

Nonlocal problems in perforated domains

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In this paper, we analyse nonlocal equations in perforated domains. We consider nonlocal problems of the form $f(x) = \int_B J(x-y)(u(y) - u(x))dy$ with x in a perforated domain $\Omega^\epsilon \subset \Omega$. Here J is a nonsingular kernel. We think about Ω^ϵ as a fixed set Ω from where we have removed a subset that we call the holes. We deal both with the Neumann and Dirichlet conditions in the holes and assume a Dirichlet condition outside Ω . In the latter case we impose that u vanishes in the holes but integrate in the whole \mathbb{R}^N ($B = \mathbb{R}^N$) and in the former we just consider integrals in \mathbb{R}^N minus the holes ($B = \mathbb{R}^N \setminus (\Omega \setminus \Omega^\epsilon)$). Assuming weak convergence of the holes, specifically, under the assumption that the characteristic function of Ω^ϵ has a weak limit, $\chi_\epsilon \rightharpoonup \mathcal{X}$ weakly* in $L^\infty(\Omega)$, we analyse the limit as $\epsilon \rightarrow 0$ of the solutions to the nonlocal problems proving that there is a nonlocal limit problem. In the case in which the holes are periodically removed balls, we obtain that the critical radius is of the order of the size of the typical cell (that gives the period). In addition, in this periodic case, we also study the behaviour of these nonlocal problems when we rescale the kernel in order to approximate local PDE problems.

Keywords: perforated domains; nonlocal equations; Neumann problem; Dirichlet problem

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1. Introduction

Let $\Omega^\epsilon \subset \mathbb{R}^N$ be a family of open bounded sets satisfying

$$\Omega^\epsilon \subset \Omega$$

for some fixed open bounded domain $\Omega \subset \mathbb{R}^N$ and $\epsilon > 0$. If $\chi_\epsilon \in L^\infty(\mathbb{R}^N)$ is the characteristic function of Ω^ϵ , we also assume that there exists $\mathcal{X} \in L^\infty(\mathbb{R}^N)$ such that

$$\chi_\epsilon \rightharpoonup \mathcal{X} \text{ weakly* in } L^\infty(\Omega). \quad (1.1)$$

This means,

$$\int_\Omega \chi_\epsilon(x) \varphi(x) dx \rightarrow \int_\Omega \mathcal{X}(x) \varphi(x) dx \quad \text{as } \epsilon \rightarrow 0$$

for all $\varphi \in L^1(\Omega)$. Note that both functions χ_ϵ and \mathcal{X} satisfy

$$0 \leq \mathcal{Y}(x) \leq 1, \text{ for all } x \in \mathbb{R}^N, \text{ and } \mathcal{Y}(x) \equiv 0 \text{ as } x \in \mathbb{R}^N \setminus \Omega,$$

for $\mathcal{Y} = \chi_\epsilon$ or \mathcal{X} .

Our main goal in this paper is to study nonlocal problems with nonsingular kernels in the *perforated domains* Ω^ϵ . We consider problems of the form

$$f(x) = \int_B J(x - y)(u^\epsilon(y) - u^\epsilon(x))dy$$

with $x \in \Omega^\epsilon \subset \Omega$. Here J is a nonsingular kernel. We deal both with the Neumann and Dirichlet problems. For the Dirichlet case we impose that u vanishes in $\mathbb{R}^N \setminus \Omega^\epsilon$ and we integrate in the whole \mathbb{R}^N ($B = \mathbb{R}^N$) while in the Neumann case we just consider integrals in \mathbb{R}^N minus $\Omega \setminus \Omega^\epsilon$ ($B = \mathbb{R}^N \setminus (\Omega \setminus \Omega^\epsilon)$) only assuming that u vanishes in $\mathbb{R}^N \setminus \Omega$. Note that for this last case we have considered nonlocal Neumann boundary conditions in the holes and a Dirichlet boundary condition in the exterior of the set Ω .

Along the whole paper, we assume that the function J that appears as the kernel in the nonlocal problem satisfies the following hypotheses:

$J \in \mathcal{C}(\mathbb{R}^N, \mathbb{R})$ is nonnegative and compactly supported with $J(0) > 0$,
(H_J)

$$J(-x) = J(x) \text{ for every } x \in \mathbb{R}^N, \text{ and } \int_{\mathbb{R}^N} J(x) dx = 1.$$

On the other hand, we only assume that $f \in L^2(\Omega)$.

Our main result, that holds both for the Dirichlet and the Neumann problem says that there exists a limit as $\epsilon \rightarrow 0$,

$$\tilde{u}^\epsilon \rightharpoonup u^*, \text{ weakly in } L^2(\Omega),$$

where $\tilde{\cdot}$ denotes the extension by zero of functions defined in subsets of \mathbb{R}^N . We characterize the nonlocal problem that verifies the weak limit u^* . Note that in the Dirichlet problem, the extension by zero and the solutions coincide, and then, to consider \tilde{u} can be omitted.

For the Dirichlet problem we have the following result:

THEOREM 1.1. *Let $\{u^\epsilon\}_{\epsilon>0}$ be the family of solutions of the nonlocal Dirichlet problem*

$$f(x) = \int_{\mathbb{R}^N} J(x - y)(u^\epsilon(y) - u^\epsilon(x))dy, \quad x \in \Omega^\epsilon \tag{1.2}$$

with

$$u^\epsilon(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega^\epsilon, \tag{1.3}$$

for $f \in L^2(\Omega)$, and assume that the characteristic functions χ_ϵ satisfy (1.1). Then, there exists $u^* \in L^2(\Omega)$ such that

$$\tilde{u}^\epsilon \rightharpoonup u^* \text{ weakly in } L^2(\Omega).$$

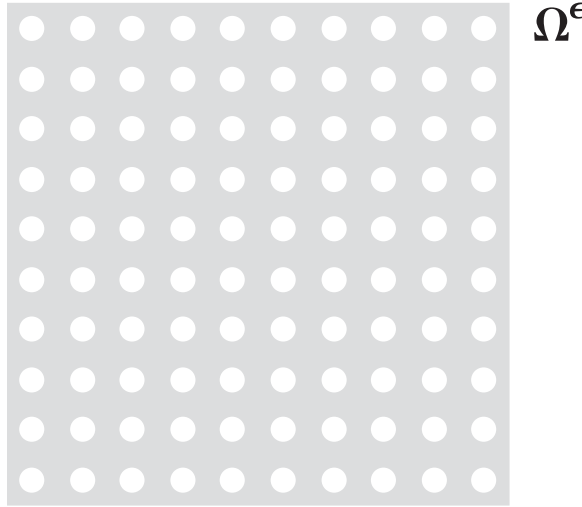


Figure 1. A periodic perforated domain $\Omega^\epsilon = (0, 1)^2 \setminus \cup B_{r^\epsilon}(x_i)$.

Moreover, the limit u^* satisfies the following nonlocal problem in Ω ,

$$\mathcal{X}(x)f(x) = \mathcal{X}(x) \int_{\mathbb{R}^N} J(x - y)(u^*(y) - u^*(x)) \, dy - (1 - \mathcal{X}(x)) u^*(x)$$

with

$$u^*(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega.$$

We want to remark that no regularity assumptions on the sets Ω^ϵ (besides measurability and the weak convergence (1.1)) is needed for our arguments.

For local operators, like the usual Laplacian, that is, for the problem $\Delta v^\epsilon = f$ in Ω^ϵ with $v^\epsilon = 0$ on $\partial\Omega^\epsilon$, the study of the behaviour of solutions in perforated domains has attracted much interest since the pioneering works [20, 28, 34]. In the classical paper [12], for example, the authors consider the Dirichlet problem for the equation $\Delta v^\epsilon = f$ in a bounded domain from where we have removed a big number of periodic small balls (the holes). That is, they consider

$$\Omega^\epsilon = \Omega \setminus \cup_i B_{r^\epsilon}(x_i)$$

where $B_{r^\epsilon}(x_i)$ is a ball centred in $x_i \in \Omega$ of the form $x_i \in 2\epsilon\mathbb{Z}^N$ with radius $0 < r^\epsilon < \epsilon \leq 1$. See figure 1 below for an example of a periodic perforated domain Ω^ϵ .

In [12] it is shown that there is a critical size of the holes (that is, a critical order of r^ϵ in ϵ) such that

$$v^\epsilon \rightarrow v^*, \quad \text{as } \epsilon \rightarrow 0,$$

with v^* given by

$$v^* = \begin{cases} \text{the solution to } \Delta v^* = f, & \text{if } r^\epsilon \ll a^\epsilon, \\ \text{the solution to } \Delta v^* - \mu v^* = f, & \text{if } r^\epsilon = a^\epsilon, \\ v^* = 0, & \text{if } r^\epsilon \gg a^\epsilon, \end{cases} \quad (1.4)$$

with Dirichlet boundary conditions $v^* = 0$ on $\partial\Omega$. Assuming $N \geq 3$, we have that the critical size of the holes is given by

$$a^\epsilon \sim \epsilon^{N/(N-2)}.$$

Note the extra term μv^* that appears in the critical case. In this particular example, for $a^\epsilon = C_0 \epsilon^{N/(N-2)}$ with C_0 a constant, μ is a positive constant that can be explicitly computed and is given by

$$\mu = \frac{S_N(N-2)}{2^N} C_0^{N-2}$$

where S_N is the surface of the sphere of radius one in \mathbb{R}^N . Also, it is proved that the convergence $v^\epsilon \rightarrow v^*$ is weak in $H^1(\Omega)$ in the first two cases and strong when $r^\epsilon \gg a^\epsilon$.

For our nonlocal problem, in the same periodic setting, we get that the critical value of r^ϵ is different from the local case and is given by

$$b^\epsilon = C_0 \epsilon$$

since in this case we obtain from theorem 1.1 that the limit u^* verifies

$$u^* = \begin{cases} \text{the solution to } \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) dy = f(x), & \text{if } r^\epsilon \ll b^\epsilon, \\ \text{the solution to } \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) dy - \nu u^*(x) = f(x), & \text{if } r^\epsilon = b^\epsilon, \\ u^* = 0, & \text{if } r^\epsilon \gg b^\epsilon. \end{cases}$$

We have weak convergence in $L^2(\Omega)$ in the first two cases and strong convergence in the last one. Here the coefficient ν that appears in the critical case is also positive and can be explicitly computed. In fact,

$$\nu = \frac{1 - \mathcal{X}}{\mathcal{X}}$$

where $\mathcal{X} \in L^\infty(\Omega)$ is just a positive constant, $\mathcal{X} = cte$, determined by the proportion of the cube which is occupied by the hole. This follows from the fact that in this periodic case we have $\chi_\epsilon \rightarrow \mathcal{X} = |Q \setminus B|/|Q|$ (here Q is the unit cube and B is a ball of radius C_0 inside the cube). In some sense, the terms \mathcal{X} and $(1 - \mathcal{X})$ in the limit problem can be seen as the effect of the holes in the original equation (1.2) and (1.3). The coefficient ν that appears in the critical case represents a kind of friction or drag caused by the perforations.

Concerning to the Neumann problem that we write as follows (see [2, 18]):

$$f(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)(u^\epsilon(y) - u^\epsilon(x)) dy, \quad x \in \Omega^\epsilon \tag{1.5}$$

with

$$u^\epsilon(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega, \tag{1.6}$$

and $f \in L^2(\Omega)$, where A^ϵ is the family of holes given by

$$A^\epsilon = \Omega \setminus \Omega^\epsilon.$$

Note that we are integrating only in $\mathbb{R}^N \setminus A^\epsilon$ in the definition of our nonlocal operator, but still assume that $u^\epsilon \equiv 0$ in $\mathbb{R}^N \setminus \Omega$. Hence, we are considering Neumann boundary conditions in the holes and a Dirichlet boundary condition outside Ω . For this problem, we need to guarantee that the quantity

$$\lambda_1^\epsilon = \inf_{u \in W_\epsilon} \frac{\frac{1}{2} \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)(u(y) - u(x))^2 dy dx}{\int_{\Omega^\epsilon} u^2(x) dx}$$

with $W_\epsilon = \{u \in L^2(\mathbb{R}^N \setminus A^\epsilon) : u(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega\}$ possesses a uniform lower bound. This is needed in order to obtain existence and uniqueness of solutions u^ϵ and is also necessary to study the asymptotic behaviour of the problem as $\epsilon \rightarrow 0$.

