because the atomic arrangements with  $2n60^\circ$  and  $(2n+1)60^\circ$  become similar. Hence, the small  $\alpha$  is probably a requirement to form the rotation of  $(2n+1)60^\circ$ . For instance, phlogopite-2O reported by Ferraris *et al.* (2001) has an  $\alpha$  value of  $2.13^\circ$ . This value is considerably smaller than those in other micas. Simulation of experimental XRD patterns from illite-smectite with limited expandability is also in agreement with the  $\alpha$  rotation discussed above (Drits *et al.*, 2006). The greater the amount of Fe and Mg in the octahedral sheet and the smaller the amount of Al in the tetrahedral sheet of the illite-like layers, the greater the probability of  $(2n+1)60^\circ$  stacking faults, probably due to the decrease in  $\alpha$  (Drits *et al.*, 2006).

In addition, the adjacent tetrahedra across the long edges of the *trans*-octahedral site are tilted and the bridging basal oxygen atoms between the tilted tetrahedra move further toward the middle of the layer by  $\Delta z$  with respect to the other basal oxygen atoms of each

tetrahedron. As a result, the basal oxygen plane is corrugated.  $\Delta z$  is near zero for homo-octahedral micas (i.e. the same cations at the three sites in the octahedral sheet). In contrast, the value of  $\Delta z$  is larger for meso-octahedral micas (i.e. different cations at *trans*- and *cis*-sites) (Brigatti and Guggenheim, 2002; Ferraris and Ivaldi, 2002). *Trans*-vacant micas like muscovite are regarded as meso-octahedral. Among the six pairs of oxygen atoms across the interlayer region around K, the oxygen–oxygen distance for the two pairs formed by the depressed oxygen atoms are significantly larger than those for the other four pairs when the stacking angle is  $0^\circ$  or  $180^\circ$ . In contrast, the grooves formed by the rows of the depressed oxygen atoms are rotated with respect to each other in the adjacent layer and therefore each depressed atom has a non-depressed atom as a pair across the interlayer when the stacking angle is  $\pm 60^\circ$  or  $\pm 120^\circ$ . As a result, the distances in the six pairs are equalized significantly and the repulsion between the nearest oxygen atoms across the interlayer would be expected to be less. Thus, it is consistent that muscovite with large  $\alpha$  and  $\Delta z$  values mainly adopts the  $2M_1$  polytype with  $\pm 120^\circ$  layer rotations.

Finally, the amount of Al substituting for Si in the tetrahedral sheets may affect the stacking angles. This substitution changes the lateral dimension of the tetrahedral sheet and  $\alpha$  as a consequence. Furthermore, substitution of Al for Si in the tetrahedral sheets leads to an undersaturated charge on the basal oxygen atoms, so that they have a tendency to repel each other. Thus, the influence of  $\alpha$  and  $\Delta z$  is probably enhanced by large amounts of Al substitution.

#### Structural features of celadonite-1M

The crystal structure of the two celadonite samples, Zaval'e and Krivoi Rog, were refined by Tspursky and Drits (1986) and Zhukhlistov (2005), respectively, using oblique-texture electron diffraction (OTED). According to the structural formulae, the negative layer charge in both structures is located in the octahedral sheets. The tetrahedral sheets have the same cation composition (0.06 atoms of Al per  $O_{10}(OH)_2$ ). Therefore, the interlayer arrangements in both structures are nearly identical. For Zaval'e,  $\alpha = 0.54^\circ$ ,  $\Delta z = 0.030 \text{ \AA}$ ,  $\langle Si, Al-O \rangle = 1.617 \text{ \AA}$ ,  $(K-O)_{\text{inner}} = (K-O)_{\text{outer}} = 3.104 \text{ \AA}$ , and for Krivoi Rog  $\alpha = 0.5^\circ$ ,  $\Delta z = 0.032 \text{ \AA}$ ,  $\langle Si, Al-O \rangle = 1.617 \text{ \AA}$ ,  $(K-O)_{\text{inner}} = 3.084$ ,  $(K-O)_{\text{outer}} = 3.107 \text{ \AA}$ . Thus, in both celadonite structures, tetrahedral rings are nearly hexagonal and the interlayer cation resides in a regular hexagonal prism. Tetrahedral cations and basal oxygen atoms of adjacent tetrahedral sheets are almost superimposed along the  $c^*$  axis. The small  $\alpha$  and  $\Delta z$  values, along with limited Al substitution in the tetrahedral sheets, are characteristic of the celadonite structure. Hence, as discussed above, the preferences for  $2n60^\circ$  by ditrigonal symmetry and for  $\pm 60^\circ$  or  $\pm 120^\circ$  by a  $\Delta z$  value are not expected in the structure and all

rotation angles may occur equally. However, as described above, the celadonite specimens contain only 180° rotation as stacking faults in the 1M stacking.

