

Existence of multiple periodic solutions to a semilinear wave equation with x-dependent coefficients

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(MS received 6 June 2018; accepted 24 March 2019)

This paper is concerned with the periodic (in time) solutions to an one-dimensional semilinear wave equation with x-dependent coefficients. Such a model arises from the forced vibrations of a nonhomogeneous string and propagation of seismic waves in nonisotropic media. By combining variational methods with saddle point reduction technique, we obtain the existence of at least three periodic solutions whenever the period is a rational multiple of the length of the spatial interval. Our method is based on a delicate analysis for the asymptotic character of the spectrum of the wave operator with x-dependent coefficients, and the spectral properties play an essential role in the proof.

Keywords: existence; periodic solutions; wave equation

2010 Mathematics subject classification: 35L71; 35B10

1. Introduction

In this paper, we consider the existence of multiple periodic solutions to the semilinear wave equation with x-dependent coefficients

$$\rho(x)u_{tt} - (\rho(x)u_x)_x = a\rho(x)u + f(t, x, u), \ t \in \mathbb{R}, \ 0 < x < \pi,$$
(1.1)

with the Dirichlet boundary conditions

$$u(t,0) = u(t,\pi) = 0, \ t \in \mathbb{R},$$
(1.2)

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and the periodic conditions

$$u(t+T,x) = u(t,x), \ u_t(t+T,x) = u_t(t,x), \ t \in \mathbb{R}, \ 0 < x < \pi,$$
(1.3)

where a > 0 is a constant and f is a given T-periodic function in time t.

Equation (1.1) originates from the forced vibrations of a bounded nonhomogeneous string and the propagation of seismic waves in nonisotropic media (see e.g. [3-5, 18-22, 27, 28] and references therein). More precisely, the vertical displacement u(t, z) at time t and depth z of a plane seismic wave is described by the equation

$$\mu(z)u_{tt} - (\nu(z)u_z)_z = 0 \tag{1.4}$$

with some initial conditions in t and boundary conditions in z, where $\mu(z)$ is the rock density and $\nu(z)$ is the elasticity coefficient. By the change of variable

$$x = \int_0^z \left(\frac{\mu(s)}{\nu(s)}\right)^{1/2} \mathrm{d}s$$

equation (1.4) is transformed into

$$\rho(x)u_{tt} - (\rho(x)u_x)_x = 0,$$

where $\rho = (\mu \nu)^{1/2}$ is called the acoustic impedance function.

It is well known that the case of $\rho \equiv C$ (a nonzero constant) corresponds to the classical wave equation, which is called the one with constant coefficients for distinguishing it from the one with x-dependent coefficients discussed here. The problem of finding periodic solutions of nonlinear wave equation with constant coefficients has received a great deal of attention since the original work [24] of Rabinowitz. By using the variational methods, he obtained the existence of periodic solutions for the weakly nonlinear homogeneous string whenever the period T is a rational multiple of the length of the spatial interval. Thereafter, many authors, such as Brézis, Chang, Nirenberg, etc., have used and developed the variational methods, topological degree and index theory to obtain a lot of results on the existence and multiplicity of periodic solutions for the problem with various nonlinearities (see e.g. [1, 6–8, 10, 12–17, 25, 26, 29, 30] and the references therein).

On the other hand, the problem of finding periodic solutions for the nonlinear wave equation with x-dependent coefficients was studied by Barbu and Pavel in [3–5] for the first time. In [5], Barbu and Pavel considered the existence and regularity of periodic solutions for such wave equation with sublinear nonlinearity under the Dirichlet boundary conditions. For the case the nonlinear term having power-law growth, Rudakov [27] proved the existence of periodic solutions under the Dirichlet boundary conditions. Later, Ji and his collaborators obtained some related results for the general Sturm–Liouville boundary value problem [18, 20], and periodic and anti-periodic boundary value problem [19, 21]. In [31], by using topological degree methods, Wang and An obtained an existence result on periodic solution of the problem with resonance and the sublinear nonlinearity. Afterwards, Ji and Li [22] obtained an existence result of periodic solution for $\eta_{\rho}(x) = 0$ under the Dirichlet boundary conditions, which actually solves an open problem posted in [5]. Recently,

Ji et al. [23] obtained the existence and multiplicity of periodic solutions for the Dirichlet–Neumann boundary value problem of a wave equation with x-dependent coefficients by using the Leray–Schauder degree theory. The restriction to such type of boundary value problem is essentially due to the possible loss of the compactness of the inverse operator on its range.