We have the following result:

THEOREM 1.2. *Let $\{u^\epsilon\}_{\epsilon>0}$ be a family of solutions of problem (1.5) with (1.6). Assume that the characteristic functions χ_ϵ satisfy (1.1) and that $\lambda_1^\epsilon \geq c > 0$ with c independent of ϵ . Then, there exists $u^* \in L^2(\Omega)$ such that*

$$\tilde{u}^\epsilon \rightharpoonup u^* \text{ weakly in } L^2(\Omega).$$

If $\mathcal{X} = 0$ in $L^\infty(\mathbb{R}^N)$ and

$$\int_{\mathbb{R}^N \setminus \Omega} J(x-y) dy \geq m > 0, \quad \text{for } x \in \Omega,$$

we have $u^*(x) = 0$ a.e. in \mathbb{R}^N .

If $\mathcal{X} \neq 0$ in $L^\infty(\mathbb{R}^N)$, the function u^* satisfies the following nonlocal problem in Ω ,

$$\mathcal{X}(x) f(x) = \mathcal{X}(x) \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) dy - \Lambda(x) u^*(x)$$

with

$$u^*(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega,$$

where $\Lambda \in L^\infty(\Omega)$ is given by

$$\Lambda(x) = \int_{\mathbb{R}^N} J(x-y) (1 - \chi_\Omega(y) + \mathcal{X}(y)) dy - \mathcal{X}(x), \quad x \in \Omega.$$

Here χ_Ω is the characteristic function of the open set Ω and \mathcal{X} is given by (1.1).

Observe that to deal with the Neumann problem we need to assume extra conditions (besides (1.1)) on the sets Ω^ϵ , namely, we need that $\lambda_1^\epsilon \geq c > 0$. Concerning this assumption, $\lambda_1^\epsilon \geq c > 0$, we will regard at λ_1^ϵ as the first eigenvalue of our

Neumann problem. Then, we will introduce a hypothesis involving the geometry of Ω^ϵ and the kernel J (see condition (\mathbf{H}_N) in §2) which ensures the validity of $\lambda_1^\epsilon \geq c$. We also include a simple example that shows that, in general, it could have happened that $\lambda_1^\epsilon = 0$ (in this case, we do not have existence of solutions to our nonlocal Neumann problem for a general datum f). In §4, we verify that this assumption (\mathbf{H}_N) holds in the classical case of periodic perforated domains. This fact also allows us to obtain the limit equation to the Neumann problem (1.5) with (1.6) in the case of a periodically perforated domain.

Note the more involved term Λ that appears in theorem 1.2. Rewriting it as

$$\Lambda(x) = \int_{\mathbb{R}^N \setminus \Omega} J(x-y) dy + \int_{\mathbb{R}^N} J(x-y)(\mathcal{X}(y) - \mathcal{X}(x)) dy, \quad x \in \Omega,$$

we see that the kernel J explicitly affects the extra term in the limit problem to the critical case for the Neumann problem. This dependence of the extra term on the kernel J does not occur in the Dirichlet problem where the coefficient ν only depends on the perturbation of the domain via \mathcal{X} .

We remark that the Dirichlet and Neumann problems can be written in a unified way as

$$f(x) = N(u^\epsilon)(x) + O^\epsilon(x)u^\epsilon,$$

see remark 2.3. Here N is a nonlocal operator (the same for both problems) and O^ϵ is equal to 1 in the Dirichlet case and a function of x that depends on ϵ but converges strongly in L^2 . This fact allows us to shorten a little the proofs of our main results, but we keep the presentation of both problems separate for the reader's convenience.

Many techniques and methods have been developed in order to understand the effect of the holes in perforated domains on the solutions of PDE problems with different boundary values. From pioneering works to recent ones we can still mention [1, 6, 8, 15, 17, 29, 30, 33, 36] and references therein that are concerned with elliptic and parabolic equations, nonlinear operators, as well as Stokes and Navier-Stokes equations from fluid mechanics. Note that this kind of problem is a 'homogenization' problem since the heterogeneous domain Ω^ϵ is replaced by a homogeneous one, Ω , in the limit. However, up to our knowledge, this is the first paper to deal with this kind of homogenization problem for a nonlocal operator with a nonsingular kernel.

For homogenization results for singular kernels, we refer to [10, 37, 38] (we emphasize that those references deal with homogenization in the coefficients involved in the equation and not with perforated domains as it is the case here). For random homogenization of an obstacle problem, we refer to [7]. We also remark that the case of stochastic homogenization (this is, the case in which the holes are randomly distributed inside Ω) is not treated here.

On the other hand, nonlocal equations with nonsingular kernel attracted some attention recently, see [2, 16, 21, 23, 24, 35] for a nonexhaustive list of references. We also mention [31, 32] where asymptotic problems in such nonlocal equations have been recently studied. Besides the applied models with such kernels (e.g., we refer to elasticity models, [25]), the mathematical interest is mainly due to the fact that, in general, there is no regularizing effect and therefore no general compactness tools are available.

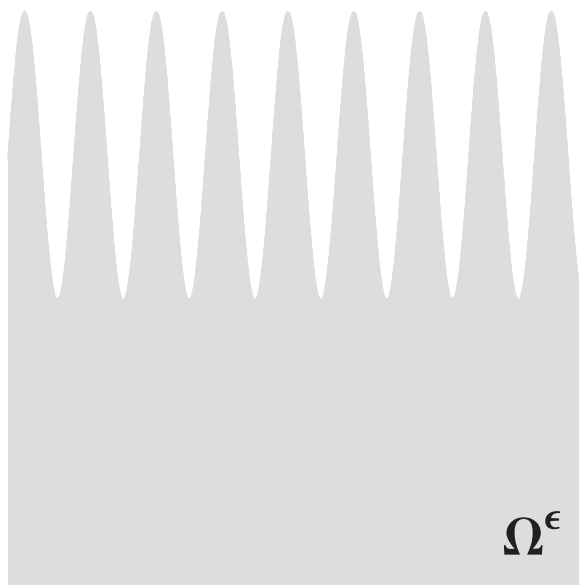


Figure 2. A domain Ω^ϵ with oscillating boundary.

Now, we comment briefly on our hypothesis and results. First, we remark that we only obtain weak convergence in L^2 of the solutions u^ϵ . This is due to the fact that the nonlocal operator does not regularize (and hence solutions u^ϵ are expected to be bounded in L^2 but nothing better) and is analogous to the fact that for the usual local case we have weak convergence in H^1 .

On the other hand, we note that our results are valid under very general assumptions on the perturbed domains. Namely, we only require that the characteristic functions of the involved domains converge weakly. Of course, this is verified in the periodic case that is our leading example (but we are not restricted to this case).

Indeed, we can allow many other different situations than perforated domains. For instance, we can consider here a family of domains Ω^ϵ whose the boundary presents a highly oscillatory behaviour as the parameter $\epsilon \rightarrow 0$. We take as a fixed domain the rectangle $\Omega = (0, 1) \times (-1, 1)$, and as a family of perturbed domains

$$\Omega^\epsilon = \{(x, y) \in \mathbb{R}^2 : x \in (0, 1), -1 < y < 0.5(1 + \sin(x/\epsilon))\}.$$

We illustrate this simple situation in figure 2 below. Note that here, the period and amplitude of the oscillations are of the same order with respect to the positive parameter ϵ . Then, if χ is the characteristic function of the open set

$$Y = \{(z, w) \in \mathbb{R}^2 : z \in (0, 2\pi), -1 < w < 0.5(1 + \sin z)\},$$

we have that the characteristic function χ_ϵ of Ω^ϵ is given by $\chi_\epsilon(x, y) = \hat{\chi}(x/\epsilon, y)$ where $\hat{\chi}$ is the periodic extension of χ with respect to its first variable (that is, in the horizontal direction). Thus, we get from Averaging theorem for oscillating

functions that

$$\chi_\epsilon \rightharpoonup \mathcal{X} = \frac{1}{2\pi} \int_0^{2\pi} \hat{\chi}(s, \cdot) ds \quad \text{weakly* in } L^\infty(\Omega)$$

as $\epsilon \rightarrow 0$. Since $\mathcal{X} \neq \chi_\Omega$, in general, we obtain from theorems 1.1 and 1.2 a nontrivial nonlocal limit problem to this case.

We quote here [3, 4, 9, 26, 27] and references therein where local problems to partial differential equations in highly oscillating domains are deeply studied.

Finally, our aim is to see how these problems behave when we introduce another parameter that controls the size of the nonlocality. In [2] (see also [18, 19]) it is shown that we can obtain solutions to local problems as limits of solutions to nonlocal problems when we rescale the kernel considering

$$J_\delta(z) = \frac{C}{\delta^{N+2}} J\left(\frac{z}{\delta}\right)$$

and letting $\delta \rightarrow 0$. Here $C = 2(\int_{\mathbb{R}^N} J(z)z_1^2)^{-1}$ is just a normalizing constant. If we apply this idea to our nonlocal problem in the Dirichlet case we are lead to consider

$$f(x) = \frac{C}{\delta^{N+2}} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{\epsilon,\delta}(y) - u^{\epsilon,\delta}(x)) dy, \quad x \in \Omega^\epsilon, \quad (1.7)$$

with $u^{\epsilon,\delta}(x) \equiv 0$, for $x \in \mathbb{R}^N \setminus \Omega^\epsilon$.

Our aim is to study the limits $\epsilon \rightarrow 0$ (to have a homogenized limit problem) and $\delta \rightarrow 0$ (to approach local problems). To perform this analysis, we restrict ourselves to the periodic case for perforated domains. That is, we consider $\Omega^\epsilon = \Omega \setminus \cup B_{r^\epsilon}(x_i)$ where $B_{r^\epsilon}(x_i)$ is a ball strictly contained in Ω centred in $x_i \in \Omega$ of the form $x_i \in 2\epsilon\mathbb{Z}^N$ with $0 < r^\epsilon < \epsilon \leq 1$.

We show that, for $u^{\epsilon,\delta}$ we have that the iterated limits

$$\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \tilde{u}^{\epsilon,\delta} = v, \quad \text{and} \quad \lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \tilde{u}^{\epsilon,\delta} = w,$$

exist, but, in general, they do not commute, that is, in general, $w \neq v$. Here, v is given by (1.4), that is,

$$v = \begin{cases} \text{the solution to } \Delta v = f, & \text{if } r^\epsilon \ll a^\epsilon, \\ \text{the solution to } \Delta v - \mu v = f, & \text{if } r^\epsilon = a^\epsilon, \\ v = 0, & \text{if } r^\epsilon \gg a^\epsilon, \end{cases}$$

with Dirichlet boundary conditions, $v = 0$ on $\partial\Omega$, and w by

$$w = \begin{cases} \text{the solution to } \Delta w = f \text{ with } w = 0 \text{ on } \partial\Omega, & \text{if } r^\epsilon \ll b^\epsilon, \\ 0, & \text{if } r^\epsilon = b^\epsilon. \end{cases} \quad (1.8)$$

Note that v is the limit for the local problem, in fact, when we compute first the limit as $\delta \rightarrow 0$ of $\tilde{u}^{\epsilon,\delta}$ we obtain a solution to a local problem with the Laplacian, then the limit $\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \tilde{u}^{\epsilon,\delta}$ coincides with the one that holds in the local case. On the other hand, when we first take the limit as $\epsilon \rightarrow 0$ from our results we

get a solution to a nonlocal problem (with a different size of the critical radius) and when we localize this problem letting $\delta \rightarrow 0$ we obtain w a solution to a local problem but different from the previous one (in general). We remark that in the limit $\lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \tilde{u}^{\epsilon, \delta}$ we have weak convergence in L^2 while for $\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \tilde{u}^{\epsilon, \delta}$ the convergence is strong in L^2 .

A similar situation (rescaling the kernel with a parameter δ in a periodically perforated domain) can be studied for the Neumann case. Now we consider

$$f(x) = \frac{C}{\delta^{N+2}} \int_{\mathbb{R}^N \setminus A^\epsilon} J\left(\frac{x-y}{\delta}\right) (u^{\epsilon, \delta}(y) - u^{\epsilon, \delta}(x)) dy, \quad x \in \Omega^\epsilon, \tag{1.9}$$

with $u^{\epsilon, \delta}(x) \equiv 0$, for $x \in \mathbb{R}^N \setminus \Omega$. In this case, we have

$$\lim_{\epsilon \rightarrow 0} P_\epsilon (\lim_{\delta \rightarrow 0} u^{\delta, \epsilon}) = v \quad \text{in } L^2(\Omega),$$

with P_ϵ an extension operator. Here the limit v is given by

$$v = \begin{cases} \text{the solution to } \Delta v = f, & \text{if } r^\epsilon \ll b^\epsilon, \\ \text{the solution to } \sum_{i,j=1}^N q_{ij} \frac{\partial^2 v}{\partial x_i \partial x_j} = \frac{|Q \setminus B|}{|Q|} f, & \text{if } r^\epsilon = b^\epsilon, \end{cases} \tag{1.10}$$

with Dirichlet boundary conditions, $v = 0$ on $\partial\Omega$. The constants q_{ij} are the *homogenized coefficients* and can be explicitly computed (see [14] and §5).

