

# GaN GROWTH BY REMOTE PLASMA MOCVD: CHEMISTRY AND KINETICS BY REAL TIME ELLIPSOMETRY

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## ABSTRACT

Cubic and hexagonal GaN layers have been grown on GaAs (001) and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (0001) substrates, respectively, by remote plasma metalorganic chemical vapor deposition (RP-MOCVD). In situ spectroscopic ellipsometry is used to monitor in real time the chemistry and kinetics of the GaN growth. The substrate/GaN interface formation is highlighted and the effect of the substrate plasma nitridation on the initial growth stage is discussed.

## INTRODUCTION

There have been numerous recent publications on the effects of the two growth steps, i.e., substrate nitridation followed by the growth of a low temperature nucleation layer, upon the material properties of GaN epilayers grown by both molecular beam epitaxy (MBE) and metalorganic chemical vapor deposition (MOCVD) [1,2]. And, it has been shown that slight differences in the substrate nitridation and/or the growth of the GaN buffer layer prior to GaN epigrowth results in different GaN growth mechanisms [3,4]. Although the number of experimental studies conducted to date is already significant, it is not yet possible to establish an all-embracing theory of the influence of the surface properties upon the GaN growth chemistry and kinetics. And, it is widely recognized that a deeper understanding of the fundamental surface kinetics governing nitride growth is necessary to move from the "GaN art" to the "GaN science". While the investigation of the GaN growth mechanism is ongoing for MBE systems, where a variety of in situ diagnostic techniques can be used, such as reflection high energy electron diffraction (RHEED), low energy electron diffraction (LEED), X-ray photoelectron spectroscopy (XPS), some difficulties are encountered for the MOCVD process, where high pressure and reactive environment prevent from the use of the above techniques.

In this paper, the remote plasma MOCVD (RP-MOCVD) of GaN buffer layers grown on both  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and GaAs substrates, also pre-nitrided by N<sub>2</sub> plasmas, is discussed. Particular emphasis is placed on the in situ real time control of the growing surface and of the substrate/GaN interface by spectroscopic ellipsometry (SE), which is a non-invasive and non-intrusive optical technique compatible with the MOCVD environment. Ellipsometry data are used to depict the chemistry and kinetics of the initial stage of the GaN growth.

## EXPERIMENT

The deposition system used in this study is the remote plasma MOCVD apparatus specifically designed for the growth and treatment of III-V materials [5] at reduced pressure (0.1 - 10 Torr) and reduced temperature. A unique feature of this apparatus is the presence of an in situ phase modulated ellipsometer (UVISEL-ISA Jobin-Yvon) compatible with the reactive environment of MOCVD reactors and, hence, able to monitor in real time the chemistry and kinetics of surface modifications with a sub-monolayer resolution. The substrates used for the growth were GaAs (001) and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (0001). The remote r.f. (13.56 MHz) plasma source was

used to produce H-atoms for the substrates cleaning, and N-atoms for the nitridation process and GaN growth. The H<sub>2</sub> remote plasma cleaning was operated at P = 1 Torr, r.f. power = 60 Watt and at a surface temperature of 250°C and 350°C for GaAs (001) and α-Al<sub>2</sub>O<sub>3</sub> (0001), respectively. The subsequent nitridation by N<sub>2</sub>-H<sub>2</sub> (3% in H<sub>2</sub>) plasmas was operated at a pressure of 0.2 Torr, a r.f. power of 200 Watt and at a temperature of 250°C and 800°C for GaAs and α-Al<sub>2</sub>O<sub>3</sub>, respectively. GaN layers were grown at 600°C by trimethylgallium and N<sub>2</sub>/H<sub>2</sub> (=1000/3 sccm) plasma at a pressure of 1 Torr and a r.f. power of 200 W. Single wavelength ellipsometry (SWE) was used, in terms of the ellipsometric angles Ψ and Δ [6], to monitor in real time the surface kinetics, and, at the growth end, ellipsometric spectra (SE) of the pseudodielectric function, <math>\langle \epsilon \rangle = \langle \epsilon\_1 \rangle + i \langle \epsilon\_2 \rangle</math>, were acquired in the energy range 1.5 - 5.5 eV. In order to determine the chemical and optical properties of surfaces, ellipsometric data were modeled by optical models based on the Bruggeman effective-medium approximation (BEMA) [6] and references for the used dielectric functions are given in ref. [7]. Optical emission spectroscopy (OES) was used to control the N-atom density interacting with the growth surface. Ex situ atomic force microscopy (AFM) and X-ray diffraction (XRD) measurements were also performed to validate the ellipsometric approach.