In this paper, we intend to pay close attention to the existence of multiple periodic solutions of wave equation with x-dependent coefficients. By combining variational methods with saddle point reduction technique, we obtain the existence of at least three periodic solutions whenever the period is a rational multiple of the length of the spatial interval. Our method is based on a delicate analysis for the asymptotic character of the spectrum of the wave operator with x-dependent coefficients, and the spectral properties play an essential role in the proof.

Denote $f(t, x, u) = (f(t, x, u))/(\rho(x))$. Throughout this paper, we assume T is a rational multiple of π which can be rewritten as

$$T = 2\pi \frac{p}{q}$$

for some relatively prime positive integers p and q. Moreover, we make the following assumptions:

(A1) $\rho(x) \in H^2(0,\pi)$ satisfies $\rho(x) > 0$ for $x \in [0,\pi]$, and

$$\rho_0 = \operatorname{ess\,inf} \eta_\rho(x) > 0,$$

where

$$\eta_{\rho}(x) = \frac{1}{2}\frac{\rho''}{\rho} - \frac{1}{4}\left(\frac{\rho'}{\rho}\right)^2$$

(A2) $\tilde{f}(t,x,u) \in C^1(\mathbb{R} \times (0,\pi) \times \mathbb{R}), \ \tilde{f}(t+T,x,u) = \tilde{f}(t,x,u),$ and

$$\hat{f}(t, x, u) = o(|u|), \text{ as } u \to 0 \text{ uniformly in } (t, x),$$
 (1.5)

and $\tilde{f}(t, x, u)$ is asymptotically linear in u at ∞ in the following sense: there exists a constant b > 0 such that

$$\tilde{f}(t, x, u) - bu = o(|u|), \text{ as } |u| \to \infty \text{ uniformly in } (t, x).$$
 (1.6)

The rest of this paper is organized as follows. In §2, we give some preliminaries and state the main result. In §3, we first characterize the solutions of problem (1.1)-(1.3) as the critical points of the corresponding variational problem. Then, with the aid of the saddle point reduction technique, we reduce the critical point of the variational problem from an infinite dimensional space to a finite dimensional subspace. In §4, we prove that reduced functional satisfies the $(PS)_c$ condition for any $c \in \mathbb{R}$. In §§ 5 and 6, we devote to the proof of the bounds of reduced functional and the main result, respectively.

2. Preliminaries and main result

Set $\Omega = (0, T) \times (0, \pi)$, and denote

$$\Psi = \{ \psi \in C^{\infty}(\Omega) : \psi(t,0) = \psi(t,\pi) = 0, \psi(0,x) = \psi(T,x), \psi_t(0,x) = \psi_t(T,x) \}.$$

Let

$$L^{r}(\Omega) = \left\{ u : \|u\|_{L^{r}(\Omega)}^{r} = \int_{\Omega} \rho(x) |u(t,x)|^{r} \mathrm{d}t \mathrm{d}x < \infty \right\}, \quad r \ge 1.$$

It is well known that Ψ is dense in $L^r(\Omega)$ for any $r \ge 1$, and $L^2(\Omega)$ is a Hilbert space with the inner product

$$\langle u, v \rangle = \int_{\Omega} \rho(x) u(t, x) \overline{v(t, x)} dt dx, \quad \forall u, \quad v \in L^{2}(\Omega).$$

We rewrite (1.1)-(1.3) on Ω in the following form

$$\rho(x)u_{tt} - (\rho(x)u_x)_x = a\rho(x)u + f(t, x, u), \quad (t, x) \in \Omega,$$
(2.1)

$$u(t,0) = u(t,\pi) = 0, \quad t \in (0,T),$$
(2.2)

$$u(0,x) = u(T,x), \quad u_t(0,x) = u_t(T,x), \quad x \in (0,\pi).$$
 (2.3)

DEFINITION 2.1. A function $u \in L^r(\Omega)$ is called a weak solution of problem (2.1)–(2.3) if it satisfies

$$\int_{\Omega} u(\rho\psi_{tt} - (\rho\psi_x)_x) dt dx - \int_{\Omega} \rho(au + \tilde{f}(t, x, u))\psi dt dx = 0, \quad \forall \psi \in \Psi.$$

Define the linear operator L_0 by

$$L_0\psi = \rho^{-1} \left(\rho\psi_{tt} - (\rho\psi_x)_x\right), \ \forall \psi \in \Psi,$$

and denote its extension in $L^2(\Omega)$ by L. It is known that L is a self-adjoint operator (see [5]), and $u \in L^2(\Omega)$ is a weak solution of problem (2.1)–(2.3) if and only if $Lu = au + \tilde{f}$.