On the other hand, the limit

$$\lim_{\delta \rightarrow 0} (\lim_{\epsilon \rightarrow 0} \tilde{u}^{\delta, \epsilon}) = w$$

exists (but this time the convergence is weak in $L^2(\Omega)$), and is given by

$$w = \begin{cases} \text{the solution to } \Delta w = f, \text{ with } w = 0 \text{ on } \partial\Omega, & \text{if } r^\epsilon \ll b^\epsilon, \\ 0, & \text{if } r^\epsilon = b^\epsilon. \end{cases} \tag{1.11}$$

The paper is organized as follows: in § 2, we give some basic results needed to deal with the Dirichlet and Neumann problem. Next, in § 3, we show the main results of the paper, theorems 1.1 and 1.2. In § 4, we deal with the case of periodically distributed holes. Finally, in § 5, we rescale the kernel.

2. Basic facts and preliminary results

In this section, we discuss the existence and uniqueness of our equations writing them in an appropriated way to pass to the limit. As we will see, the first eigenvalue of the nonlocal Dirichlet and Neumann problems plays an important rule in order to do that.

2.1. The Dirichlet problem

First, we deal with the nonlocal Dirichlet problem. For $f \in L^2(\Omega)$, we consider

$$f(x) = \int_{\mathbb{R}^N} J(x - y)(u^\epsilon(y) - u^\epsilon(x))dy, \quad x \in \Omega^\epsilon \tag{2.1}$$

with

$$u^\epsilon(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega^\epsilon. \tag{2.2}$$

Observe that existence and uniqueness of our problem follows considering the variational problem

$$\min_{u \in W_\epsilon} \frac{1}{4} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x - y)(u(y) - u(x))^2 dy dx - \int_{\mathbb{R}^N} f(x)u(x) dx$$

with $W_\epsilon = \{u \in L^2(\Omega^\epsilon) : u^\epsilon(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega^\epsilon\}$. It follows from **(H_J)** that the unique minimizer (that we call u^ϵ) is a solution to (2.1) and (2.2) since it holds that

$$\begin{aligned} 0 &= -\frac{1}{2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x - y)(u^\epsilon(y) - u^\epsilon(x))(\varphi(y) - \varphi(x)) dy dx - \int_{\mathbb{R}^N} f(x)\varphi(x) dx \\ &= \int_{\mathbb{R}^N} \varphi(x) \int_{\mathbb{R}^N} J(x - y)(u^\epsilon(y) - u^\epsilon(x)) dy dx - \int_{\mathbb{R}^N} f(x)\varphi(x) dx \end{aligned} \tag{2.3}$$

for any $\varphi \in L^2(\mathbb{R}^N)$ with $\varphi(x) \equiv 0$ for $x \in \mathbb{R}^N \setminus \Omega^\epsilon$. Note that, if $u^\epsilon \in L^2(\Omega^\epsilon)$ is a solution to (2.1) and (2.2), then it is the unique minimizer. Taking $\varphi = u^\epsilon$ in (2.3) we get

$$\begin{aligned} - \int_{\Omega^\epsilon} f(x)u^\epsilon(x) dx &= \frac{1}{2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x - y)(u^\epsilon(y) - u^\epsilon(x))^2 dy dx \\ &\geq \beta_1^\epsilon \int_{\Omega^\epsilon} (u^\epsilon(x))^2 dx \end{aligned} \tag{2.4}$$

where β_1^ϵ is the first eigenvalue associated with this operator in the space W_ϵ . It is given by

$$\beta_1^\epsilon = \inf_{u \in W_\epsilon} \frac{1/2 \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x - y)(u(y) - u(x))^2 dy dx}{\int_{\Omega^\epsilon} u^2(x) dx}. \tag{2.5}$$

From [2, proposition 2.3] we know that β_1^ϵ is strictly positive. Therefore, due to (2.4), we get

$$\|u^\epsilon\|_{L^2(\Omega^\epsilon)} \leq \frac{1}{\beta_1^\epsilon} \|f\|_{L^2(\Omega)}.$$

Thus, since $\beta_1^\epsilon \geq c > 0$ with c independent of ϵ (see lemma 3.5 below), we also obtain

$$\|u^\epsilon\|_{L^2(\Omega^\epsilon)} \leq K$$

for some positive constant K depending only on f (and so, independent of ϵ). Then, along a subsequence if necessary,

$$\tilde{u}^\epsilon \rightharpoonup u^* \quad \text{weakly in } L^2(\Omega) \tag{2.6}$$

as $\epsilon \rightarrow 0$ where $\tilde{\cdot}$ denotes the extension by zero applied to functions defined in subsets of \mathbb{R}^N .

Thus, if χ_ϵ is the characteristic function of Ω^ϵ and \tilde{u}^ϵ is the extension by zero of u^ϵ to \mathbb{R}^N , we can use $\int_{\mathbb{R}^N} J(x - y) dy = 1$ to rewrite (2.3) as

$$\int_{\Omega} \chi_\epsilon(x) \varphi(x) f(x) dx = \int_{\Omega} \chi_\epsilon(x) \varphi(x) N(\tilde{u}^\epsilon)(x) dx - \int_{\Omega} \varphi(x) \tilde{u}^\epsilon(x) dx \tag{2.7}$$

for any $\varphi \in L^2(\Omega)$ where $N : L^2(\Omega) \mapsto L^2(\Omega)$ is the nonlocal operator set by

$$N(u)(x) = \int_{\Omega} J(x - y) u(y) dy, \quad x \in \Omega^\epsilon. \tag{2.8}$$

2.2. The Neumann problem

Now let us discuss the nonlocal Neumann problem. Given $f \in L^2(\Omega)$, we consider here the following equation

$$f(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y) (u^\epsilon(y) - u^\epsilon(x)) dy, \quad x \in \Omega^\epsilon \tag{2.9}$$

with

$$u^\epsilon(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega. \tag{2.10}$$

Here we are calling A^ϵ the set $A^\epsilon = \Omega \setminus \Omega^\epsilon$.

In this problem (2.9) with (2.10), we have taken nonlocal Neumann boundary conditions in the holes and a Dirichlet boundary condition in the exterior of the set Ω . The arguments that we use in this section are similar to the ones for the Dirichlet case.

Since we want to consider general functions $f \in L^2(\Omega)$, in what follows we need that the first eigenvalue associated to this problem, that is given by

$$\lambda_1^\epsilon = \inf_{u \in W_\epsilon} \frac{1/2 \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y) (u(y) - u(x))^2 dy dx}{\int_{\Omega^\epsilon} u^2(x) dx} \tag{2.11}$$

with $W_\epsilon = \{u \in L^2(\mathbb{R}^N \setminus A^\epsilon) : u(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega\}$, is strictly positive.

In fact, when $\lambda_1^\epsilon = 0$ then our problem (2.9) may not have solutions as the following example shows: Assume that $\lambda_1^\epsilon = 0$ and take f such that $\int_{\mathbb{R}^N \setminus A^\epsilon} f \phi_1 \neq 0$

being ϕ_1 an eigenfunction associated with λ_1^ϵ , that is, ϕ_1 verifies,

$$0 = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(\phi_1(y) - \phi_1(x))dy.$$

The existence of such an eigenfunction can be proved as in [22]. Then, assume that there is a solution u^ϵ to the problem

$$f(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(u^\epsilon(y) - u^\epsilon(x))dy, \quad x \in \Omega^\epsilon.$$

Multiplying by ϕ_1 and integrating in $\mathbb{R}^N \setminus A^\epsilon$ we get

$$\begin{aligned} 0 &\neq \int_{\mathbb{R}^N \setminus A^\epsilon} f(x)\phi_1(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} \phi_1(x) \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(u^\epsilon(y) - u^\epsilon(x))dy \, dx \\ &= \int_{\mathbb{R}^N \setminus A^\epsilon} u^\epsilon(x) \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(\phi_1(y) - \phi_1(x))dy \, dx = 0, \end{aligned} \tag{2.12}$$

a contradiction.

A condition under which the first eigenvalue λ_1^ϵ is strictly positive is given in our next result. Let us introduce this condition that we call **(H_N)**.

(H_N) We assume that there exists a finite family of sets $B_0, B_1, \dots, B_L \subset \mathbb{R}^N \setminus A^\epsilon$ such that $B_0 = \mathbb{R}^N \setminus \Omega$,

$$(\mathbb{R}^N \setminus A^\epsilon) \subset \bigcup_{i=0}^L B_i \quad \text{and} \quad \alpha_j = \frac{1}{4} \min_{x \in B_j} \int_{B_{j-1}} J(x - y) \, dy > 0.$$

This condition says that we can cover $\mathbb{R}^N \setminus A^\epsilon$ by a finite family of sets B_i starting with $B_0 = \mathbb{R}^N \setminus \Omega$ in such a way that every point in some B_j sees the set B_{j-1} with uniformly positive probability. This property holds if A^ϵ is not very thick (see example 3.7 at § 3).

LEMMA 2.1. *Under the previous condition **(H_N)** on the set A^ϵ , there exists a positive constant λ such that*

$$\int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)|u(y) - u(x)|^2 \, dy \, dx \geq \lambda \int_{\mathbb{R}^N \setminus A^\epsilon} |u(x)|^2 \, dx,$$

for every $u \in W_\epsilon = \{u \in L^2(\mathbb{R}^N \setminus A^\epsilon) : u(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega\}$.

Proof. Following [2, Chapter 6] we cover the domain $\mathbb{R}^N \setminus A^\epsilon$ with a finite family of disjoint sets, B_j $j = 0, 1, \dots, L$ (the existence of such sets is guaranteed by our

hypothesis) and define

$$\alpha_j = \frac{1}{4} \min_{x \in B_j} \int_{B_{j-1}} J(x-y) dy, \quad 1 = \int_{\mathbb{R}^N} J(s) ds. \tag{2.13}$$

Now, for $u \in W_\epsilon$, we have

$$\begin{aligned} & \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)|u(y) - u(x)|^2 dy dx \\ & \geq \int_{B_j} \int_{B_{j-1}} J(x-y)|u(y) - u(x)|^2 dy dx \end{aligned}$$

for $j = 1, \dots, L$, and

$$\begin{aligned} & \int_{B_j} \int_{B_{j-1}} J(x-y)|u(y) - u(x)|^2 dy dx \\ & \geq \frac{1}{4} \int_{B_j} \left(\int_{B_{j-1}} J(x-y) dy \right) |u(x)|^2 dx - \int_{B_{j-1}} \left(\int_{B_j} J(x-y) dx \right) |u(y)|^2 dy \\ & \geq \alpha_j \int_{B_j} |u(x)|^2 dx - \int_{B_{j-1}} |u(y)|^2 dy. \end{aligned}$$

Then, taking $B_0 = \mathbb{R}^N \setminus \Omega$ (notice that $u = 0$ in B_0), we can iterate this inequality to get that

$$\int_{B_j} |u(x)|^2 dx \leq C_j \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)|u(y) - u(x)|^2 dy dx$$

where

$$C_1 = \frac{1}{\alpha_1}, \quad C_j = \frac{1}{\alpha_j}(1 + C_{j-1}) \quad j = 2, \dots, L.$$

Therefore, adding in j , we have the Poincaré type inequality

$$\int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)|u(y) - u(x)|^2 dy dx \geq \lambda \int_{\mathbb{R}^N \setminus A^\epsilon} |u(x)|^2 dx, \tag{2.14}$$

with

$$\lambda = \left(\sum_{j=1}^L C_j \right)^{-1} \sim \prod_{j=1}^L \alpha_j,$$

as we wanted to show. □

REMARK 2.2. Note that if we have that for each ϵ condition (\mathbf{H}_N) is satisfied with the number of sets, L , independent of ϵ and a uniform constant α such that $\alpha_j \geq \alpha > 0$, then there is a positive constant c independent of ϵ such that $\lambda_1^\epsilon \geq c > 0$. In § 4, we will see that this property holds, for example, for periodically perforated

domains in the case that the characteristic function $\mathcal{X} \in L^\infty(\mathbb{R}^N)$ given in (1.1) is not the null function.