## RESULTS AND DISCUSSION

Figure 1 shows typical SWE real time (Ψ,Δ) trajectories recorded at 3 eV during GaN growth on both α-Al<sub>2</sub>O<sub>3</sub> (0001) and GaAs (001) substrates previously nitrided. An easily provided information is the GaN layer thickness which is controlled in real time.

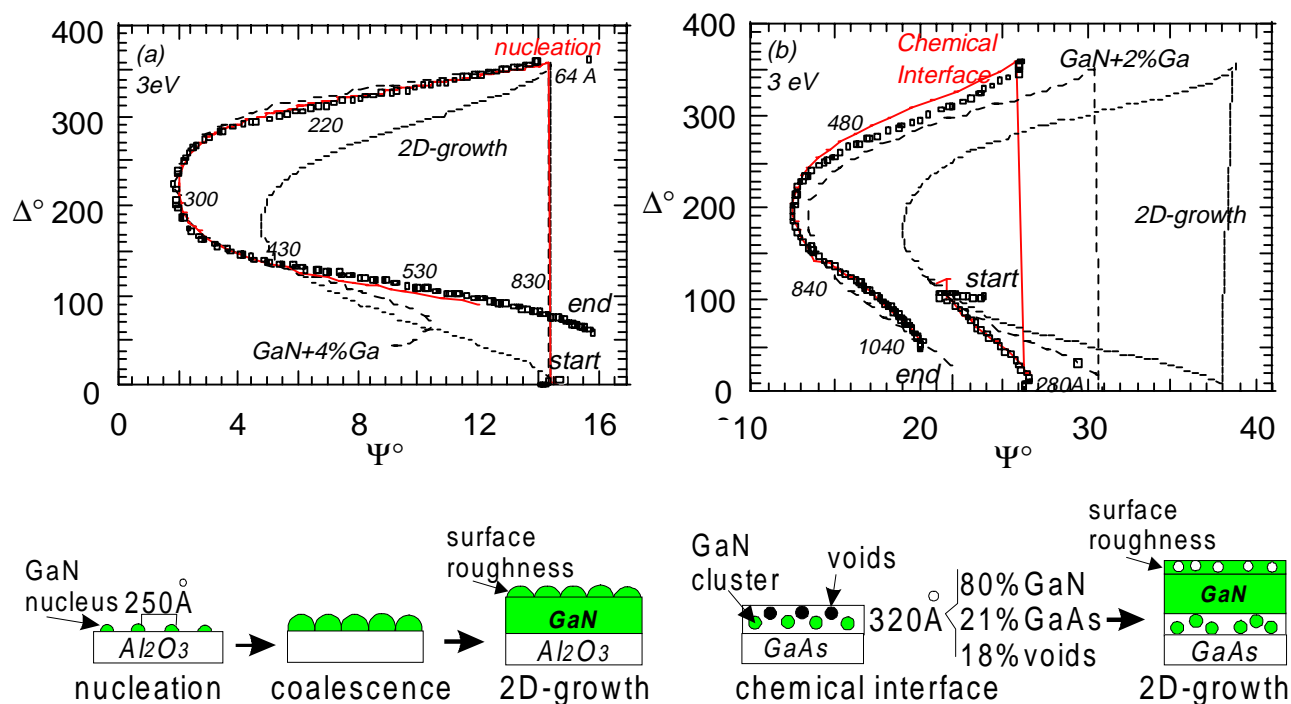


FIGURE 1. Time evolution of the optical response in the (Ψ,Δ) space at 3 eV during GaN growth on (a) α-Al<sub>2</sub>O<sub>3</sub> and (b) GaAs. Points are for experimental data; full line is for the best-fit model which is, for GaN on α-Al<sub>2</sub>O<sub>3</sub>, nucleation followed by 2D-homogeneous growth and, for GaN on GaAs, chemical interface appearance followed by 2D-homogeneous growth. Simulated trajectories by different models (see text) are also reported for comparison.

