

GRAVITATIONAL CHANNELS AND THE COSMOLOGICAL CONSTANT Λ

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A new model of filamentary matter fields and voids is proposed. This is a gravitational version of the MIT bag model of hadrons (see a review of DeTar and Donoghue 1983). Bekenstein and Milgrom(1984) have first proposed a gravitational bag model. Their bag is closed but our bag, called a channel, is open.

We start from the action in the weak limit of gravity

$$I = \int_{-\infty}^{\infty} dt \int d\mathbf{x} \left\{ \sum_n m_n \left[\frac{1}{2} \dot{\mathbf{x}}_n^2 - \phi(\mathbf{x}_n) \right] \delta(\mathbf{x} - \mathbf{x}_n) - \frac{1}{8\pi G} (\nabla\phi)^2 + B \right\}, \quad (1)$$

where ϕ is the gravitational potential and $B = -c^4 \Lambda / (8\pi G)$. Varying $\mathbf{x}_n(t)$, $\phi(\mathbf{x})$ and V under the conditions that $\delta \mathbf{x}_n(t) \rightarrow 0$ as $|t| \rightarrow \infty$ and that $\delta \phi(\mathbf{x}) \rightarrow 0$ as $|\mathbf{x}| \rightarrow \infty$, we obtain the equations of motion for masses, Poisson equation for ϕ

$$\nabla^2 \phi = 4\pi G \sum_n m_n \delta(\mathbf{x} - \mathbf{x}_n), \quad (2)$$

and the natural boundary conditions at the free surface (S_1 in Figure 1)

$$\frac{\partial \phi}{\partial n} = 0 \quad \text{and} \quad \frac{\partial \phi}{\partial t} = U, \quad (3)$$

where $U = (8\pi G B)^{1/2} = c^2 (-\Lambda)^{1/2}$.

The condition $\frac{\partial \phi}{\partial n} = 0$ means that the lines of force are confined within the channels, i.e. filamentary structures. The potential ϕ is equivalent to the velocity potential in the incompressible irrotational flow

which is surrounded by a gas of constant pressure and has the sources of total flux $4\pi G m_n$ at $\mathbf{x} = \mathbf{x}_n$. The radius r_0 of the distant cross-section (S_2 in Figure 1) is obtained from Gauss's Theorem $2\pi r_0^2 U = 4\pi G m$, so that $r_0 = (Gm^2/2\pi B)^{1/4}$.

The exact solution for the two-dimensional channel from a point mass has been obtained using the standard method of conformal mapping (Milne-Thompson 1938). The force F on a test particle at the center line changes from $F \propto 1/r$ (two-dimensionally Newtonian) to $F = \text{const.}$ at a distance of the half-width of the channel. This would be the case with the three-dimensional channel from a disk galaxy.

Assuming that the law of force is $-Gm_1 m_2 (1 + r_{12}^2/A^2)^{1/2} / r_{12}^2$ between two masses in the disks of nearby spirals (M31, M33, NGC2403, NGC6946), we have found, from fitting the computed and observed rotation curves, that the scale length A satisfies a close relation $A = 2.7 \text{ kpc} * (m/10^{10} M_{\odot})^{1/2}$, where m is the total mass of a spiral. The acceleration $Gm/A^2 = 1.9 \times 10^{-8} \text{ cms}^{-2}$ is coincident with Milgrom's(1983) limiting acceleration if $H_0 = 50 \text{ kms}^{-1} \text{ Mpc}^{-1}$. Equating $r_0 = A$, we have $\Lambda = -1.8 \times 10^{-5} \text{ cm}^{-2}$.

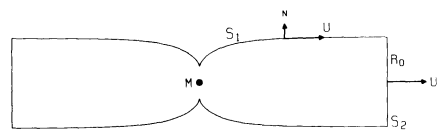


Figure 1. An axisymmetric channel of two opposite equal fluxes from a point mass.

References

Bekenstein, J. and Milgrom, M.:1984, Astrophys. J. **286**, 7.
 Milgrom, M.: 1983, Astrophys. J. **270**, 365, 371, and 384.