Now, let us observe that existence and uniqueness to (2.9) with (2.10) follows considering the variational problem

$$\min_{u \in W_\epsilon} \frac{1}{4} \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(u(y) - u(x))^2 dy dx - \int_{\mathbb{R}^N \setminus A^\epsilon} f(x)u(x) dx.$$

In fact, it follows from hypothesis **(H_N)** (we can use lemma 2.1) that

$$\|u\|^2 := \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(u(y) - u(x))^2 dy dx$$

is a norm equivalent to the usual L^2 -norm in W_ϵ . Hence, the functional involved in the minimization problem is lower semicontinuous, coercive and convex in W_ϵ , and then possesses a unique minimizer (that we call u^ϵ). As for the Dirichlet case, this minimizer is a solution to (2.9) and (2.10). Moreover, if $u^\epsilon \in L^2(\Omega^\epsilon)$ satisfies (2.9) and (2.10), then it is a minimizer. Multiplying the equation by u^ϵ we get

$$\begin{aligned} - \int_{\Omega^\epsilon} f(x)u^\epsilon(x) dx &= \frac{1}{2} \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(u^\epsilon(y) - u^\epsilon(x))^2 dy dx \\ &\geq \lambda_1^\epsilon \int_{\Omega^\epsilon} (u^\epsilon(x))^2 dx. \end{aligned} \tag{2.15}$$

Here λ_1^ϵ is the first eigenvalue (2.11) associated with this operator in the space W_ϵ . Therefore, assuming that $\lambda_1^\epsilon \geq c > 0$ with c independent of ϵ (see remark 2.2), we get that

$$\int_{\Omega^\epsilon} (u^\epsilon(x))^2 dx$$

is bounded by a constant that depends only on f but is independent of ϵ . Hence, along a subsequence if necessary,

$$\tilde{u}^\epsilon \rightharpoonup u^*$$

weakly in $L^2(\Omega)$ as $\epsilon \rightarrow 0$ where $\tilde{\cdot}$ denotes the extension by zero on functions defined in the open set Ω^ϵ . Hence, if χ_ϵ is the characteristic function of Ω^ϵ and \tilde{u}^ϵ is the extension by zero of u^ϵ to \mathbb{R}^N , we can write the weak form of the equation as

$$\int_{\Omega} \chi_\epsilon(x) \varphi(x) f(x) dx = \int_{\Omega} \chi_\epsilon(x) \varphi(x) N(\tilde{u}^\epsilon)(x) dx - \int_{\Omega} \varphi(x) \tilde{u}^\epsilon(x) O^\epsilon(x) dx \tag{2.16}$$

for any $\varphi \in L^2(\Omega)$. Here, N is the nonlocal operator introduced in (2.8), and O^ϵ is the following function defined on \mathbb{R}^N

$$O^\epsilon(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y) dy. \tag{2.17}$$

REMARK 2.3. Notice that both problems (2.1) and (2.9) with (2.10) can be written on the same form (2.16). For the Dirichlet problem, we take O^ϵ constant and equal to 1. On the other hand, for the Neumann equation, we take O^ϵ as (2.17). In this way, we can pass to the limit in our problems, proving theorems 1.1 and 1.2.

3. Proof of theorems 1.1 and 1.2.

Now we are ready to proceed with the proof of theorems 1.1 and 1.2. We first consider the nonlocal Dirichlet problem.

Proof of theorem 1.1. We need to pass to the limit in (2.7). In order to do that, we evaluate

$$N(\tilde{u}^\epsilon)(x) = \int_{\Omega} J(x - y)\tilde{u}^\epsilon(y) \, dy, \quad x \in \mathbb{R}^N. \tag{3.1}$$

From (2.6), we have that

$$N(\tilde{u}^\epsilon)(x) \rightarrow N(u^*)(x) = \int_{\Omega} J(x - y) u^*(y) \, dy, \quad \text{as } \epsilon \rightarrow 0, \tag{3.2}$$

for each $x \in \mathbb{R}^N$. Hence, since u^ϵ is uniformly bounded in $L^2(\Omega^\epsilon)$ and

$$|N(\tilde{u}^\epsilon)(x)| \leq |\Omega|^{1/2} \|J\|_\infty \|u^\epsilon\|_{L^2(\Omega^\epsilon)}, \tag{3.3}$$

we obtain by Dominated Convergence theorem that

$$\int_{\Omega} N(\tilde{u}^\epsilon)(x) \varphi(x) \, dx \rightarrow \int_{\Omega} N(u^*)(x) \varphi(x) \, dx$$

for all $\varphi \in L^2(\Omega)$, and then, $N(\tilde{u}^\epsilon) \rightharpoonup N(u^*)$ weakly in $L^2(\Omega)$. Indeed, we can prove that

$$N(\tilde{u}^\epsilon) \rightarrow N(u^*) \quad \text{strongly in } L^2(\Omega). \tag{3.4}$$

It follows from (3.2) and (3.3) that $(N(\tilde{u}^\epsilon)(x))^2 \rightarrow (N(u^*)(x))^2$ for all $x \in \mathbb{R}^N$ with

$$|N(\tilde{u}^\epsilon)(x)|^2 \leq |\Omega| \|J\|_\infty^2 \|u^\epsilon\|_{L^2(\Omega^\epsilon)}^2.$$

Consequently, using again the Dominated Convergence theorem, we get

$$\|N(\tilde{u}^\epsilon)\|_{L^2(\Omega)} \rightarrow \|N(u^*)\|_{L^2(\Omega)}$$

as $\epsilon \rightarrow 0$ proving (3.4) since we are working in Hilbert spaces.

We now can combine (1.1), (2.6), (2.7) and (3.1) to obtain

$$\int_{\Omega} \mathcal{X}(x) \varphi(x) f(x) \, dx = \int_{\Omega} \mathcal{X}(x) \varphi(x) N(u^*)(x) \, dx - \int_{\Omega} \varphi(x) u^*(x) \, dx \tag{3.5}$$

for any $\varphi \in L^2(\Omega)$. We still can rewrite (3.5) as

$$\begin{aligned} \int_{\Omega} \mathcal{X}(x) \varphi(x) f(x) \, dx &= \int_{\Omega} \mathcal{X}(x) \varphi(x) \left(\int_{\mathbb{R}^N} J(x-y) (u^*(y) - u^*(x)) \, dy \right) \, dx \\ &\quad - \int_{\Omega} \varphi(x) u^*(x) (1 - \mathcal{X}(x)) \, dx \end{aligned} \tag{3.6}$$

which implies

$$\mathcal{X}(x) f(x) = \mathcal{X}(x) \int_{\mathbb{R}^N} J(x-y) (u^*(y) - u^*(x)) \, dy - (1 - \mathcal{X}(x)) u^*(x), \quad \text{a.e. } \Omega, \tag{3.7}$$

with $u^*(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega$.

Finally, we show that u^* is the unique solution of (3.7). This fact implies the convergence of whole sequence \tilde{u}^ϵ . To do that, we consider the set D where the function \mathcal{X} vanishes

$$D = \{x \in \Omega : \mathcal{X}(x) = 0\}. \tag{3.8}$$

From (3.7), we have $u^*(x) = 0$ for all $x \in D$. Hence, if $D = \Omega$ a.e., $u^*(x) = 0$ a.e. $x \in \mathbb{R}^N$ is the unique solution. Thus, let us consider that $\Omega \setminus D$ is a nontrivial measurable set. We can rewrite equation (3.7) as

$$f(x) = \int_{\mathbb{R}^N} J(x-y) (u^*(y) - u^*(x)) \, dy - \frac{(1 - \mathcal{X}(x))}{\mathcal{X}(x)} u^*(x), \quad \text{a.e. } \Omega \setminus D, \tag{3.9}$$

with $u^*(x) \equiv 0, x \in \mathbb{R}^N \setminus (\Omega \setminus D)$. Hence, if u^* and v^* are solutions of (3.9), $w^* = u^* - v^*$ satisfies

$$0 = \int_{\mathbb{R}^N} J(x-y) (w^*(y) - w^*(x)) \, dy - \frac{(1 - \mathcal{X}(x))}{\mathcal{X}(x)} w^*(x), \quad \text{a.e. } \Omega \setminus D. \tag{3.10}$$

Consequently, $(1 - \mathcal{X})/\mathcal{X} w^* \in L^2(\Omega \setminus D)$, and then,

$$\int_{\Omega \setminus D} \frac{1 - \mathcal{X}(x)}{\mathcal{X}(x)} (w^*(x))^2 \, dx = -\frac{1}{2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x-y) (w^*(y) - w^*(x))^2 \, dy \, dx \leq 0.$$

Since $(1 - \mathcal{X}(x))/\mathcal{X}(x) \geq 0$ in $\Omega \setminus D$, we get

$$\frac{1 - \mathcal{X}(x)}{\mathcal{X}(x)} (w^*(x))^2 = 0 \quad \text{a.e. } x \in \Omega \setminus D.$$

This combined with [2, proposition 2.2] completes the proof. □

If we still assume in the hypotheses of theorem 1.1 that the characteristic functions χ_ϵ converges weakly star to $\mathcal{X}(x) = 0$ a.e. in Ω , we get that the limit function u^* must be null with strong convergence in L^2 .

COROLLARY 3.1. Under the hypotheses of theorem 1.1 with $\mathcal{X}(x) = 0$ a.e. $x \in \Omega$, we obtain $u^*(x) = 0$ a.e. $x \in \mathbb{R}^N$ with

$$\tilde{u}^\epsilon \rightarrow 0, \quad \text{strongly in } L^2(\Omega).$$

Proof. From the limit equation in theorem 1.1 is very easy to see that $u^* = 0$. We obtain strong convergence just observing that the norm of \tilde{u}^ϵ also converges to 0 as $\epsilon \rightarrow 0$. Indeed, from (2.7) with $\varphi = \tilde{u}^\epsilon$ we have

$$\begin{aligned} \|\tilde{u}^\epsilon\|_{L^2(\Omega)}^2 &= \int_{\Omega} \tilde{u}^\epsilon(x)N(\tilde{u}^\epsilon)(x) - \int_{\Omega} \tilde{u}^\epsilon(x) f(x) \, dx \rightarrow \int_{\Omega} u^*(x)N(u^*)(x) \\ &\quad - \int_{\Omega} u^*(x) f(x) \, dx = 0. \end{aligned} \quad \square$$

REMARK 3.2. Let us suppose that we are in the classic situation in Homogenization Theory in which the family of perforated domains Ω^ϵ possesses an extension operator such as

$$P_\epsilon \in \mathcal{L}(L^2(\Omega^\epsilon); L^2(\Omega)) \cap \mathcal{L}(H^1(\Omega^\epsilon); H^1(\Omega))$$

with bounded norm (independently of ϵ). Let us also assume that the data are smooth, by this we mean that $f \in H^1(\Omega)$ and $J \in C^1(\mathbb{R}^N, \mathbb{R})$. Then, it follows for solutions to the nonlocal Dirichlet problem (2.1) and (2.2) that $u^\epsilon \in H^1(\Omega^\epsilon)$ and satisfies

$$\frac{\partial u^\epsilon}{\partial x_i}(x) = \int_{\mathbb{R}^N} J(x - y) \left(\frac{\partial u^\epsilon}{\partial x_i}(y) - \frac{\partial u^\epsilon}{\partial x_i}(x) \right) \, dy - \frac{\partial f}{\partial x_i}(x), \quad \text{a.e. } \Omega^\epsilon.$$

Consequently, since $\|\partial f / \partial x_i\|_{L^2(\Omega^\epsilon)}$ is uniformly bounded, from the same arguments used in lemma 3.5, we also get $\|u^\epsilon\|_{H^1(\Omega^\epsilon)} \leq K$ for some constant $K > 0$ independent of ϵ . Hence, $P_\epsilon u^\epsilon$ is uniformly bounded in $H^1(\Omega)$, from where we can extract a subsequence (up to a sequence) such that

$$P_\epsilon u^\epsilon \rightarrow u_0 \quad \text{strongly in } L^2(\Omega), \tag{3.11}$$

for some $u_0 \in H^1(\Omega)$.

Remark that for nonlocal problems with nonsingular kernels it does not hold that the Dirichlet datum is taken continuously, that is, we do not have $u^\epsilon(x) \rightarrow 0$ as $x \rightarrow \partial\Omega^\epsilon$ (even for solutions in $C(\Omega^\epsilon)$), see [5]. Therefore, the extension $P_\epsilon u^\epsilon$ of u^ϵ to the holes does not coincide with the extension by zero, \tilde{u}^ϵ , that we have considered here.

