Characterization and Rietveld refinements of new dense ceramics $Ba_{3-x}Sr_{x}Tb_{3-x}Ce_{x}O_{9}$ (x = 1 and 1.5) perovskites

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 $Ba_{3-x}Sr_xTb_{3-x}Ce_xO_9$ (x = 1 and 1.5) ceramics (BSTC) with a relative density of 93% and a grain size distribution of $0.2-3 \,\mu\text{m}$ were prepared by the mixed-oxides reaction route. The crystalline structures, microstructures, valence states, and electrical properties of two ceramics were analyzed using X-ray powder diffraction (XRPD), scanning electron microscopy (SEM), X-ray photoelectron spectroscopy (XPS), electron paramagnetic resonance (EPR), and electrical measurements. Rietveld analyses of XRPD patterns show that BSTC1 is indexed as a trigonal structure with the space group R-3c, and BSTC3/2 is indexed as an orthorhombic perovskite structure with the space group Pmcn. The EPR, XPS, and electrical conductivity results confirm that Ce and Tb ions in BSTC exist as Ce⁴⁺ and mixed-valence states of Tb^{4+}/Tb^{3+} , respectively. At room temperature, the two BSTC ceramics exhibit a similar semiconducting behavior. The relationships between electrical conductivity and temperature/frequency are provided. The defect chemistry is discussed. © 2020 International Centre for Diffraction Data. [doi:10.1017/S0885715620000056]

Key words: $Ba_2SrTb_2CeO_9$ and $Ba_{1.5}Sr_{1.5}Tb_{1.5}Ce_{1.5}O_9$, X-ray powder diffraction, electric properties, electron paramagnetic resonance, X-ray photoelectron spectroscopy

I. INTRODUCTION

Perovskite-type oxides, such as BaCeO₃ and SrCeO₃, are state-of-the-art high-temperature proton conductors as electrolytes in solid oxide fuel cells (SOFC) because of its high protonic conductivity when exposed to a humidified hydrogencontaining atmosphere at temperatures higher than 300 °C (Uchida et al., 1983; Scherban et al., 1988; Iwahara et al., 2004; Tolchard and Grande, 2007; Fu and Weng, 2014; Knight et al., 2015). In order to improve proton conductivity, several kinds of rare-earth-doped BaCeO3 and SrCeO3 ceramics have been developed: (1) BaCe_{1-x}RE_xO_{3- δ}, via the B-site replacement of Ce⁴⁺ by acceptor-type trivalent rare-earth ions (RE³⁺) such as RE³⁺ = Y³⁺, Pr³⁺, Nd³⁺, Sm³⁺, Gd³⁺, Eu³⁺, Tb³⁺, and Yb³⁺ (Matsumoto et al., 1999; Wang et al., 2004; Wu et al., 2004; Sharova et al., 2005; Malavasi et al., 2008); (2) Ba (Ce_{0.8-v}Pr_vGd_{0.2})O_{2.9}, via the B-site co-doping with double rare-earth ions Pr³⁺ and Gd³⁺ (Mukundan et al., 2001); and (3) BaCe_{0.2}Zr_{0.7}RE_{0.1}O_{3- δ}, via the B-site co-doping with Zr⁴⁺ and RE^{3+} (= Y^{3+} , Sm^{3+}) (Barison *et al.*, 2008; Ricote *et al.*, 2012; Kannan et al., 2013; Choi et al., 2014). A more complex system $Ba_{1-x}Sr_xCe_{0.5}Zr_{0.35}Y_{0.1}Sm_{0.05}O_{3-\delta}$ with a cubic perovskite structure (the space group Pm-3m), in which Sr^{2+} and Zr⁴⁺/Y³⁺/Sm³⁺ are incorporated into the A- and B-sites in BaCeO₃, respectively, shows high conductivity and highdensity proton conductivity (Radenahmad et al., 2016). From another perspective, the commercial application of materials has been hindered by the technical difficulty of

creating fully dense ceramics with good mechanical properties (Dahl et al., 2007). Thus, dense SrCeO₃- or BaCeO₃-based ceramics with simultaneous occupations of both the A- and B-sites may be promising candidates for high-temperature proton conductors in SOFC.