One of the possible factors responsible for the absence of  $\pm 60^\circ$  and  $\pm 120^\circ$  stacking rotation in celadonite may be related to the lateral periodicity of the layers. For example, if  $b/a \neq \sqrt{3}$ , both rotated (except 180°) and non-rotated layers have a different orthogonal unit cell. Thus, the two-dimensional (2D) periodicities within the rotated layer pair will be non-commensurate. To adjust to a common 2D periodicity, the structure of the rotated layers will be modified slightly creating a certain strain which may restrict the occurrence of stacking faults with  $\pm 60^\circ$  or  $\pm 120^\circ$  rotation. However, in the studied samples,  $b/a$  is 1.731 for Taiheizan and 1.727 for Zaval'e (Table 1). These values are so close to  $\sqrt{3}$  that the stacking faults with  $\pm 60^\circ$  or  $\pm 120^\circ$  rotation will not interfere with 2D periodicity of the stacking sequence.

#### *The relationship between stacking sequence and ribbon-like morphology of celadonite*

Zhukhlistov *et al.* (1979) reported the occurrence of two types of celadonite crystals with different morphologies from Krivoi Rog. One has a plate-like shape, which was formed by metasomatic reaction from Fe-rich biotite. Another has a large ribbon-like shape, which grew in an empty space by direct crystallization in solution. They are similar in chemical composition and cell dimensions, but OTED indicated that the stacking of the plate-like crystals is more disordered than the ribbon-like crystals. Zhukhlistov *et al.* (1977) and Zhukhlistov (2005) used the latter type for the structure analysis by OTED and the present study also investigated this type. The result of Zhukhlistov *et al.* (1979) suggests that the stacking feature of celadonite is significantly influenced by the condition during crystal growth.

Güven (2001) considered the relationships between the morphological features of *trans*-vacant (*tv*) 1M mica crystallites and the relative growth rates along different lateral directions within the mica layer. He concluded that a pseudo-hexagonal plate-like habit of the *tv* 1M crystal develops when the growth rates are equal on the (010), (110), and  $(\bar{1}10)$  fronts, and that laths or fibers develop when the growth rate on the (010) fronts is slower than on the (110) and  $(\bar{1}10)$  fronts. He showed that the atomic configurations exposed on the (010), (110), and  $(\bar{1}10)$  planes are different: the (010) plane exposes pairs of octahedral OH groups whereas the (110) and  $(\bar{1}10)$  planes expose octahedral *cis*-sites occupied by cations and numerous broken bonds (Güven, 2001). The (110) and  $(\bar{1}10)$  fronts are therefore expected to grow faster than the (010) fronts, resulting in laths or fibers with their longest dimensions parallel to the *a* axis. According to Güven (2001), not only the crystal but the *tv* 2:1 unit layer itself also has such a growth character. This was also confirmed experimentally by Baronnet *et*

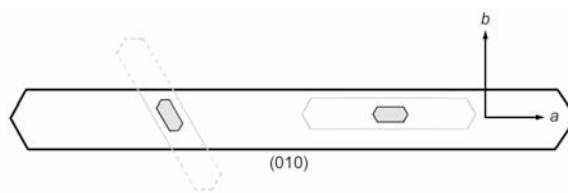


Figure 5. A model to show limited (left) and unlimited (right) layer growth of nuclei with  $(\pm 60^\circ, \pm 120^\circ)$  and  $(0^\circ, 180^\circ)$  layer rotations, respectively, on a ribbon-like crystal. See the text for details.

*al.* (1976). Crystallites are expected in which constituent *tv* 2:1 layers have parallel (1M stacking) or antiparallel (180° layer rotation) orientations which can grow more rapidly than those containing layers with  $\pm 60^\circ$  and/or  $\pm 120^\circ$  layer rotations. As a consequence, the crystals with 1M and only 180° stacking faults may occupy larger volumes than those containing  $\pm 60^\circ$  and/or  $\pm 120^\circ$  stacking faults in the specimen. This is illustrated in Figure 5. Suppose, for example, that several new layers with various orientations are nucleating on a crystal with a ribbon-like form. These new layers are also assumed to have the ability to elongate along the *a* axis (note that this is the axis for the two-dimensional unit layer), as mentioned above. In this situation, new layers with  $\pm 60^\circ$  or  $\pm 120^\circ$  rotations (the left one) expand on the preceding layer along the *a* axis but soon reach the side-end of the ribbon-like crystal. Then, growth occurs only on the (010) fronts and becomes very slow. On the other hand, new layers with  $0^\circ$  and  $180^\circ$  (the right one) can expand more widely owing to the parallelism with the preceding layer or crystal. As a consequence, the latter is dominant on the crystal and a large portion of the crystal has stacking with  $0^\circ$  (1M stacking) and  $180^\circ$  layer rotation. Such a quantitative dominance may be reflected in the XRD and HRTEM results in the present investigation.

The morphology of illite has been shown experimentally to be influenced considerably by the growth conditions (Güven, 2001), and can thus affect the dominant stacking sequence through the change of morphology. For example, glauconite has a similar composition to celadonite except that it is deficient in interlayer cations (Rieder *et al.*, 1999). Glauconite often occurs in marine sediments and its crystallites grow very slowly by direct precipitation at low temperature in diagenetic environments (Odom, 1984). Under such conditions, *tv* 2:1 layers may grow with a plate-like form rather than a ribbon-like form. Sakharov *et al.* (1990) noted that stacking faults involving  $2n60^\circ$  and  $(2n+1)60^\circ$  layer rotations occur with equal probability in glauconite, and these can be attributed to the plate-like form of the *tv* 2:1 layers in such growth environments.

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