Now, we observe that we can pass to the limit (weakly in L^2) in the identity

$$\tilde{u}^\epsilon(x) = \chi_\epsilon(x)P_\epsilon u^\epsilon(x), \quad \text{a.e. } \Omega.$$

Note that \tilde{u}^ϵ is the extension by zero of u^ϵ in the holes and that $\chi_\epsilon P_\epsilon u^\epsilon$ first extends u^ϵ to the holes and then multiply by χ_ϵ . From the limit of $\tilde{u}^\epsilon = \chi_\epsilon P_\epsilon u^\epsilon$ we obtain

the following relation between the function $u^* \in L^2(\Omega)$ given by theorem 1.1, and $u_0 \in H^1(\Omega)$ introduced in (3.11)

$$u^*(x) = \mathcal{X}(x) u_0(x), \quad \text{a.e. } \Omega.$$

For the particular case of a periodically perforated domain, we have that $\mathcal{X}(x)$ is a constant and therefore we obtain that $u^*(x) = \mathcal{X} u_0(x) \in H^1(\Omega)$. This regularity result for the limit u^* can also be obtained from the limit problem that is satisfied since we assumed that $f \in H^1(\Omega)$ and $J \in C^1(\mathbb{R}^N, \mathbb{R})$.

Now let us consider the Neumann problem (2.9) with (2.10). We show theorem 1.2.

Proof of theorem 1.2. Under the assumption $\lambda_1^\epsilon \geq c > 0$ with c independent of ϵ , we have already seen that $\|u^\epsilon\|_{L^2(\Omega^\epsilon)} \leq c^{-1} \|f\|_{L^2(\Omega)}$ from where we obtain (up to a subsequence) the weak convergence

$$\tilde{u}^\epsilon \rightharpoonup u^*.$$

First, we observe that the additional assumption $\mathcal{X}(x) = 0$ implies $u^*(x) = 0$ in \mathbb{R}^N . In fact, it is a consequence of the lower semicontinuity of the norm and the limit

$$\left| \int_{\Omega} \chi_\epsilon(x) \varphi(x) u^\epsilon(x) \, dx \right| \leq \|\chi_\epsilon \varphi\|_{L^2(\Omega^\epsilon)} \|u^\epsilon\|_{L^2(\Omega^\epsilon)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0 \quad \forall \varphi \in L^2(\Omega).$$

Now let us get the limit problem to the other cases. In order to do that, we need to pass to the limit in (2.16). First, we pass to the limit in O^ϵ given by (2.17). We show that

$$O^\epsilon \rightarrow O_0 \quad \text{strongly in } L^2(\Omega) \tag{3.12}$$

with

$$O_0(x) = \int_{\mathbb{R}^N} J(x - y)(1 - \chi_\Omega(y) + \mathcal{X}(y)) \, dy.$$

Since

$$O^\epsilon(x) = \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y) \, dy = \int_{\mathbb{R}^N} J(x - y) \, dy - \int_{\Omega} J(x - y) \, dy + \int_{\Omega^\epsilon} J(x - y) \, dy$$

for any $x \in \mathbb{R}^N$, we just need to observe that, from (1.1), we get

$$\hat{O}^\epsilon(x) = \int_{\Omega^\epsilon} J(x - y) \, dy = \int_{\Omega} \chi_\epsilon(y) J(x - y) \, dy. \rightarrow \hat{O}_0(x) = \int_{\Omega} \mathcal{X}(y) J(x - y) \, dy$$

for all $x \in \mathbb{R}^N$. Thus, since $|O^\epsilon(x)| \leq 1$ for every $\epsilon > 0$ and $x \in \mathbb{R}^N$, we can argue as in (3.4) to obtain from Dominated Convergence theorem that

$$\hat{O}^\epsilon \rightarrow \hat{O}_0 \quad \text{strongly in } L^2(\Omega).$$

Consequently, we conclude (3.12).

Now we can combine (1.1), (3.4) and (3.12) to pass to the limit in (2.16) obtaining

$$\int_{\Omega} \mathcal{X}(x) \varphi(x) f(x) \, dx = \int_{\Omega} \mathcal{X}(x) \varphi(x) N(u^*)(x) \, dx - \int_{\Omega} \varphi(x) u^*(x) \int_{\mathbb{R}^N} J(x-y)(1 - \chi_{\Omega}(y) + \mathcal{X}(y)) \, dy \, dx$$

for any $\varphi \in L^2(\Omega)$. This can be rewritten as

$$\begin{aligned} & \int_{\Omega} \mathcal{X}(x) \varphi(x) f(x) \, dx \\ &= \int_{\Omega} \mathcal{X}(x) \varphi(x) \int_{\mathbb{R}^N} J(x-y) (u^*(y) - u^*(x)) \, dy \, dx \\ &+ \int_{\Omega} \varphi(x) u^*(x) \left(\mathcal{X}(x) - \int_{\mathbb{R}^N} J(x-y)(1 - \chi_{\Omega}(y) + \mathcal{X}(y)) \, dy \right) \, dx. \end{aligned}$$

Observe that $u^*(x) \equiv 0$ wherever $x \in \mathbb{R}^N \setminus \Omega$.

Therefore, there exists $u^* \in L^2(\Omega)$ with $\tilde{u}^\epsilon \rightharpoonup u^*$ weakly in $L^2(\Omega)$ and u^* satisfying the following nonlocal problem

$$\mathcal{X}(x) f(x) = \mathcal{X}(x) \int_{\mathbb{R}^N} J(x-y) (u^*(y) - u^*(x)) \, dy - \Lambda(x) u^*(x) \tag{3.13}$$

with $u^*(x) \equiv 0, x \in \mathbb{R}^N \setminus \Omega$, where $\Lambda \in L^\infty(\Omega)$ is given by

$$\Lambda(x) = \int_{\mathbb{R}^N} J(x-y) (1 - \chi_{\Omega}(y) + \mathcal{X}(y)) \, dy - \mathcal{X}(x), \quad x \in \Omega. \tag{3.14}$$

Uniqueness for the limit problem can be obtained as in the proof of theorem 1.1. This fact implies the convergence of whole sequence \tilde{u}^ϵ . □

We still notice that a stronger hypothesis is needed to obtain strong convergence in L^2 of the extended solutions to zero for the singular case $\mathcal{X}(x) = 0$ a.e. \mathbb{R}^N . In this respect, we have the following result.

COROLLARY 3.3. *Let $\Gamma : \Omega \mapsto \mathbb{R}$ be the function given by*

$$\Gamma(x) = \int_{\mathbb{R}^N \setminus \Omega} J(x-y) \, dy. \tag{3.15}$$

Suppose $\mathcal{X}(x) = 0$ a.e. Ω in (1.1) and assume

$$\Gamma(x) \geq m > 0, \quad \text{in } \Omega \tag{3.16}$$

for some $m > 0$. Then, the solutions u^ϵ of the Neumann problem (2.9) with (2.10) satisfy

$$\|u^\epsilon\|_{L^2(\Omega^\epsilon)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0.$$

Proof. First, let us take $\varphi = \chi_\epsilon u^\epsilon$ in (2.16). Then,

$$\int_\Omega \tilde{u}^\epsilon(x) f(x) \, dx = \int_\Omega \tilde{u}^\epsilon(x) N(\tilde{u}^\epsilon)(x) \, dx - \int_\Omega \tilde{u}^\epsilon(x)^2 O^\epsilon(x) \, dx. \tag{3.17}$$

Since Γ is strictly positive in Ω , hypothesis (\mathbf{H}_N) holds, and then

$$\tilde{u}^\epsilon \rightharpoonup 0, \quad \text{weakly in } L^2(\Omega),$$

as $\epsilon \rightarrow 0$. Consequently, due to (3.17), we obtain

$$\lim_{\epsilon \rightarrow 0} \int_\Omega \tilde{u}^\epsilon(x)^2 O^\epsilon(x) \, dx = 0. \tag{3.18}$$

On the other hand, we have that

$$\left| \int_\Omega \tilde{u}^\epsilon(x)^2 O^\epsilon(x) \, dx \right| \geq \|u^\epsilon\|_{L^2(\Omega^\epsilon)} \inf_{x \in \Omega} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y) \, dy. \tag{3.19}$$

Hence, we conclude the proof from (3.16), (3.18) and (3.19). □

REMARK 3.4. Notice that we still can reach the same results described in theorems 1.1 and 1.2 for the more general situation with $f^\epsilon \rightarrow f$ strongly in $L^2(\Omega)$ since $\|f^\epsilon\|_{L^2(\Omega)}$ remains uniformly bounded in ϵ and

$$\int_\Omega \chi_\epsilon(x) f^\epsilon(x) \, dx \rightarrow \int_\Omega \mathcal{X}(x) f(x) \, dx, \quad \text{as } \epsilon \rightarrow 0.$$

Now let us show that the sequence of first eigenvalues to the Dirichlet problem converges to a positive value. Thus, this sequence possesses a positive lower bound.

LEMMA 3.5. *Let β_1^ϵ be the family of first eigenvalues introduced in (2.5). Then, there exist $\beta_1 > 0$, such that*

$$\beta_1^\epsilon \rightarrow \beta_1, \quad \text{as } \epsilon \rightarrow 0.$$

Consequently, there exist $\epsilon_0 > 0$ and $c > 0$ with

$$\beta_1^\epsilon > c > 0, \quad \forall \epsilon \in (0, \epsilon_0).$$

Proof. We start observing that β_1^ϵ is a bounded sequence with respect to ϵ . Indeed, it follows from [2, proposition 2.3] that $0 < \beta_1^\epsilon < 1$. Thus, we can extract a subsequence, still denoted by β_1^ϵ , such that $\beta_1^\epsilon \rightarrow \beta_1$. Since β_1^ϵ is strictly positive, β_1 is nonnegative.

Now let ϕ^ϵ be the associated eigenfunction to β_1^ϵ with $\|\phi^\epsilon\|_{L^2(\Omega^\epsilon)} = 1$ (the existence of such eigenfunction can be proved as in [22]). Hence,

$$-\beta_1^\epsilon \int_{\Omega^\epsilon} \phi^\epsilon(x) \varphi(x) dx = \int_{\Omega^\epsilon} \varphi(x) \int_{\mathbb{R}^N} J(x-y)(\phi^\epsilon(y) - \phi^\epsilon(x)) dy dx$$

for any $\varphi \in L^2(\Omega)$ with ϕ^ϵ vanishing in $\mathbb{R}^N \setminus \Omega^\epsilon$. Therefore, using the extension by zero to whole space and the characteristic function χ_ϵ of Ω^ϵ , we get

$$-\beta_1^\epsilon \int_{\Omega} \tilde{\phi}^\epsilon(x) \varphi(x) dx = \int_{\Omega} \chi_\epsilon(x) \varphi(x) N(\tilde{\phi}^\epsilon)(x) dx - \int_{\Omega} \varphi(x) \tilde{\phi}^\epsilon(x) dx. \tag{3.20}$$

Due to $\|\phi^\epsilon\|_{L^2(\Omega^\epsilon)} = 1$, we can extract a subsequence, still denoted by $\tilde{\phi}^\epsilon$, such that

$$\tilde{\phi}^\epsilon \rightharpoonup \phi^* \quad \text{weakly in } L^2(\Omega) \tag{3.21}$$

with $\phi^* \equiv 0$ as $x \in \mathbb{R}^N \setminus \Omega$.

First, let us discuss the case $\phi^* = 0$. From (3.20) with $\varphi = \tilde{\phi}^\epsilon$, we have that

$$(1 - \beta_1^\epsilon) = \int_{\Omega} \tilde{\phi}^\epsilon(x) N(\tilde{\phi}^\epsilon)(x) dx, \tag{3.22}$$

since $\|\phi^\epsilon\|_{L^2(\Omega^\epsilon)} = 1$. Arguing as in (3.1), we can show from assumptions (3.21) and the fact that $\phi^*(x) = 0$ a.e. in \mathbb{R}^N that, as $\epsilon \rightarrow 0$,

$$N(\tilde{\phi}^\epsilon) \rightarrow 0 \quad \text{strongly in } L^2(\Omega).$$

Hence, it follows from (3.21) and (3.22) that $\beta_1^\epsilon \rightarrow 1$ proving our result. Therefore, we can suppose $\phi^* \neq 0$ in $L^2(\Omega)$.

Now, we can argue as in (3.5) to pass to the limit in (3.20) obtaining the following limit equation

$$\begin{aligned} -\beta_1 \int_{\Omega} \phi^*(x) \varphi(x) dx &= \int_{\Omega} \mathcal{X}(x) \varphi(x) N(\phi^*)(x) dx - \int_{\Omega} \varphi(x) \phi^*(x) dx \\ &= \int_{\Omega} \mathcal{X}(x) \varphi(x) \int_{\mathbb{R}^N} J(x-y) (\phi^*(y) - \phi^*(x)) dy dx \\ &\quad - \int_{\Omega} (1 - \mathcal{X}(x)) \phi^*(x) \varphi(x) dx, \end{aligned} \tag{3.23}$$

which can be rewritten as

$$-\beta_1 \phi^*(x) = \mathcal{X}(x) \int_{\mathbb{R}^N} J(x-y) (\phi^*(y) - \phi^*(x)) dy - (1 - \mathcal{X}(x)) \phi^*(x), \quad a.e. \Omega, \tag{3.24}$$

with $\phi^*(x) \equiv 0$, $x \in \mathbb{R}^N \setminus \Omega$.