In this work, two ceramics with nominal compositions $Ba_{1-x/3}Sr_{x/3}Tb_{1-x/3}Ce_{x/3}O_3$ (x = 1 and 1.5) (BSTC) were prepared using the mixed-oxides method. The refined crystal structures of BSTC are identified using Rietveld refinements. They exhibit higher densification and have perovskite structures with the general formula Ba₂SrTb₂CeO₉ and Ba_{1.5}Sr_{1.5}Tb_{1.5}Ce_{1.5}O₉, similar to Sr₂CaFe₂WO₉ and Sr₂PbFe₂TeO₉ perovskites, respectively (El Hachmi et al., 2018). Scanning electron microscopy (SEM), X-ray photoelectron spectroscopy (XPS), electron paramagnetic resonance (EPR), and electrical measurements were employed to discuss the crystalline structure and valence states of rare-earth ions Ce and Tb.

II. EXPERIMENTAL

A. Synthesis

Ceramic raw materials were reagent-grade BaCO₃ (99.5%), SrCO₃ (99.5%), Tb₄O₇ (99.9%), and CeO₂ (99.9%) powders. Two ceramics were prepared according to the nominal formulas of $Ba_{1-x/3}Sr_{x/3}Tb_{1-x/3}Ce_{x/3}O_3$ (x = 1 and 1.5) (abbreviated as BSTC1 and BSTC3/2, respectively) using a conventional mixed-oxides method. The stoichiometric mixtures in accordance with the above metal ratios Ba:Sr: Tb:Ce were carefully ground in an agate mortar, and then were calcined at 1100 °C for 5 h in air. After furnace cooling, the

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calcined mixtures were reground. Aqueous PVA (polyvinyl alcohol) solution was added into the calcined mixtures. The resulting powders were pressed uniaxially at 200 MPa into disk-like pellets of 12 mm in diameter. The discs were placed in a pile on the bottom of the Al₂O₃ crucible in an electric furnace. The discs were heated at a rate of 100 °C h⁻¹ and sintered at 1400 °C for 12 h in air, cooled at a rate of -200 °C h⁻¹ to 700 °C, and then furnace cooled to room temperature (RT) to form ceramics. In addition, BaCeO₃, BaTbO₃, SrCeO₃, and SrTbO₃ ceramics were prepared under the same conditions as BSTC.

B. X-ray powder diffraction

Diffraction data were collected at RT on a DX-2700 X-ray diffractometer (Dandong Haoyuan), with the Bragg–Brentano geometry, using Cu*K* α radiation ($\lambda = 1.5418$ Å) with 35 kV and 30 mA, divergence slit of 1°; anti-scatter slit of 1°; receiving slit of 0.3 mm. The diffraction pattern was scanned between 5° $\leq 2\theta \leq 120^{\circ}$ in 0.02° 2 θ intervals with a fixed-time counting of 3 s step⁻¹.

The Rietveld refinements were used to refine the crystal structure. The peak shape was described by a pseudo-Voigt function. The starting data needed for Rietveld refinements for BSTC1 and BSTC3/2 are the atomic positions and unit-cell parameters from BaCeO₃ (88592-ICSD) and SrCeO₃ (PDF#01-074-8250), respectively.

C. Scanning electron microscopy, backscattered electron, and energy-dispersive X-ray spectroscopy

Microstructures of polished and thermally etched ceramic surfaces were examined using an EVOMA 10 scanning electric microscope (SEM) (Zeiss) operated at 15 kV. To confirm the nature of single phase and the absence of secondary phases in BSTC, SEM investigations in the backscattered electron (BSE) mode were performed. Aztec 2.3 energy-dispersive X-ray (EDX) spectrometer (Oxford, UK) was attached to the SEM for compositional analyses.

D. Electrical measurements

Both surfaces of polished ceramic disks (10 mm in diameter and 0.8 mm in thickness) were sputtered with Au and silver paste to form electrodes for electric measurements. Temperature dependences of the electrical conductivity (σ), the dielectric permittivity (ϵ'), and the dielectric loss (tan δ) at a frequency of 1 kHz were measured from -75 to 200 °C at a heating rate of 2 °C min⁻¹ using a Concept 41 Dielectric/Impedance spectrometer (Novocontrol) with an applied voltage of 1 V. Frequency dependences of σ , ϵ' , and tan δ were measured at RT. The accuracy in measurements of ϵ' , tan δ , and temperature control is less than 5%, 3 × 10^{-5} , and ± 0.3 °C, respectively.