For the study of periodic solutions of problem (2.1)–(2.3), we need to use the following complete orthonormal system of eigenfunctions $\{\phi_m(t)\varphi_n(x): m \in \mathbb{Z}, n \in \mathbb{N}\}$ in $L^2(\Omega)$ (see [32]), where

$$\phi_m(t) = T^{-(1/2)} e^{i\mu_m t}, \quad \mu_m = 2m\pi T^{-1}, \quad m \in \mathbb{Z},$$

and λ_n , $\varphi_n(x)$ are given by the Sturm–Liouville problem

$$-\left(\rho(x)\varphi_n'(x)\right)' = \lambda_n^2 \rho(x)\varphi_n(x), \ \varphi_n(0) = \varphi_n(\pi) = 0, \ n \in \mathbb{N}.$$
 (2.4)

LEMMA 2.2 ([5]). Assume that $\rho(x)$ satisfies (A1), then the eigenvalues of problem (2.4) have the form

$$\lambda_n = n + \theta_n \text{ with } \theta_n \to 0 \text{ as } n \to \infty,$$

where

$$0 < \frac{\rho_2}{n} \leqslant \sqrt{n^2 + \rho_0} - n \leqslant \theta_n \leqslant \sqrt{n^2 + \rho_1} - n \leqslant \frac{\rho_1}{2n}, \tag{2.5}$$

and $\rho_1 = \frac{2}{\pi} \int_0^{\pi} \eta_{\rho}(x) dx$, $\rho_2 = \sqrt{\rho_0 + 1} - 1$.

By Lemma 2.2, the eigenvalues of operator L can be rewritten as

$$\lambda_{nm} = \lambda_n^2 - \mu_m^2 = p^{-2}(np + \theta_n p - mq)(np + \theta_n p + mq).$$

Thus, when $np \neq |m|q$, it is easy to verify that $|\lambda_{nm}| \to \infty$ as $n, m \to \infty$. On the other hand, when np = |m|q, by (2.5) we have

$$2\rho_2 \leftarrow \left(\frac{\rho_2}{n}\right)^2 + 2\rho_2 \leqslant \lambda_{nm} = \theta_n(2n + \theta_n) \leqslant \rho_1 + \left(\frac{\rho_1}{2n}\right)^2 \to \rho_1$$

as $n \to \infty$.

Denote the set of eigenvalues of operator L by

$$\Lambda(L) = \{\lambda_{nm} : \lambda_{nm} = \lambda_n^2 - \mu_m^2\}.$$

The above statement shows that $\Lambda(L)$ has at least one accumulation point in $[2\rho_2, \rho_1]$. Therefore, we have the following lemma.

LEMMA 2.3. Let assumption (A1) hold. Then

- (i) L has at least one essential spectral point, and all of them belong to $[2\rho_2, \rho_1]$;
- (ii) If $\lambda \in \Lambda(L)$ and $\lambda \notin [2\rho_2, \rho_1]$, then λ is isolated and its multiplicity is finite.
- If $(a, b) \cap \Lambda(L) \neq \emptyset$ and b satisfies $\rho_1 < a < b$, by lemma 2.3, we can define

$$b^- = \max\{\lambda \in \Lambda(L) : \lambda < b\}, \quad b^+ = \min\{\lambda \in \Lambda(L) : \lambda > b\}$$

It is obvious that $b^- < b < b^+$.

Denote $e = b^+ - a$. The main result of this paper is as follows.

THEOREM 2.4. Assume $a, b \notin \Lambda(L), \rho_1 < a < b^+ - b, (a,b) \cap \Lambda(L) \neq \emptyset$, and (A1)–(A2) hold. If \tilde{f} is increasing in u and there exists a constant $\kappa > 0$ such that

$$\frac{\partial f}{\partial u}(t,x,u)\leqslant e-\kappa,\quad \forall (t,x,u)\in\mathbb{R}\times(0,\pi)\times\mathbb{R}$$

Then the problem (2.1)–(2.3) has at least three T-periodic solutions.