If $\mathcal{X}(x) \equiv 1$ a.e. Ω , it follows from (3.24) that β_1 is the first eigenvalue of the self-adjoint operator $T : L^2(\Omega) \mapsto L^2(\Omega)$ given by

$$T(\phi) = \int_{\mathbb{R}^N} J(x-y) (\phi(y) - \phi(x)) dy.$$

Hence, from [2, proposition 2.3] we get $\beta_1 > 0$, and the proof is completed.

Thus, let us assume $\mathcal{X} \neq 1$ in $L^\infty(\Omega)$, and consider the set D given by the vanishing points of \mathcal{X} introduced in (3.8). From (3.24), we get $\phi^*(x)(1 - \beta_1) = 0$ for all $x \in D$. If $\beta_1 = 1$, the proof is complete. If it is not the case, we have $\phi^*(x) \equiv 0$ in $x \in D$. Let us suppose $\phi^*(x) \equiv 0$ in D . In fact, without loss of generality, we can assume $\phi^*(x) \neq 0$ in $\Omega \setminus D$. From (3.24)

$$\frac{1 - \mathcal{X}(x) - \beta_1}{\mathcal{X}(x)} \phi^*(x) = \int_{\mathbb{R}^N} J(x - y)(\phi^*(y) - \phi^*(x)) \, dy, \quad \text{a.e. } \Omega \setminus D, \quad (3.25)$$

with $\phi^*(x) \equiv 0$ wherever $x \in \mathbb{R}^N \setminus (\Omega \setminus D)$. Since $\phi^* \in L^2(\Omega)$, it follows from (3.25) that $(1 - \mathcal{X} - \beta_1)/\mathcal{X} \phi^*$ also belongs to $L^2(\Omega \setminus D)$. Then, we also obtain from (3.25) that

$$\begin{aligned} & \int_{\Omega \setminus D} \frac{1 - \mathcal{X}(x) - \beta_1}{\mathcal{X}(x)} (\phi^*(x))^2 \, dx \\ &= -\frac{1}{2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J(x - y)(\phi^*(y) - \phi^*(x))^2 \, dy dx < 0. \end{aligned}$$

Consequently, there exists an open set $\tilde{D} \subset \Omega \setminus D$ such that

$$1 - \mathcal{X}(x) - \beta_1 < 0, \quad \forall x \in \tilde{D}.$$

Hence, since $0 \leq \mathcal{X}(x) \leq 1$, we conclude that $\beta_1 > 0$ finishing the proof. □

REMARK 3.6. Notice that $\beta_1 = 1$ is not an eigenvalue of the operator

$$L^2(\Omega) \ni \phi \rightarrow \mathcal{X}(\cdot) \int_{\mathbb{R}^N} J(\cdot - y)(\phi(y) - \phi(\cdot)) \, dy - (1 - \mathcal{X}(\cdot))\phi(\cdot) \in L^2(\Omega).$$

Indeed, it is equivalent to the existence of a nonzero function $\phi^* \in L^2(\Omega)$ satisfying

$$-\phi^*(x) = \int_{\mathbb{R}^N} J(x - y)(\phi^*(y) - \phi^*(x)) \, dy, \quad \text{a.e. } \Omega \setminus D, \quad (3.26)$$

with $\phi^*(x) \equiv 0$ in $\mathbb{R}^N \setminus (\Omega \setminus D)$ where $D \subset \Omega$ is the set introduced in (3.8). Now, it is known from [2, proposition 2.2] that (3.26) implies $\phi^*(x) \equiv 0$ in $\Omega \setminus D$. Since $\phi^*(x) = 0$ in D , we get the contradiction.

Finally, we point out why we have not proceeded as in the proof of lemma 3.5 to obtain a positive lower bound for the eigenvalues of the nonlocal Neumann problem. In fact, we cannot pass to the limit in the eigenvalue problem associated with (2.9) with (2.10)

$$-\lambda_1^\epsilon \int_{\Omega^\epsilon} \phi^\epsilon(x) \varphi(x) \, dx = \int_{\Omega^\epsilon} \varphi(x) \int_{\mathbb{R}^N \setminus A^\epsilon} J(x - y)(\phi^\epsilon(y) - \phi^\epsilon(x)) \, dy dx, \quad x \in \Omega^\epsilon,$$

with ϕ^ϵ and $\varphi \in L^2(\mathbb{R}^N)$ vanishing in $\mathbb{R}^N \setminus \Omega$. We can extract a convergent subsequence of eigenvalues and eigenfunctions, but we are not able to evaluate their limits in order to show that they are nontrivial. This drives us to hypothesis

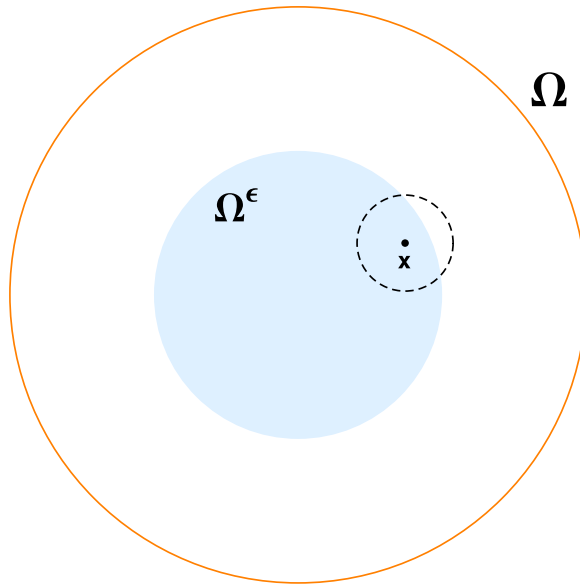


Figure 3. An example where the first eigenvalue to the Neumann problem is zero.

(**H_N**) and the following example that shows that in fact there are configurations of the holes for which the first eigenvalue is zero.

EXAMPLE 3.7. As illustrated in figure 3, let $\Omega = B_6$ and $\Omega^\epsilon = B_3$ the balls of radius 6 and 3 centred at the origin in \mathbb{R}^N respectively, hence the annulus $A^\epsilon = \Omega \setminus \Omega^\epsilon$ is the hole. Now assume that the kernel J satisfies hypothesis (**H_J**) and has the unit ball B_1 as its support. Then, we have that the function

$$u(x) = \begin{cases} 1, & x \in \Omega^\epsilon \\ 0, & x \in \mathbb{R}^N \setminus \Omega \end{cases}$$

satisfies

$$\frac{\frac{1}{2} \int_{\mathbb{R}^N \setminus A^\epsilon} \int_{\mathbb{R}^N \setminus A^\epsilon} J(x-y)(u(y) - u(x))^2 dy dx}{\int_{\Omega^\epsilon} u^2(x) dx} = 0$$

since $\text{supp}(J) = B_1$ implies $J(x-y)(u(y) - u(x)) = 0$ whenever $x, y \in \mathbb{R}^N \setminus A^\epsilon$. Consequently, hypothesis (**H_J**) it is not enough to guarantee the first eigenvalue of the Neumann problem introduced in (2.11) is strictly positive as in the Dirichlet case.

4. Periodically Perforated Domains

In this section, we discuss theorems 1.1 and 1.2 in the particular situation where the holes in $\Omega \subset \mathbb{R}^N$ are periodically distributed.

First, we deal with the critical case, it means, that one in which the size and distribution of holes possess the same order. Next, we consider the other cases.

4.1. Size and distribution with same order

Let $Q \subset \mathbb{R}^N$ be the *representative cell*

$$Q = (0, l_1) \times (0, l_2) \times \cdots \times (0, l_N).$$

We perforate Ω removing from it a set A^ϵ of periodically distributed holes given as follows: Take any open set $A \subset Q$ such that $T = Q \setminus A$ is measurable set satisfying $|T| \neq 0$. Denote by $\tau_\epsilon(A)$ the set of all translated images of $\epsilon\bar{A}$ of the form $\epsilon(kl + A)$ where $k \in \mathbb{Z}^N$ and $kl = (k_1l_1, \dots, k_Nl_N)$. Now define

$$A^\epsilon = \Omega \cap \tau_\epsilon(A).$$

We introduce our perforated domain as

$$\Omega^\epsilon = \Omega \setminus A^\epsilon. \tag{4.1}$$

Note that when considering Ω^ϵ we have removed from Ω a large number of holes of size $|\epsilon\bar{A}|$ which are ϵ -periodically distributed. A^ϵ represent the sets of holes inside Ω . It contains the interior holes, that ones that are fully contained into Ω , as well as part of each hole that intersects the boundary $\partial\Omega$.

We now pass to the limit in the characteristic function χ_ϵ to obtain the limit equations to Dirichlet and Neumann problems in the family of perforated domains given by (4.1). To do that let χ_A be the characteristic function of the open set $A \subset Q$ extended periodically in \mathbb{R}^N . Hence, if χ_{A^ϵ} is the characteristic function of A^ϵ , for each $x \in A^\epsilon$ there exist $k \in \mathbb{Z}^N$ such that

$$\chi_{A^\epsilon}(x) = \chi_A\left(\frac{x - \epsilon kl}{\epsilon}\right) = \chi_A(x/\epsilon).$$

Therefore, if χ_Ω and χ_ϵ are the characteristic functions of Ω and Ω^ϵ respectively, we have the following relationship

$$\chi_\epsilon(x) = \chi_\Omega(x) - \chi_{A^\epsilon}(x). \tag{4.2}$$

It follows from the Average theorem [11, theorem 2.6] that

$$\chi_{A^\epsilon} \rightharpoonup \frac{1}{|Q|} \int_Q \chi_A(s) ds = \frac{|A|}{|Q|}, \quad \text{as } \epsilon \rightarrow 0, \tag{4.3}$$

weakly star in $L^\infty(\Omega)$. Hence, from (4.2), (4.3) and $|T| = |Q \setminus A| = |Q| - |A|$ we obtain that

$$\chi_\epsilon \rightharpoonup \frac{|T|}{|Q|} \quad \text{weakly}^* \text{ in } L^\infty(\Omega).$$

In this way, we can set

$$\mathcal{X}(x) = \frac{|T|}{|Q|} \chi_\Omega(x) \quad \text{in } \mathbb{R}^N \tag{4.4}$$

at assumption (1.1).

Thus, as a consequence of theorem 1.1 and (4.4), the extended solutions \tilde{u}^ϵ of the Dirichlet problem (2.1) and (2.2) weakly converge in $L^2(\Omega)$ to the solution u^* to

$$\frac{|T|}{|Q|}f(x) = \frac{|T|}{|Q|} \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) \, dy - \left(\frac{|Q| - |T|}{|Q|} \right) u^*(x), \quad x \in \Omega, \tag{4.5}$$

with $u^*(x) \equiv 0$, for $x \in \mathbb{R}^N \setminus \Omega$, which can be rewritten as

$$f(x) = \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) \, dy - \left(\frac{|Q| - |T|}{|T|} \right) u^*(x), \quad x \in \Omega, \tag{4.6}$$

whenever $|T| \neq 0$.

Now to get the limit equation to the Neumann problem (2.9) and (2.10), we first have to see that condition (\mathbf{H}_N) is verified. To do that, given $\delta > 0$ small we define the sets

$$B_0 = \mathbb{R}^N \setminus \Omega$$

and

$$B_j = \{x \in \Omega^\epsilon : \delta j < d(x, \partial\Omega) < \delta(j+1)\} \quad j = 1, \dots, L.$$

Note that

$$(\mathbb{R}^N \setminus A^\epsilon) \subset \bigcup_{j=0}^L B_j$$

Notice also that, for $x \in B_j$ we have

$$\begin{aligned} \int_{B_{j-1}} J(x-y) \, dy &\geq \int_{\{y \in \Omega \setminus A^\epsilon : \delta(j-1) + \delta/2 < d(y, \partial\Omega) < \delta j\}} J(x-y) \, dy \\ &\geq \left(\min_{|z| \leq 1 - \delta/2} J(z) \right) \times |\{y \in \Omega \setminus A^\epsilon : \delta(j-1) \\ &\quad + \delta/2 < d(y, \partial\Omega) < \delta j\}| > 0. \end{aligned} \tag{4.7}$$

Here we have used that J is continuous with $J(0) > 0$.