E. Electron paramagnetic resonance

EPR spectra were measured using an A300 electron-spin resonance spectrometer system (Bruker BioSpin GMBH) at an X-band frequency of 9.148 GHz. The EPR of six samples, (1) 1 mg of BaTbO₃ + 99 mg of BaCeO₃, (2) 1 mg of SrTbO₃ + 99 mg of SrCeO₃, (3) 5 mg of BSTC1 + 95 mg of BaCeO₃,

(4) 5 mg of BSTC3/2 + 95 mg of BaCeO₃, (5) 100 mg of BaCeO₃, and (6) 100 mg of SrCeO₃, were measured at RT. The *g*-factor is calculated by the formula $hv_0 = g\beta H$, where the Planck's constant is $h = 6.0626 \times 10^{-34}$ J s, v_0 is the frequency, $\beta = 9.262 \times 10^{-24}$ J T⁻¹, and *H* is the magnetic field strength.

F. X-ray photoelectron spectroscopy

XPS measurements were performed at RT using an Escalab 250 Xi X-ray photoelectron spectrometer (Thermo Electron). XPS raw data were processed by smoothing multiple times. The core-level binding energy was calibrated using the C 1s peak located at 285 eV.

III. RESULTS AND DISCUSSION

A. Structure determination

For X-ray powder diffraction (XRPD) measurements, two factors are considered: (1) in general cases, the high-angle upper-limit $2\theta_{\rm m}$ is taken to be 120° or larger, which is apt to yield satisfactory refined results in the crystal structure refinement (Chen et al., 1995) and (2) the possible ordering of Ba^{2+}/Sr^{2+} at the A-sites or Ce/Tb at the B-sites in the lattice will result in extra peaks, forming a so-called superstructure. The superstructure peaks with higher intensity are easier to be observed in the lower 2θ range of $10-20^\circ$, such as for A-site-ordered perovskites (Li_{1/2}Nd_{1/2})TiO₃ (Takahashi et al., 1991) and (La_{1/2}Na_{1/2})TiO₃ (Ioshiyuki et al., 1992), as well as B-site-ordered perovskites Pb₂CaTeO₆ (Artner and Weil, 2019), Ba₂CaMoO₆ (Nguyen et al., 2019), and $Ba_{2-x}La_xFeMoO_6$ (Hussain *et al.*, 2019). For these two factors, the 2θ ranges in XRPD spectra of BSTC1 and BSTC3/ 2 are expanded from 5° to 120° and no diffraction peaks between 5° and 20° were observed. Their XRPD patterns (black line) show perovskite-like diffraction feature, as shown in Figure 1. On the basis of the trigonal $BaCeO_3$ (88592-ICSD) and the orthorhombic SrCeO₃ (PDF# 01-074-8250), the XRPD patterns of BSTC1 and BSTC3/2 can be refined by the Rietveld method, and their crystalline structures are indexed as a trigonal perovskite structure with the space group R-3c and an orthorhombic perovskite structure with the space group *Pmcn*, respectively. Table I gives the details of Rietveld refinements including lattice parameters, cell volumes, crystal system, space group, reliability factors $(R_{\rm p} \text{ and } R_{\rm wp})$, and Uiso parameters. Table II gives positional parameters and occupancy at RT. The general formulas of BSTC1 and BSTC3/2 are expressed by Ba₂SrTb₂CeO₉ and $Ba_{1,5}Sr_{1,5}Tb_{1,5}Ce_{1,5}O_{9}$, respectively. Pure phases were obtained for these two compositions, as no secondary crystalline phase was detected in the measured XRD patterns.