3. Variational problem and its reduction

In what follows, we always assume a, b satisfy the conditions in theorem 2.4. Since $a \notin \Lambda(L)$ and $a > \rho_1$, then there exists a constant $\delta > 0$ such that

$$|\lambda_{nm} - a| \ge \delta > 0, \quad n \in \mathbb{N}, \quad m \in \mathbb{Z}.$$
(3.1)

We define the working space

$$E = \left\{ u \in L^{2}(\Omega) : ||u||_{E}^{2} = \sum_{n,m} |\lambda_{nm} - a| |\alpha_{nm}|^{2} < \infty \right\},\$$

where α_{nm} denote the Fourier coefficients of $u \in L^2(\Omega)$, i.e.,

$$u = \sum_{n,m} \alpha_{nm} \varphi_n(x) \phi_m(t), \quad \alpha_{nm} = \int_{\Omega} \rho u \varphi_n \overline{\phi}_m dx dt.$$

The estimate (3.1) indicates $\|\cdot\|_E$ is a norm. Furthermore, E is a Hilbert space equipped with the inner product $\langle u, v \rangle_0 = \sum_{n,m} |\lambda_{nm} - a| \alpha_{nm} \overline{\beta}_{nm}$, where α_{nm} and β_{nm} are the Fourier coefficients of u and v, respectively.

From (3.1), we have

$$||u||_{L^{2}(\Omega)}^{2} = \sum_{n,m} |\alpha_{nm}|^{2} \leqslant \delta^{-1} \sum_{n,m} |\lambda_{nm} - a| |\alpha_{nm}|^{2} = \delta^{-1} ||u||_{E}^{2}, \qquad (3.2)$$

which implies the continuous embedding $E \hookrightarrow L^2(\Omega)$. Moreover, for $1 \leq r \leq 2$, the continuous embedding $L^2(\Omega) \hookrightarrow L^r(\Omega)$ implies that there exists a constant C = C(r) such that

$$\|u\|_{L^r(\Omega)} \leqslant C \|u\|_E, \quad 1 \leqslant r \leqslant 2. \tag{3.3}$$

Since $a \in (\rho_1, b^+ - b)$, then

$$e = b^+ - a > b. (3.4)$$

If $\lambda_{nm} > b$, we have

$$\lambda_{nm} - a \ge e > b, \quad n \in \mathbb{N}, \quad m \in \mathbb{Z}.$$
(3.5)

Define the energy functional

$$\Phi(u) = \frac{1}{2} \langle (L-a)u, u \rangle - \int_{\Omega} \rho \widetilde{F}(t, x, u) dt dx, \quad \forall u \in E,$$
(3.6)

where $\widetilde{F}(t, x, u) = \int_0^u \widetilde{f}(t, x, s) ds$. In addition, by (1.5) and the assumption on \widetilde{f} in theorem 2.4, it is easy to see $\widetilde{F}(t, x, u) \ge 0$ for any $u \in \mathbb{R}$.

Consequently, Φ is a C^1 functional on E and

$$\langle \Phi'(u), v \rangle = \langle (L-a)u, v \rangle - \int_{\Omega} \rho \tilde{f}(t, x, u) v dt dx, \quad \forall u, \quad v \in E.$$
(3.7)

Then u is a weak solution of problem (2.1)–(2.3) if and only if $\Phi'(u) = 0$. Therefore, the solutions of problem (2.1)–(2.3) are characterized as critical points of the functional Φ . Since \tilde{f} is C^1 , then Φ is a C^2 functional on E, and

$$\langle \Phi''(u)w,v\rangle = \langle (L-a)w,v\rangle - \int_{\Omega} \rho \frac{\partial \tilde{f}}{\partial u}(t,x,u)vwdtdx, \quad \forall u, \ v, \quad w \in E.$$

In particular,

$$\langle \Phi''(u)v,v\rangle = \langle (L-a)v,v\rangle - \int_{\Omega} \rho \frac{\partial \tilde{f}}{\partial u}(t,x,u)v^2 dt dx, \quad \forall u, \quad v \in E.$$
(3.8)

It is not difficult to see that Φ is neither bounded from above nor from below, which shows that we can not obtain the critical points of Φ by a simple minimization or maximization. Here we shall prove our main result by virtue of the following saddle point reduction technique developed by Amann [1], Castro and Lazer [9], and Arcoya and Costa [2], etc.