Also observe that the number of sets, L , as well as the lower bound for the α_j , depend only on δ and therefore they can be chosen independently of ϵ .

Therefore, we can conclude from lemma 2.1 that there exists ϵ_0 and $c > 0$ such that

$$\lambda_1^\epsilon > c > 0, \quad \forall \epsilon \in (0, \epsilon_0),$$

and then, due to theorem 1.2, we obtain that the limit equation of (2.9) and (2.10) is

$$\begin{aligned} f(x) &= \int_{\mathbb{R}^N} J(x-y)(u^*(y) - u^*(x)) \, dy \\ &\quad - \left(\frac{|Q| - |T|}{|T|} \right) u^*(x) \int_{\mathbb{R}^N \setminus \Omega} J(x-y) \, dy, \quad x \in \Omega, \end{aligned} \tag{4.8}$$

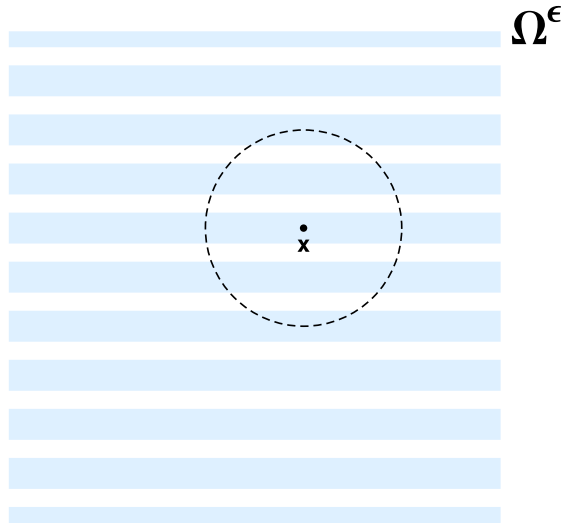


Figure 4. A nonconnected perforated domain.

with $u^*(x) \equiv 0$ for all $x \in \mathbb{R}^N \setminus \Omega$, where the term Λ defined in (3.14) can be calculated by

$$\begin{aligned} \Lambda(x) &= \int_{\mathbb{R}^N} J(x - y) \left(1 - \chi_\Omega(y) + \frac{|T|}{|Q|} \chi_\Omega(y) \right) dy - \frac{|T|}{|Q|} \chi_\Omega(x) \\ &= \left(\frac{|Q| - |T|}{|Q|} \right) \int_{\mathbb{R}^N} J(x - y) (1 - \chi_\Omega(y)) dy. \end{aligned}$$

REMARK 4.1. Notice that what makes distinction between the limit problems (4.5) and (4.8) here is just the coefficient (3.15).

Now, let us present an example that shows that we do not need the set Ω^ϵ to be connected. This has to be contrasted with what happens in the local case (where connectedness of Ω^ϵ plays a crucial role).

EXAMPLE 4.2. We consider $\Omega = (-1, 1) \times (-1, 1)$ the periodic case in which the hole is given by the scaled version of the strip $A = [0, 1] \times [1/3, 2/3]$, in the representative cell $(0, 1) \times (0, 1)$. That is, we remove from Ω the set of all translated images of ϵA of the form $\epsilon(k + A)$ where $k \in \mathbb{Z}^N$. Note that Ω^ϵ is a square from where we have removed a large number of horizontal strips of size $|\epsilon A|$ which are ϵ -periodically distributed. See figure 4 which illustrates this situation. Considering B_1, \dots, B_L as horizontal strips of width δ one can easily check that hypothesis (\mathbf{H}_N) holds (note that here δ and L can be chosen independently of ϵ).

4.2. Size and distribution with different orders

Now let us take holes in the open set Ω with distribution and size of different orders. To do that we consider the sets Q , A and B as in § 4.1, and define by $\hat{\tau}_\epsilon(A)$ the

union of all translated images of $\epsilon^\gamma \bar{A}$ of the form

$$\epsilon kl + \epsilon^\gamma A = \epsilon(kl + \epsilon^{\gamma-1}A)$$

where $\gamma > 0$, $k \in \mathbb{Z}^N$ and $kl = (k_1l_1, \dots, k_Nl_N)$. Here the hole set are given by

$$A^\epsilon = \Omega \cap \hat{\tau}_\epsilon(A)$$

and the perforated domain by

$$\Omega^\epsilon = \Omega \setminus A^\epsilon.$$

When $\gamma < 1$, the order of distributions is bigger than the size of the holes at $\epsilon = 0$. On the other hand, the size order of the holes is larger than their distributions as $\gamma > 1$. The case $\gamma = 1$ has been considered in § 4.1 and the sets are of the same order of the size and distribution of a typical cell.

If we assume $\gamma > 1$, the set of distributed holes vanishes in Ω as $\epsilon \rightarrow 0$, and then, Ω^ϵ fills the whole of Ω which implies

$$\chi_\epsilon \rightharpoonup \chi_\Omega \quad \text{weakly* in } L^\infty(\Omega) \tag{4.9}$$

where χ_Ω is the characteristic function of Ω .

Then, as a consequence of theorems 1.1 and 1.2, we get the same limit equation to both Dirichlet and Neumann problems (2.1) and (2.9) respectively. We get the nonlocal Dirichlet problem in the open set Ω

$$f(x) = \int_{\mathbb{R}^N} J(x - y)(u^*(y) - u^*(x)) dy, \quad x \in \Omega, \tag{4.10}$$

with $u^*(x) \equiv 0$, in $\mathbb{R}^N \setminus \Omega$.

Notice that hypothesis (\mathbf{H}_N) to the Neumann problem can be easily verified here as in (4.7). Among other things, it means that under the assumptions (4.9), (\mathbf{H}_J) and (\mathbf{H}_N) , the nonlocal Neumann problem (2.9) behaves as the nonlocal Dirichlet problem for ϵ small enough.

Finally, if $\gamma < 1$, it is not difficult to see that the set of the holes fill the whole of Ω when ϵ goes to zero implying that

$$\chi_\epsilon \rightharpoonup 0 \quad \text{weakly* in } L^\infty(\Omega).$$

Hence, it follows from theorems 1.1 and 1.2 that the extended solutions \tilde{u}^ϵ to both Neumann and Dirichlet problems weakly converge to zero in $L^2(\Omega)$ as $\epsilon \rightarrow 0$. Moreover, from corollary 3.1, we still obtain strong convergence to zero in the Dirichlet case. That is, if u^ϵ satisfies the Dirichlet problem (2.1) and (2.2), then

$$\tilde{u}^\epsilon \rightarrow 0 \quad \text{strongly in } L^2(\Omega)$$

as $\epsilon \rightarrow 0$.

With the additional condition (3.16), we also have $\tilde{u}^\epsilon \rightarrow 0$ strongly in $L^2(\Omega)$ for the Neumann problem (2.9) with (2.10). See corollary 3.3.

5. Rescaling the kernel

In this section, our aim is to see how these problems behave when we rescale the kernel in order to approximate local equations. Hence, now we have two different parameters ϵ (that controls the size of the holes) and δ (that is used to rescale the kernel). Our aim is to study the limits $\epsilon \rightarrow 0$ (to have an homogenized limit problem) and $\delta \rightarrow 0$ (to approach local problems).

To describe accurately these limits along this section we need to restrict ourselves to the periodic case. We assume that we have a bounded domain Ω from where we have removed a big number of periodic small balls (the holes). That is, we consider $\Omega^\epsilon = \Omega \setminus \cup B_{r^\epsilon}(x_i)$ where $B_{r^\epsilon}(x_i)$ is a ball centred in $x_i \in \Omega$ of the form $x_i \in 2\epsilon\mathbb{Z}^N$ with $0 < r^\epsilon < \epsilon \leq 1$. To simplify, we only remove here balls that are strictly contained in Ω . Moreover, we will assume here that J is radially symmetric.

Our aim is to show that, in general, the involved limits do not commute, but it holds that both limits exist. For $u^{\epsilon,\delta}$ a solution to the nonlocal problem in Ω^ϵ with a kernel rescaled with δ we have that

$$\lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \tilde{u}^{\epsilon,\delta} = w,$$

and

$$\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} \tilde{u}^{\epsilon,\delta} = v,$$

but, in general

$$w \neq v.$$

5.1. The Dirichlet case

We start by considering

$$f(x) = \frac{C}{\delta^{N+2}} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{\epsilon,\delta}(y) - u^{\epsilon,\delta}(x)) dy, \quad x \in \Omega^\epsilon, \tag{5.1}$$

with

$$u^{\epsilon,\delta}(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega^\epsilon, \tag{5.2}$$

and C is a normalizing constant given by

$$C = \left(\frac{1}{2} \int_{\mathbb{R}^N} J(x) x_1^2 dx \right)^{-1} \tag{5.3}$$

where x_1 is the first coordinate of $x \in \mathbb{R}^N$.

Existence and uniqueness of the solutions $u^{\epsilon,\delta}$ of (5.1) are guaranteed in § 2.1 for any $f \in L^2(\Omega)$, ϵ and $\delta > 0$. Thus, we proceed with the analysis of the behaviour of $u^{\epsilon,\delta}$ as ϵ and δ go to zero.

Performing the change of variable

$$z = \frac{(x - y)}{\delta}$$

and using Taylor expansion, we obtain for a smooth u

$$\begin{aligned} C \frac{1}{\delta^{2+N}} \int_{\mathbb{R}^N} J\left(\frac{x - y}{\delta}\right) (u(y) - u(x)) dy &= \frac{C}{\delta^2} \int_{\mathbb{R}^N} J(z) (u(x - \delta z) - u(x)) dz \\ &= \frac{C}{2} \left[\sum_{i=1}^N \partial_i^2 u(x) \right] \int_{\mathbb{R}^N} J(z) z_i^2 dz + O(\delta) = \Delta_x u(x) + O(\delta). \end{aligned} \tag{5.4}$$

Expression (5.4) makes a connection between the nonlocal problem with the kernel J_δ and the following boundary value problem

$$\begin{cases} \Delta v^\epsilon(x) = f(x), & x \in \Omega^\epsilon \\ v^\epsilon(x) = 0, & x \in \partial\Omega^\epsilon. \end{cases} \tag{5.5}$$

Now, we can use the results in [2, § 3.2.2] (see also [19]) to obtain that

$$\|u^{\delta,\epsilon} - v^\epsilon\|_{L^2(\Omega)} \rightarrow 0, \quad \text{as } \delta \rightarrow 0. \tag{5.6}$$

As we can see from [12], problem (5.5) is the prototype for the study of the Laplacian in perforated domains. As we have mentioned in the introduction, the limit as $\epsilon \rightarrow 0$ is given by (1.4). Hence, we have that

$$\|v^\epsilon - v\|_{L^2(\Omega)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0.$$

Now, we just observe that we have obtained the following result.

THEOREM 5.1. *If $u^{\delta,\epsilon}$, v are given by (5.1) and (1.4) respectively, it holds that*

$$\lim_{\epsilon \rightarrow 0} (\lim_{\delta \rightarrow 0} u^{\delta,\epsilon}) = v \quad \text{in } L^2(\Omega),$$

Now, we let first $\epsilon \rightarrow 0$ and then $\delta \rightarrow 0$. To compute the first limit we rewrite our problem (5.1) as follows:

$$\frac{\delta^2}{C} f(x) = \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x - y}{\delta}\right) (u^{\epsilon,\delta}(y) - u^{\epsilon,\delta}(x)) dy, \quad x \in \Omega^\epsilon.$$

Note that the involved kernel satisfies

$$\frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{z}{\delta}\right) dz = 1.$$

This property was used in the proof of theorem 1.1 in § 3. Arguing as in § 3 with a fixed δ we obtain that there exists $u^{*,\delta} \in L^2(\Omega)$ such that

$$\tilde{u}^{\epsilon,\delta} \rightharpoonup u^{*,\delta} \text{ weakly in } L^2(\Omega).$$

Moreover, the limit $u^{*,\delta}$ satisfies the following nonlocal problem in Ω

$$\begin{aligned} \frac{\delta^2}{C} \mathcal{X}(x) f(x) &= \mathcal{X}(x) \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{*,\delta}(y) \\ &\quad - u^{*,\delta}(x)) \, dy - (1 - \mathcal{X}(x)) u^{*,\delta}(x) \end{aligned} \tag{5.7}$$

with $u^{*,\delta}(x) \equiv 0$, for $x \in \mathbb{R}^N \setminus \Omega$. Note that, since we are considering periodic holes, we have that \mathcal{X} is a constant given by

$$\mathcal{X} = \begin{cases} 1 & \text{if } r^\epsilon \ll C_0\epsilon, \\ \frac{|Q \setminus B|}{|Q|} & \text{if } r^\epsilon = C_0\epsilon. \end{cases} \tag{5.8}$$

Recall that $|Q \setminus B|$ is the measure of the complement of the ball B in the cube Q . Hence, here $|Q \setminus B|/|Q|$ is the proportion of the cube that is inside Ω^ϵ . Also, recall that in this case the critical size is $b^\epsilon := C_0\epsilon$.