For a better understanding of the ( $\Psi, \Delta$ ) data, it has to be considered that at 3 eV (below energy gap value, where the GaN absorption is almost zero) a closed and cyclical trajectory should be recorded for a well ordered and pure GaN layer. Thus, deviations from this ideal cyclical curve are indicative of non-homogeneous growth dynamics and/or defected absorbing layers. For the ( $\Psi, \Delta$ ) trajectories simulation, different models have been considered such as the 2D-homogeneous growth, the nucleation followed by homogeneous growth and the stoichiometry deviation in terms of Ga enrichment. Among these models, only the “nucleation followed by homogeneous growth with roughness” model fits the whole experimental trajectory in fig. 1(a). Specifically, the growth process of the GaN epilayer on  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> can be distinguished into two regimes, i.e., an initial hemispherical nucleation regime followed by the GaN bulk growth (see scheme in fig. 1a). The nucleation density on the starting surface produces a mean separation of  $\sim 250$  Å between nuclei, whose radius increases according to the law  $R(t) = 1.8(\text{Å}/\text{sec}) \cdot t$  until they come into contact. As soon as the complete nuclei coalescence is obtained, the homogeneous GaN growth starts with a steady-state growth rate of about 1 Å/sec and a gradual roughening of the growing surface also confirmed by SE modeling (see below) and AFM measurements.

For the GaN growth on GaAs (fig. 1b), the regression analysis fit evidences the appearance of a chemical interface between the substrate and the GaN film. For the interface formation regime, free fit parameters include the volume fractions of the interface components and the thickness of the interface layer. The simulation shows that the interface layer is modeled as a three-component GaAs/GaN/voids composite whose composition varies with time, i.e., the GaN volume fraction increases at the expense of GaAs until 100% GaN is reached. Then, the GaN bulk 2D-growth starts, and the GaN layer thickness increases linearly with time at about 1 Å/sec. The voids content at the interface simulates either the presence of structural defects and the grain boundaries.

Figure 2 shows the experimental SE spectra, at room temperature, of the real,  $\langle \epsilon_1 \rangle$ , and imaginary,  $\langle \epsilon_2 \rangle$  part of the pseudodielectric function of the GaN buffer layers recorded at the end of the growth. In the figure, the solid lines are for the fitting results according to the models shown in the inset. A two-layer model (substrate/bulk GaN/surface roughness) provides the best fit for the SE spectrum of GaN grown on  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>. The surface roughness developed during the growth is simulated by an outer layer which is a mixture of GaN+voids as depicted in fig. 2.