Under the current measuring conditions, no superstructure peaks were detected for BSTC. However, the superstructure may be imperceptible based on the XRPD, and the potential ordering in the structure is possibly undetectable by XRD. A long-time record of diffraction data benefits observations of superstructure peaks (Chen *et al.*, 1997). In order to investigate the possible ordering of Ba²⁺/Sr²⁺ at the A-sites or Ce/ Tb at the B-sites in BSTC, the XRPD spectra of both samples were measured between $5^{\circ} \le 2\theta \le 20^{\circ}$ in smaller 2θ intervals (0.01°) with a slow counting of 6 s step⁻¹ (not presented here).



Figure 1. (Colour online) Final Rietveld plots for (a) BSTC1 and (b) BSTC3/2 ceramics. The upper patterns illustrate the observed data (black line) and the calculated pattern (red line). The vertical green markers show the calculated positions of Bragg reflections. The lower blue curve is the difference diagram.

Extra peaks were not observed. Thus, Ba^{2+}/Sr^{2+} at the A-sites or Ce/Tb at the B-sites in the perovskite lattice are considered to be randomly disordered (Chen *et al.*, 1989, 1995; Ioshiyuki

TABLE I. Details of Rietveld refinements for $Ba_2SrTb_2CeO_9$ and $Ba_{1.5}Sr_{1.5}Ce_{1.5}Tb_{1.5}O_9$ perovskites at RT.

Composition	Ba2SrTb2CeO9	Ba _{1.5} Sr _{1.5} Ce _{1.5} Tb _{1.5} O ₉	
Wavelength (Å)	$\lambda_1 = 1.54056 (CuK\alpha_1)$	$\lambda_1 = 1.54056 (Cu K \alpha_1)$	
	$\lambda_2 = 1.54439 (Cu K \alpha_2)$	$\lambda_2 = 1.54439 (Cu K \alpha_2)$	
2θ step scan increment (°)	0.02	0.02	
2θ range (°)	10-120	10-120	
Profile function	Pseudo-Voigt	Pseudo-Voigt	
Crystal system	Trigonal	Orthorhombic	
Space group	R-3c	Pmcn	
a (Å)	6.0531(3)	8.5665(6)	
b (Å)	6.0531	6.0651(6)	
<i>c</i> (Å)	6.0531	6.1016(6)	
V_0 (Å ³)	157.86(1)	317.02(2)	
α	60.2918(28)	90.00	
β	60.2918	90.00	
γ	60.2918	90.00	
Rp	9.92%	7.78%	
R _{wp}	11.99%	9.69%	
χ^2	2.304	1.847	
Ζ	6	12	
Uiso(Ba)	0.01306	0.00168	
Uiso(Sr)	0.01306	0.00168	
Uiso(Tb)	0.01204	0.00186	
Uiso(Ce)	0.01204	0.00186	
Uiso(O)	0.03262	0.01326	
		0.01688	

et al., 1992; Sahoo *et al.*, 2016); their ordering is very weak if existed (Chen *et al.*, 1997). BSTC1 and BSTC3/2 can be thought to possess disordered perovskite structures.

The four ceramic phases prepared according to the same conditions as BSTC for structural and valence state analyses. The observed and simulated XRPD patterns of BaCeO₃, BaTbO₃, SrCeO₃, and SrTbO₃ are shown in Figure 2. BaCeO₃, SrCeO₃, and SrTbO₃ exhibit orthorhombic perovskite structures, corresponding to PDF# 01-070-6741, PDF# 01-074-8250, and PDF# 01-089-5513, respectively, whereas BaTbO₃ has a tetragonal structure (PDF# 01-074-4289).

B. Microstructure and evidence of single phase

The BSE and SEM images, as well as EDX spectra for BSTC1 and BSTC3/2, are shown in Figure 3. Both samples show nonuniform microstructures with a grain size distribution from 0.2 to 3 μ m, but they are denser. Their relative density (ρ_r), which is referred to as the ratio of the volumetric mass density to the theoretical density (Lu *et al.*, 2019), was determined to be 93%. This reveals that dual doping with Sr and Tb in BaCeO₃ is apt to the densification of ceramics.