Now, for the limit as $\delta \rightarrow 0$ we consider two cases. First, when $r^\epsilon \ll b^\epsilon$, we have $\mathcal{X} = 1$ and (5.7) reduces to

$$f(x) = \frac{C}{\delta^{N+2}} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{*,\delta}(y) - u^{*,\delta}(x)) \, dy,$$

and from the results in [2, § 3.2.2] we get that

$$u^{*,\delta} \rightarrow w, \quad \text{in } L^2(\Omega),$$

as $\delta \rightarrow 0$, where w is the solution to $\Delta w = f$ with $w = 0$ on $\partial\Omega$.

In the case $r^\epsilon = b^\epsilon$ we have

$$\frac{\delta^2}{C} f(x) = \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{*,\delta}(y) - u^{*,\delta}(x)) \, dy - \frac{(1 - \mathcal{X})}{\mathcal{X}} u^{*,\delta}(x).$$

Multiplying by $u^{*,\delta}$ and integrating we get

$$\begin{aligned} \int_{\mathbb{R}^N} \frac{\delta^2}{C} f(x) u^{*,\delta}(x) \, dx &= -\frac{1}{2} \frac{1}{\delta^N} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{*,\delta}(y) - u^{*,\delta}(x))^2 \, dy \, dx \\ &\quad - \frac{(1 - \mathcal{X})}{\mathcal{X}} \int_{\mathbb{R}^N} (u^{*,\delta}(x))^2 \, dx. \end{aligned}$$

It follows that

$$\begin{aligned} &\frac{(1 - \mathcal{X})}{\mathcal{X}} \int_{\mathbb{R}^N} (u^{*,\delta}(x))^2 \, dx \\ &\leq \left| \int_{\mathbb{R}^N} \frac{\delta^2}{C} f(x) u^{*,\delta}(x) \, dx \right| \leq C(f) \delta^2 \left(\int_{\mathbb{R}^N} (u^{*,\delta}(x))^2 \, dx \right)^{1/2}, \end{aligned}$$

and we conclude that

$$u^{*,\delta} \rightarrow 0, \quad \text{in } L^2(\Omega),$$

as $\delta \rightarrow 0$. Hence, we have obtained that

$$\|u^{*,\delta} - w\|_{L^2(\Omega)} \rightarrow 0, \quad \text{as } \delta \rightarrow 0,$$

where w is given by

$$w = \begin{cases} \text{the solution to } \Delta w = f \text{ with } w = 0 \text{ on } \partial\Omega, & \text{if } r^\epsilon \ll b^\epsilon, \\ 0, & \text{if } r^\epsilon = b^\epsilon. \end{cases} \quad (5.9)$$

We have obtained the following theorem.

THEOREM 5.2. *If $u^{\delta,\epsilon}$, w are given by (5.1) and (5.9) respectively, it holds that*

$$\lim_{\delta \rightarrow 0} (\lim_{\epsilon \rightarrow 0} \tilde{u}^{\delta,\epsilon}) = w \quad \text{weakly in } L^2(\Omega),$$

Now, let us see when v and w coincide. We have to distinguish several cases according to the size of the holes. Notice $b^\epsilon \geq a^\epsilon$ whenever ϵ is small enough.

Case 1. $r^\epsilon = b^\epsilon$. In this case, we have that $w = 0$ and $v = 0$. Therefore $w = v$ in this case.

Case 2. $a^\epsilon \ll r^\epsilon \ll b^\epsilon$. In this case, w is the solution to $\Delta w = f$ and $v = 0$. Therefore $w \neq v$ in this case.

Case 3. $r^\epsilon = a^\epsilon$. In this case, w is the solution to $\Delta w = f$ and v the solution to $\Delta v - \mu v = f$. Therefore, also $w \neq v$ in this case.

Case 4. $r^\epsilon \ll a^\epsilon$. In this case, w and v coincide and are given by the unique solution to $\Delta w = f$.

5.2. The Neumann case

Now we consider

$$f(x) = \frac{C}{\delta^{N+2}} \int_{\mathbb{R}^N \setminus A^\epsilon} J\left(\frac{x-y}{\delta}\right) (u^{\epsilon,\delta}(y) - u^{\epsilon,\delta}(x)) dy, \quad x \in \Omega^\epsilon, \quad (5.10)$$

with

$$u^{\epsilon,\delta}(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega, \quad (5.11)$$

and C is the same normalizing constant that we used for the Dirichlet case and is given by (5.3).

Existence and uniqueness of the solutions $u^{\epsilon,\delta}$ of (5.10) are guaranteed in §2.2 for any $f \in L^2(\Omega)$, ϵ and $\delta > 0$, provided there is a positive constant c independent of ϵ such that

$$\lambda_1^\epsilon \geq c > 0.$$

We have seen that this holds in our case, that is, for periodically perforated domains.

Let us proceed with the analysis of the behaviour of $u^{\epsilon,\delta}$ as ϵ and δ go to zero. From (5.4) using results from [18] we have

$$\|u^{\delta,\epsilon} - v^\epsilon\|_{L^2(\Omega^\epsilon)} \rightarrow 0, \quad \text{as } \delta \rightarrow 0.$$

being v^ϵ the solution to the local problem

$$\begin{cases} \Delta v^\epsilon(x) = f(x), & x \in \Omega^\epsilon \\ v^\epsilon(x) = 0, & x \in \partial\Omega, \\ \frac{\partial v^\epsilon(x)}{\partial \eta} = 0, & x \in \partial\Omega^\epsilon \cap \Omega. \end{cases} \tag{5.12}$$

Now, using results from [14, theorem 2.16] (see also [13]) we get that the limit as $\epsilon \rightarrow 0$ in this local problem is given by

$$v = \begin{cases} \text{the solution to } \Delta v = f, & \text{if } r^\epsilon \ll b^\epsilon, \\ \text{the solution to } \sum_{i,j=1}^N q_{ij} \frac{\partial^2 v}{\partial x_i \partial x_j} = \frac{|Q \setminus B|}{|Q|} f, & \text{if } r^\epsilon = b^\epsilon, \end{cases} \tag{5.13}$$

with Dirichlet boundary conditions, $v = 0$ on $\partial\Omega$. Here q_{ij} are called the *homogenized coefficients* and are given by

$$q_{ij} = \frac{1}{|Q|} \left[\int_{Q \setminus B} \delta_{ij} \, dy - \int_{Q \setminus B} \frac{\partial X^i}{\partial y_j}(y) \, dy \right],$$

where δ_{ij} is the Kronecker delta and X^i (for $i = 1, \dots, N$) are the solutions of the system

$$\begin{cases} \Delta X^i = 0 & \text{in } Q \setminus B, \\ \partial_\eta X^i = \eta_i & \text{on } \partial B, \\ X^i & Q\text{-periodic} \end{cases}$$

with $\int_{Q \setminus B} X^i(y) \, dy = 0$. Here the critical size of the holes is $b^\epsilon := C_0\epsilon$,

Now let P_ϵ be an extension operator, that is, $P_\epsilon \in \mathcal{L}(L^2(\Omega^\epsilon); L^2(\Omega)) \cap \mathcal{L}(V_\epsilon; H_0^1(\Omega))$ where $V_\epsilon = \{v \in H^1(\Omega) : v = 0 \text{ on } \partial_{\text{ext}}\Omega^\epsilon\}$ (the existence of such extension operator was proved in [12]). Then, from [14, theorem 2.16] we have

$$\|P_\epsilon v^\epsilon - v\|_{L^2(\Omega)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0.$$

Thus, we conclude that

THEOREM 5.3. *If $u^{\delta,\epsilon}$, v are given by (5.10) and (5.13) respectively, it holds that*

$$\lim_{\epsilon \rightarrow 0} P_\epsilon \left(\lim_{\delta \rightarrow 0} u^{\delta,\epsilon} \right) = v \quad \text{in } L^2(\Omega).$$

Now, we reverse the order in which we take limits and let first $\epsilon \rightarrow 0$ and then $\delta \rightarrow 0$. First, as we did in the Dirichlet case we write our equation as

$$\frac{\delta^2}{C} f(x) = \frac{1}{\delta^N} \int_{\mathbb{R}^N \setminus A^\epsilon} J\left(\frac{x-y}{\delta}\right) (u^{\epsilon,\delta}(y) - u^{\epsilon,\delta}(x)) dy, \quad x \in \Omega^\epsilon,$$

since again in this case we want to use that the involved kernel satisfies

$$\frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{z}{\delta}\right) dz = 1.$$

Arguing as in § 3 with a fixed δ we obtain that there exists a limit $u^{*,\delta} \in L^2(\Omega)$ such that

$$\tilde{u}^{\epsilon,\delta} \rightharpoonup u^{*,\delta} \quad \text{weakly in } L^2(\Omega).$$

Moreover, the limit $u^{*,\delta}$ satisfies the following nonlocal problem in Ω

$$\begin{cases} \frac{\delta^2}{C} \mathcal{X}(x) f(x) = \mathcal{X}(x) \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^*(y) - u^*(x)) dy - \Lambda(x) u^*(x) \\ u^*(x) \equiv 0, \quad x \in \mathbb{R}^N \setminus \Omega, \end{cases} \tag{5.14}$$

where $\Lambda \in L^\infty(\Omega)$ is given by

$$\Lambda(x) = \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (1 - \chi_\Omega(y) + \mathcal{X}(y)) dy - \mathcal{X}(x), \quad x \in \Omega.$$

Here χ_Ω is the characteristic function of the open set Ω and \mathcal{X} is the constant given by (5.8), that is,

$$\mathcal{X} = \begin{cases} \frac{|Q \setminus B|}{|Q|} & \text{if } r^\epsilon = C_0\epsilon, \\ 1 & \text{if } r^\epsilon \ll C_0\epsilon. \end{cases}$$

Note that also in this case the critical size is $b^\epsilon := C_0\epsilon$.

Now, for the limit as $\delta \rightarrow 0$ we consider two cases. First, when $r^\epsilon \ll b^\epsilon$, we have $\mathcal{X} = 1$ and from the results in [2, § 3.2.2] we get that

$$u^{*,\delta} \rightarrow w, \quad \text{in } L^2(\Omega),$$

as $\delta \rightarrow 0$, where w is the solution to $\Delta w = f$ with $w = 0$ on $\partial\Omega$.

In the case $r^\epsilon = b^\epsilon$ we have

$$\frac{\delta^2}{C} f(x) = \frac{1}{\delta^N} \int_{\mathbb{R}^N} J\left(\frac{x-y}{\delta}\right) (u^{*,\delta}(y) - u^{*,\delta}(x)) dy - \Lambda u^{*,\delta}(x).$$

Multiplying by $u^{*,\delta}$ and integrating we get, arguing as in the Dirichlet case and using lemma 2.1, that

$$u^{*,\delta} \rightarrow 0, \quad \text{in } L^2(\Omega),$$

as $\delta \rightarrow 0$.

Hence, we have obtained that $u^{*,\delta}$ converges as $\delta \rightarrow 0$ to w that is given by

$$w = \begin{cases} \text{the solution to } \Delta w = f, \text{ with } w = 0 \text{ on } \partial\Omega, & \text{if } r^\epsilon \ll b^\epsilon, \\ 0, & \text{if } r^\epsilon = b^\epsilon. \end{cases} \quad (5.15)$$

Therefore, we have the following theorem.

THEOREM 5.4. *If $u^{\delta,\epsilon}$, w are given by (5.10) and (5.15) respectively, it holds that*

$$\lim_{\delta \rightarrow 0} \left(\lim_{\epsilon \rightarrow 0} \tilde{u}^{\delta,\epsilon} \right) = w \quad \text{weakly in } L^2(\Omega).$$

Now, let us see when v and w coincide. We have to distinguish only two cases according to the size of the holes:

Case 1. $r^\epsilon = b^\epsilon$. In this case we have that $w = 0$ and v is the solution to

$$\sum_{i,j=1}^N q_{ij} \frac{\partial^2 v}{\partial x_i \partial x_j} = \frac{|Q \setminus B|}{|Q|} f.$$

Note that $w \neq v$ in this case.

Case 2. $r^\epsilon \ll b^\epsilon$. In this case, w and v coincide and are given by the unique solution to $\Delta w = f$.

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