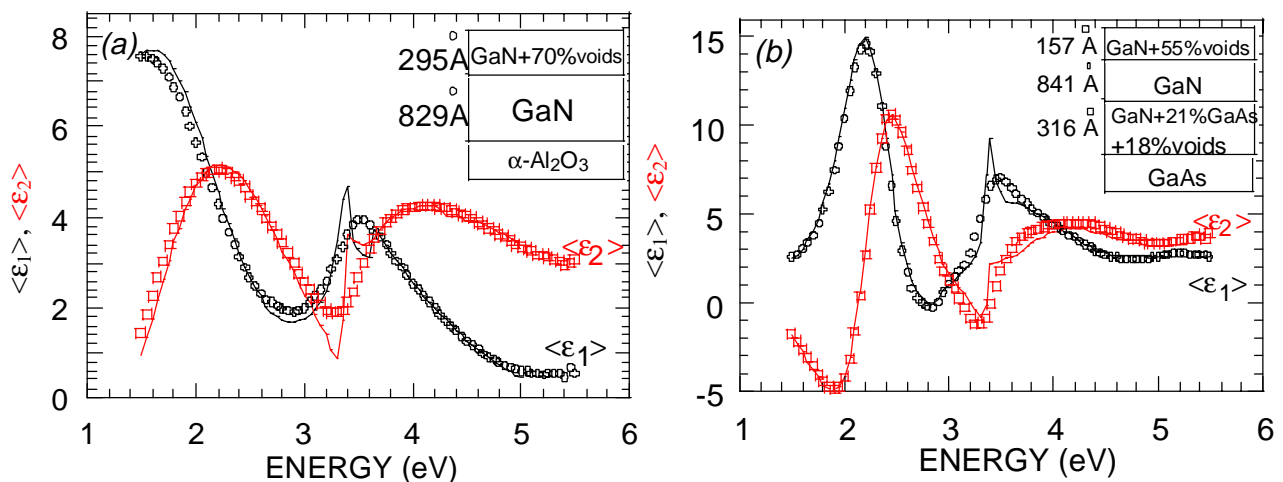


FIGURE 2. Real,  $\langle \epsilon_1 \rangle$ , and imaginary,  $\langle \epsilon_2 \rangle$  part of the pseudodielectric function of the GaN layers grown on (a)  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and (b) GaAs. The best-fit curves (full line) by the BEMA models shown in the inset are also reported. The standard deviations are of  $\pm 1\%$  on constituents volume fraction and of  $\pm 10\%$  on layer thickness.

The best-fit model for the SE spectrum of GaN on GaAs employs a three-layer model (substrate/interface/bulk GaN/surface roughness) that accounts for the growth dynamics of fig. 1b: at first the interface formation between GaAs substrate and GaN film, and then the bulk GaN film growth. The evolution of the interface composition is approximated in the SE model by an average value (BEMA mixture of GaN+21% GaAs+18% voids). The comparison of the fit and experimental data in both GaN/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> and GaN/GaAs evidence the absence of a sharp excitonic transition, so indicating a very defective GaN buffer layer. Interestingly, cathodoluminescence (CL) measurements have shown a very low intensity of the near-band edge emission and a well evident deep level emission at around 550 nm. Also, AFM images have shown a grain-like morphology with an average diameter of  $\sim$ 500 Å. These features have already been described by Tarsa et al.[8] and, they are, accordingly, related to the growth dynamics under N-stable conditions.

Thus, the amount of N-atoms impinging on the growing films is a crucial parameter for the GaN growth. In the present RP-MOCVD apparatus, the performance of the nitrogen remote plasma source has been evaluated by OES measurements, performed in the afterglow at the substrate position. There, the OES spectra are dominated by the N<sub>2</sub> 1st positive system [7], which indicates a considerable amount of N-atoms and of electronically and vibrationally excited N<sub>2</sub> molecules impinging on the growth surface. Any contribute of ions is absent and the importance of this can be understood by considering that the ionized portion of the N<sub>2</sub> plasma beam has been suggested to have detrimental effect on the GaN film morphology and crystallinity [9].

Figure 3 shows the XRD measurements for GaN on both  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (0001) and GaAs (001) substrates from which the crystallographic structure has been derived. The XRD pattern of GaN grown on GaAs shows only the (002) and (004) reflections, so demonstrating this film to be purely cubic and (001) oriented monocrystalline. The FWHM of the (002) reflection is 22 min for a film  $\sim$ 100 nm thick, and this FWHM value is, to our knowledge, among the narrowest reported so far [10,11]. The GaN (002) reflection at 39.95° corresponds to the GaN with a lattice constant of 4.509 Å, which is in good agreement with the calculated value of 4.503 Å [12]. Also, from XRD, there is no evidence of a ternary alloy GaAs<sub>x</sub>N<sub>1-x</sub> at the interface, and a detailed analysis of the (002) peak shape, which is sensitive to compositional disorder, suggests that, indeed, phase separation of GaAs and GaN [11] does seem to occur at the interface, in agreement with the ellipsometric data. Here, it is important to underline that, in the present case, the pure cubic phase is grown without arsenic background pressure in the initial stage as reported by others [10,11,13].