LEMMA 3.1. Let H be a real Hilbert space with the norm $\|\cdot\|_H$, $\Phi \in C^1(H, \mathbb{R})$, and H_1 , H_2 and H_3 be closed subset of H such that $H = H_1 \oplus H_2 \oplus H_3$. If there exists a constant $\gamma > 0$ satisfying

$$\langle \Phi'(u+w+v_1) - \Phi'(u+w+v_2), v_1 - v_2 \rangle \leq -\gamma \|v_1 - v_2\|_H^2, \ \forall v_1, v_2 \in H_1, u \in H_2, w \in H_3,$$

and

$$\langle \Phi'(u+w_1+v) - \Phi'(u+w_2+v), w_1 - w_2 \rangle \geq \gamma \|w_1 - w_2\|_H^2, \ \forall v \in H_1, u \in H_2, w_1, w_2 \in H_3.$$

Then

(i) There exists a unique continuous mapping $h: H_2 \to H_1 \oplus H_3$, such that

$$\Phi(u+h(u)) = \max_{v \in H_1} \min_{w \in H_3} \Phi(u+v+w) = \min_{w \in H_3} \max_{v \in H_1} \Phi(u+v+w);$$

(ii) Define $\widetilde{\Phi}(u) = \Phi(u+h(u))$ for any $u \in H_2$, then $\widetilde{\Phi} \in C^1(H_2, \mathbb{R})$, and

$$\langle \Phi'(u), v \rangle = \langle \Phi'(u+h(u)), v \rangle, \forall u, v \in H_2;$$

- (iii) If $u \in H_2$ is a critical point of $\widetilde{\Phi}$, then u + h(u) is a critical point of Φ . On the other hand, if u + v is a critical point of Φ , then v = h(u) and u is a critical point of $\widetilde{\Phi}$, where $u \in H_2$, $v \in H_1 \oplus H_3$;
- (iv) If Φ satisfies the Palais–Smale condition $(PS)_c$ at the level $c \in \mathbb{R}$, then the functional $\tilde{\Phi}$ also satisfies the $(PS)_c$ condition.

By an observation of lemma 3.1, it needs to decompose the working space E into suitable orthogonal subspaces. Noting $a, b \notin \Lambda(L)$ and $a > \rho_1$, the working space E can be decomposed into the following orthogonal subspaces

$$E_1 = \{ u \in E : \lambda_{nm} < a, \ n \in \mathbb{N}, m \in \mathbb{Z} \},\$$

$$E_2 = \{ u \in E : a < \lambda_{nm} < b, \ n \in \mathbb{N}, m \in \mathbb{Z} \},\$$

$$E_3 = \{ u \in E : \lambda_{nm} > b, \ n \in \mathbb{N}, m \in \mathbb{Z} \}.\$$

Thus we write $E = E_1 \oplus E_2 \oplus E_3$. Moreover, by $(a, b) \cap \Lambda(L) \neq \emptyset$ and lemma 2.3, we have $E_2 \neq \emptyset$ and dim $(E_2) < \infty$.

Existence of multiple periodic solutions to a semilinear wave equation 2593 For any $u \in E_1$, we write $u = \sum_{\lambda_{nm} < a} \alpha_{nm} \varphi_n \phi_m$, then

$$\langle (L-a)u, u \rangle = -\sum_{\lambda_{nm} < a} |\lambda_{nm} - a| |\alpha_{nm}|^2 = -||u||_E^2.$$
 (3.9)

For any $u \in E_2 \oplus E_3$, by a similar calculation we have

$$\langle (L-a)u, u \rangle = ||u||_E^2.$$
 (3.10)

With the help of (3.9) and (3.10), we have the following lemma.