On the basis of BSE investigations, no difference in brightness was observed for all of the grains, suggesting the absence of secondary phases in BSTC. To further clarify the single-phase nature of BSTC, EDX investigations were made to provide evidence for compositional distributions in different grains and grain boundaries. For BSTC1, the ratios of Ba to Sr and Tb to Ce at a fine grain, a triple-grain boundary, and a coarse grain are 2.0–2.1 and 1.8–2.0, respectively, which are close to the theoretical values of Ba/Sr = Tb/Ce = 2. Similarly, the ratios of Ba/Sr = 0.9-1.0 and Tb/Ce = 0.8-1.0for BSTC3/2 are also close to the theoretical value of Ba/Sr = Tb/Ce = 1. These EDXS results reveal the homogeneous concentration distributions of Ba, Sr, Tb, and Ce in BSTC. That is to say, Ba/Sr and Tb/Ce are completely incorporated into the A-sites and the B-sites in the BSTC perovskite lattice, respectively, and no secondary phase is present; both BSTC1 and BSTC3/2 are single-phase solutions. Figure 4 gives the schematic diagram of the crystal structure for BSTC1 as a representative.

C. Electrical properties

Figure 5 shows a plot of the electrical conductivity (σ) as a function of temperature (*T*) at 1 kHz for BSTC1 and BSTC3/2. The σ increases rapidly with increasing *T*. When *T*=197 °C, the σ value of BSTC1 (σ =5.8×10⁻⁶ S cm⁻¹) is slightly greater than that of BSTC3/2 (σ =4.5×10⁻⁶ S cm⁻¹). This reveals that the difference in *x* has only a minor effect on the electrical conductivity of both samples. The σ at 25 °C increases slowly with increasing frequency (*f*), as shown in Figure 5 inset, which suggests that hole conduction is predominant, rather than electron conduction.

Figure 6 shows the temperature dependences of the dielectric permittivity (ϵ') at 1 kHz and the dielectric loss (tan δ) at RT for BSTC1 and BSTC3/2. Both ϵ' and tan δ at 1 kHz increase rapidly with increasing *T* because of high electrical conducting behavior. This further confirms that BSTC1 and BSTC3/2 are semiconductors.

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TABLE II. Positional parameters and occupancy for Ba₂SrTb₂CeO₉ (BSTC1) and Ba_{1.5}Sr_{1.5}Ce_{1.5}Tb_{1.5}O₉ (BSTC3/2) perovskites at RT.

Atom	x	у	Z	Occupancy	Site symmetry	Ζ
BSTC1						6
Ba	0.25	0.25	0.25	0.6667	2	
Sr	0.25	0.25	0.25	0.3333	2	
Ce	0	0	0	0.3333	2	
Tb	0	0	0	0.6667	2	
01	0.8134	0.6866	0.25	1	6	
BSTC3/2						
Ba/Sr	0.25	0.0143	-0.0338	0.5	4	12
Ce/Tb	0	0.5	0	0.5	4	
01	0.25	0.549	0.0029	1	4	
02	-0.0425	0.7028	0.3007	1	8	

D. EPR investigations

The EPR technique can gather an insight into the valence state of rare-earth ions and vacancies in ceramics (Lu, 2015; Lu *et al.*, 2016a). The EPR spectra of BaCeO₃, SrCeO₃, BaTbO₃, SrTbO₃, BSTC1, and BSTC3/2 at RT are shown Figure 7. Ce³⁺ (4 f^{1}) Kramers ion in compounds is EPR silent at RT because of its short spin-lattice relaxation time. Ce⁴⁺



Figure 2. (Colour online) Observed and simulated XRPD patterns of (a) $BaCeO_3$, (b) $BaTbO_3$, (c) $SrCeO_3$, and (d) $SrTbO_3$.

 $(4f^{0})$ non-Kramers ion is also EPR silent in theory. Hence, no EPR signal was observed for BaCeO₃ and SrCeO₃.

Our experiments clarified that the EPR cannot be observed for 100 mg of samples containing Tb because of the strong EPR from Tb⁴⁺. For this reason, 1 mg of BaTbO₃ or SrTbO₃ and 5 mg of BSTC1 or BSTC3/2 are dispersed into 99 and 95 mg of BaCeO₃ or SrCeO₃ for EPR observations. A very broad singlet signal at g = 2.020-2.029 appears in the four ceramics BaTbO₃, SrTbO₃, BSTC1, and BSTC3/2. This signal originates from Tb⁴⁺ (4 f^7) Kramers ions because there is no EPR response for Ba²⁺, Sr²⁺, Ce⁴⁺, and O²⁻. For 5% Tb-doped BaTiO₃ ceramics, however, Tb⁴⁺-related signal appears at $g = \sim 6.5$ (Lu, 2015; Lu *et al.*, 2016b). This reveals that when TbO₆ octahedrons act as the perovskite skeleton, the Tb⁴⁺-related EPR signal exhibits different *g* values. Thus, the EPR investigations provide evidence of the existence of a large number of Tb⁴⁺ ions in BSTC1 and BSTC3/2.