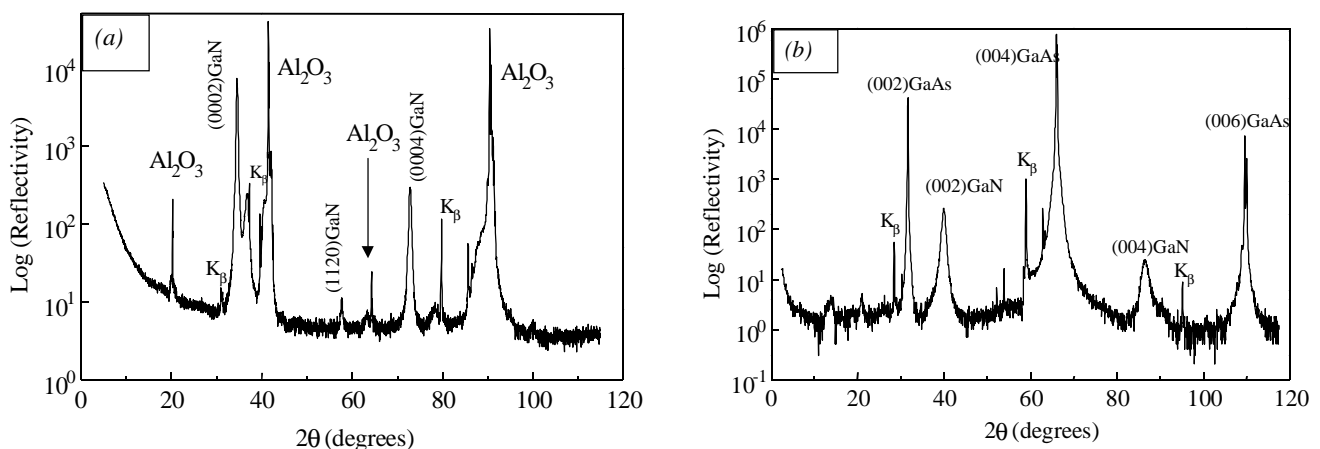


FIGURE 3. Typical XRD patterns of (a) GaN/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (0001) and (b) GaN/GaAs(001) samples.

In the present case, we have grown pure cubic GaN on a plasma nitrated GaAs surface differently from other workers [14,15] who obtained hexagonal GaN on nitrated GaAs. This capability of our method is probably related to the fact that the remote plasma nitridation of the GaAs substrate occurs in absence of ion bombardment and under controlled conditions, so giving a very smooth (RMS = 3.1 Å) and homogeneous nitrated layer of 130 Å in thickness [7,16], which acts as a template for the subsequent growth of the cubic phase. However, from the SE analysis of the GaN layer grown at 600°C (see model in fig. 2b), there is no memory of the layer coming from the pre-nitridation of the GaAs surface at 250°C, i.e., the nitrated layer can not be discriminate after growth. This is, probably, because of its recrystallization and partial decomposition of the underlying GaAs surface during rise to growth temperature of 600°C, thus resulting in the “new” 316 Å interface layer of (GaN+21%GaAs+18% voids).