For doped BaTiO₃ ceramics, oxygen vacancies (V_O) can be detected in the low-temperature rhombohedral phase of $T \le -100$ °C (Lu *et al.*, 2016c). It is reported that SrCeO₃ may remain orthorhombic, the space group *Pbnm* from 1.2 K up to the 1 atm melting point of 2266 K. (Knight *et al.*, 2015). Our experiments confirm that no additional V_O-related EPR signal was observed for BSTC1 and BSTC3/2 (not presented here) because no phase transition was observed in these two ceramics.

E. XPS investigations

XPS spectra of Ce 3d and Tb 3d core levels of BSTC1, BSTC3/2, SrCeO₃, BaCeO₃, SrTbO₃, and Tb₄O₇ are shown in Figure 8. The Ce^{3+} and Ce^{4+} species in compounds can be differentiated by XPS, with distinct line shapes corresponding to various final states: $Ce(III) = v_0 + v' + u_0 + u'$ and Ce (IV) = v + v'' + v''' + u + u'' + u'''. The u''' component, as a satellite peak, is a fingerprint of Ce⁴⁺ state, which arises from the so-called shake-down effect (Schneider et al., 1981; Bêche et al., 2008; Jaiswal et al., 2016; Xiong et al., 2016). It can be seen from Figure 8(a) that the Ce $3d_{3/2}$ and Ce $3d_{5/2}$ 2 photoelectron lines of all Ce-containing samples exhibit the characteristic features of tetravalent cerium, in good agreement with the reports from other authors (Braaten et al., 1989; Douillard et al., 1994; Bêche et al., 2008). The final states corresponding to core-level binding energy are as follows: Ce $3d^{9}4f^{0} \text{ O } 2p^{6} \text{ final state: } u''' = 916.3 \pm 0.1 \text{ eV}, v''' = 897.9 \pm 0.1 \text{ eV}$ 0.1 eV; Ce $3d^{9}4f^{2}$ O $2p^{4}$ final state: $u = 900.4 \pm 0.2$ eV, v =882.4 ± 0.2 eV; Ce $3d^{9}4f^{1}$ O $2p^{5}$ final state: $u'' = 907.0 \pm$

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Figure 3. (Colour online) BSE and SEM images of the polished and thermally etched surfaces as well as corresponding EDX investigations at a fine grain, a triple-grain boundary junction, and a coarse grain for (a) BSTC1 and (b) BSTC3/2 ceramics.

0.3 eV, $v'' = 888.1 \pm 0.1$ eV; The Ce³⁺ spectrum (Blanco *et al.*, 2002; Bêche et al., 2008; Jaiswal *et al.*, 2016; Xiong *et al.*, 2016; Liang, 2019) is absent for BSTC. Thus, Ce is considered to be present in the form of Ce⁴⁺ and Ce³⁺ is negligible in BSTC1 and BSTC3/2.

For terbium ions in compounds, the most intense photoemission peak corresponds to the 3*d* core level, peaking at ~1277 eV for Tb $3d_{3/2}$ and ~1242 eV for Tb $3d_{5/2}$, respectively (Van Den Bossche *et al.*, 1994; Cao *et al.*, 2010; Blanco *et al.*, 2002). Tb³⁺ and Tb⁴⁺ due to their 3*d*-spectrum overlapping cannot be clearly identified. However, the corelevel spectrum of Tb⁴⁺ was reported to show relatively intense 3*d* satellites with binding energy ca. 10 eV higher than those of the main peaks (Van Den Bossche *et al.*, 1994; Martínez-Arias *et al.*, 2005; Blanco *et al.*, 2002), but no 3*d* satellite of Tb³⁺ is present, as observed for TbF₃ with Tb³⁺ (Van Den Bossche *et al.*, 1994).

All XPS spectra of BSTC1, BSTC3/2, SrTbO₃, and Tb₄O₇ show two Tb 3*d* binding energy peaks at 1276.7 \pm 0.1 and 1242.1 \pm 0.1 eV [Figure 8(b)], as mentioned above. A clear satellite at 1253.3 \pm 0.1 eV confirms the existence of Tb⁴⁺ in these ceramics. The intensity of this satellite for



Figure 4. (Colour online) A schematic diagram of the trigonal crystal structure for BSTC1. Red balls stand for O^{2-} ions. Ba²⁺ and Sr²⁺ ions are represented by green/blue balls, whose surface consists of 2/3 green and 1/3 blue, representing the molar ratio of Ba²⁺/Sr²⁺ = 2:1. Tb and Ce ions by purple/yellow balls, and the corresponding colors reflect the molar ratio of Tb/Ce = 2:1.

SrTbO₃ is obviously higher than the other three samples, which arises from the following two evidences: (1) the molar concentration of Tb in SrTbO₃ is highest and (2) the high electrical conductivity (Figure 5) should be accompanied by oxygen vacancies, which are caused by some Tb³⁺ ions in BSTC.

F. Discussion on the structure, valence states of Tb and Ce ions, and defect chemistry

For BSTC1 and BSTC3/2, the EPR and XPS results provide evidence of the existence of Tb⁴⁺ (Figures 7 and 8), whereas the XPS results indirectly confirm the existence of some Tb³⁺ ions at the B-sites (Figure 8). Table III gives ionic radii versus coordinate number (CN) (Shannon, 1976; Lu, 2015). For perovskite, the reduplicative orientations by BO₆ octahedron form the skeleton of the perovskite lattice and the A-site ions locate on the interstitial space of BO₆ skeleton. More stable Ce⁴⁺ at the B-sites in BSTC cannot be reduced to Ce³⁺ because of the existence of Tb⁴⁺ ions. The unit-cell volume (V₀) of the orthorhombic BSTC3/2 (V₀ = 317.02 Å³) is nearly the same as that of the orthorhombic SrCeO₃ (V₀ = 317.47 Å³) reported (Ranlov *et al.*, 1995). On the basis of the orthorhombic SrCeO₃ lattice, dual doping



Figure 5. (Colour online) Plot of the electrical conductivity (σ) as a function of temperature for (a) BSTC1 and (b) BSTC3/2 ceramics, measured at 1 kHz. The inset depicts σ versus *f* at 25 °C.



Figure 6. (Colour online) Temperature dependence of ϵ' at 1 kHz for BSTC1 and BSTC3/2. The inset in (a) depicts tan δ versus *T*.

with Ba and Tb results in only a little change in V_0 . This is because the expansion in V_0 caused by the A-site Ba²⁺ ions is close to the contraction in V_0 caused by the B-site Tb⁴⁺ ions on the basis of ionic size comparisons between Ba²⁺ and Sr²⁺ as well as between Tb⁴⁺ and Ce⁴⁺. The V_0 of the orthorhombic BSTC3/2 ($V_0 = 317.02 \text{ Å}^3$) is evidently greater than the volume of two unit cells of the trigonal BSTC1 ($2V_0 = 315.72 \text{ Å}^3$) (see Table I), which arises from the higher Ce concentration at the B-sites in BSTC3/2.

The cation and vacancy defects in BSTC based on the SrCeO₃ lattice can be expressed by $Ba_{Sr}^{\times}, Sr_{Sr}^{\times}, Ce_{Ce}^{\times}, Tb_{Ce}^{\times}, Tb_{Ce}^{\prime}, Tb_{Ce}^{\prime}, and O$ vacancies ($V_{0}^{\bullet\bullet}$). Tb ions at the B-sites in perovskites are confirmed to coexist in the form of the mixed-valence states of Tb³⁺/Tb⁴⁺ (Lu, 2015; Lu and Peng, 2016; Lu *et al.*, 2016a). The higher values in the electrical conductivity (Figure 5) and the dielectric loss (Figure 6) arise from the creation of oxygen vacancies, which are compensated by some Tb³⁺ ions at the B-sites based on a general



Figure 7. (Colour online) EPR spectra of BaCeO₃, SrCeO₃, BaTbO₃, SrTbO₃, BSTC1, and BSTC3/2.