The XRD pattern of GaN grown on the (0001) plane of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (fig. 3a) shows the main peak at 34.55° corresponding to the (0002) reflection from wurtzite GaN with a lattice constant in the c-direction of 5.18 Å [12]. This confirms that the film is grown with its c-plane parallel to the substrate. The (0002) peak has a FWHM value of 8.6 min, which is well comparable with the best value of ~ 5 min reported in literature [17]. The feasibility of our system for the deposition of GaN buffer layer is the consequence of the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nitridation which has been performed under real time control by ellipsometry. In fact,  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nitridation leads to different surface chemistry and morphology depending on exposure time [18,19]. Figure 4 shows the time evolution of the  $\langle \epsilon_2 \rangle$  value during  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nitridation. The decrease of  $\langle \epsilon_2 \rangle$  is indicative of nitrogen incorporation, and the SE spectrum (spectrum B) recorded at the  $\langle \epsilon_2 \rangle$  minimum can be modeled by a thin (~7 Å) layer of a BEMA mixture of AlN+10% Al<sub>2</sub>O<sub>3</sub> which, however, can optically simulate an AlNO<sub>x</sub> layer also. AFM measurements of this Al<sub>2</sub>O<sub>3</sub> nitrated surface show an atomically flat surface with a RMS value of 1.6 Å. For longer exposure time (>10 min in the figure), the  $\langle \epsilon_2 \rangle$  trend reverses as a consequence of roughening and of free aluminum formation, this last also confirmed by XPS measurements. The formation of free Al, besides roughening [19], could be the reason why long nitridation time results in a degradation of GaN layer quality [18,19].

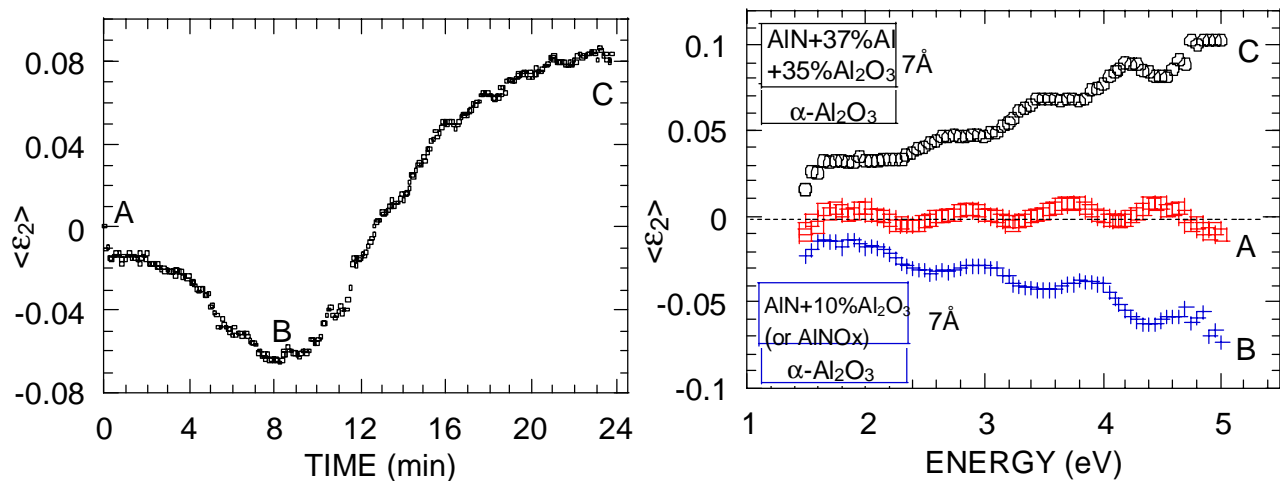


FIGURE 4. Time evolution of the  $\langle \epsilon_2 \rangle$  value at 4.5 eV during  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nitridation (on the left) and SE spectra (on the right) at significant points of (A) H<sub>2</sub> plasma cleaned  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> surface, (B) 8 min nitrated  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> surface, and (C) 25min nitrated and damaged  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>. The inset shows the best-fit BEMA models of spectra B and C.



Thus, the real time control by SE of the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nitridation has allowed the GaN growth to start soon after the minimum in  $\langle \epsilon_2 \rangle$  (end-point of nitridation) has been reached.

## CONCLUSIONS

In conclusion, spectroscopic ellipsometry has been applied as a real-time probe for the remote plasma MOCVD of GaN film in order to understand and optimize GaN growth and substrates nitridation processes. Ellipsometry has been shown to provide information on the end-point of the substrate (GaAs and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>) nitridation, on the substrate/GaN film interface and on the growth kinetics. The RP-MOCVD growth of GaN under N<sub>2</sub> remote plasma preactivation conditions results in cubic GaN on GaAs and in hexagonal GaN on  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>.

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