Figure 8. (Colour online) XPS spectra corresponding to (a) Ce 3d core levels of BSTC1, BSTC3/2, SrCeO3, BaCeO3 and (b) Tb 3d core levels of BSTC1, BSTC3/2, SrTbO3, Tb4O7.

lattice electroneutrality condition. The sixfold coordinated Tb³⁺ can coexist with Tb^{4+} and Ce^{4+} because of a complementary ionic radii relation $2R_{Ce(IV)} \approx R_{Tb(IV)} + R_{Tb(III)}$ (Table III).

$$4BaO + Tb_4O_7 \rightarrow 4Ba_{Ba}^{\times} + 2Tb_{Ce}^{\times} + 2Tb_{Ce}^{\prime} + 11O_0$$
$$+ V_0^{\bullet\bullet}$$
(1)

The tolerance factor (t) of ABO₃-type perovskite is expressed by the following equation:

$$t = \frac{r_{\rm A} + r_{\rm O}}{\sqrt{2}(r_{\rm B} + r_{\rm O})} \tag{2}$$

where r_A , r_B , and r_O are the ionic radii of the cation A, B, and the anion O^{2-} , respectively. The ideal perovskite without distortion should show t = 1.0, and the t value of most perovskitetype oxides is empirically in the range of 0.75-1.0 (Fu et al., 2014). For $Ba_{1-x/3}Sr_{x/3}Tb_{1-x/3}Ce_{x/3}O_3$ (BSTC), the tolerance factor may be approximately expressed by the following equation:

$$t = \frac{\left[(1 - x/3)r_{\text{Ba}} + x/3 \cdot r_{\text{Sr}}\right] + r_{\text{O}}}{\sqrt{2}\left\{\left[(1 - x/3)r_{\text{Tb}} + x/3 \cdot r_{\text{Ce}}\right] + r_{\text{O}}\right\}}$$
(3)

The tolerance factors of BSTC1 and BSTC3/2 are determined to be t = 0.956 and 0.939, respectively. These high t values show that the two BSTC phases are both nearly free of distortion of the perovskite structure.

TABLE III. Ionic radii versus coordinate number (CN).

Ions	CN	<i>r</i> (Å)	
Ba ²⁺	12	1.61	
Sr ²⁺	12	1.44	
Tb ³⁺	12	1.203	
Tb ³⁺	6	0.923	
Tb ⁴⁺	6	0.76	
Ce ⁴⁺	6	0.87	
Ce ³⁺	12	1.34	

IV. CONCLUSIONS

Two ceramics $Ba_{1-x/3}Sr_{x/3}Tb_{1-x/3}Ce_{x/3}O_3$ (x = 1 and 1.5) (BSTC) with dense microstructures ($\rho_r = 93\%$, GS = 0.2–3 μ m) were prepared at 1400 °C in air using the mixed-oxides method. Dual doping with Sr and Tb in the BaCeO₃ perovskite results in a trigonal structure for BSTC1 and an orthorhombic perovskite structure for BSTC3/2. Ba/Sr ions at the A-sites and Ce/Tb ions at the B-sites in the perovskite lattice are randomly disordered and exhibit homogeneous concentration distributions. The EPR, XPS, and electrical results confirm that Ce and Tb ions in BSTC exist as Ce⁴⁺ and the mixed-valence state of Tb⁴⁺/Tb³⁺, respectively. At RT, two BSTC ceramics exhibit a similar semiconducting behavior, and the difference in x had little effect on their σ -T relations. The electrical conductivity (σ) of BSTC with x = 1.5 increases rapidly from 5.7×10^{-9} to 4.5×10^{-6} S cm⁻¹ with increasing T from -50 to 200 °C, while the σ at RT increases slowly with increasing frequency, from 3.0×10^{-8} at 1 Hz to 2.9×10^{-7} S cm⁻¹ at 10⁶ Hz. Direct evidence of Tb³⁺, oxygen vacancies, and possible local ordering still need to be further explored.

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