LEMMA 3.2. Let the assumptions in theorem 2.4 hold. Then there exists a constant $\gamma > 0$ such that

$$\langle \Phi'(u+v) - \Phi'(u+w), v-w \rangle \leqslant -\gamma ||v-w||_E^2, \quad \forall u \in E_2 \oplus E_3, v, w \in E_1,$$

and

$$\langle \Phi'(u+v) - \Phi'(u+w), v-w \rangle \ge \gamma \|v-w\|_E^2, \quad \forall u \in E_1 \oplus E_2, \ v, \ w \in E_3.$$

Proof. For any $u, v, w \in E$ and $s \in \mathbb{R}$, we have

$$\langle \Phi'(u+v) - \Phi'(u+w), v-w \rangle = \int_0^1 \langle \Phi''(u+w+s(v-w))(v-w), v-w \rangle \mathrm{d}s.$$
(3.11)

From (3.8), we obtain

$$\langle \Phi''(u+w+s(v-w))(v-w), v-w \rangle = \langle (L-a)(v-w), v-w \rangle$$

$$-\int_{\Omega} \rho \frac{\partial \tilde{f}}{\partial u}(t, x, u+w+s(v-w))(v-w)^2 \mathrm{d}t \mathrm{d}x.$$
 (3.12)

The assumption on \tilde{f} in theorem 2.4 shows $(\partial \tilde{f}/\partial u) \ge 0$. For the case $v, w \in E_1$, $u \in E_2 \oplus E_3$, by (3.9) we have

$$\langle \Phi'(u+v) - \Phi'(u+w), v-w \rangle \leqslant - \|v-w\|_E^2.$$

For the case $v, w \in E_3, u \in E_1 \oplus E_2$, by (3.10) we have

$$\langle (L-a)(v-w), v-w \rangle = \|v-w\|_E^2.$$
 (3.13)

Moreover, since $0 \leq (\partial \tilde{f}/\partial u)(t, x, u) \leq e - \kappa$, with the aid of (3.5), a direct calculation yields

$$\int_{\Omega} \rho \frac{\partial \hat{f}}{\partial u} (v - w)^2 \mathrm{d}t \mathrm{d}x \leqslant (e - \kappa) \|v - w\|_{L^2(\Omega)}^2 \leqslant \frac{(e - \kappa)}{e} \|v - w\|_E^2.$$
(3.14)

Inserting (3.13), (3.14) into (3.12), from (3.11) we have

$$\langle \Phi'(u+v) - \Phi'(u+w), v-w \rangle \ge \frac{\kappa}{e} \|v-w\|_E^2.$$

By setting $\gamma = \min\{1, (\kappa/e)\}\)$, we arrive at the assertion.

Lemma 3.2 verifies all the conditions in lemma 3.1. Therefore, there exists a unique continuous mapping $h: E_2 \to E_1 \oplus E_3$ such that

$$\bar{\Phi}(u) = \Phi(u+h(u)) = \max_{v \in E_1} \min_{w \in E_3} \Phi(u+v+w) = \min_{w \in E_3} \max_{v \in E_1} \Phi(u+v+w). \quad (3.15)$$

Moreover, lemma 3.1 shows the reduced functional $\widetilde{\Phi} \in C^1(E_2, \mathbb{R})$, and $u \in E_2$ is a critical point of $\widetilde{\Phi}$ if and only if u + h(u) is a critical point of Φ . Thus, the critical points of Φ on the infinite dimensional space E are transformed into the ones of $\widetilde{\Phi}$ on the finite dimensional subspace E_2 . In what follows, we shall apply the variational methods to obtain critical points of the functional $\widetilde{\Phi}$ on E_2 .

4. Verification of the $(PS)_c$ condition

We will acquire the critical points of Φ via variational methods, thus it is needed to verify Φ satisfies $(PS)_c$ condition for any $c \in \mathbb{R}$. Furthermore, lemma 3.1 shows that it suffices to verify that Φ satisfies $(PS)_c$ condition, which means, any sequence $\{u_i\} \subset E$ satisfying $\Phi(u_i) \to c$ and $\Phi'(u_i) \to 0$ as $i \to \infty$ has a convergent subsequence for any $c \in \mathbb{R}$. To this goal, we need the following lemma which provides two estimates for the quadratic forms on different subspaces of E.

LEMMA 4.1. Let a, b satisfy the assumptions in theorem 2.4. Then there exist $\gamma_1, \gamma_2 > 0$ such that

$$\langle (L-a-b)u, u \rangle \leqslant -\gamma_1 \|u\|_E^2, \ \forall u \in E_1 \oplus E_2,$$

$$(4.1)$$

$$\langle (L-a-b)u, u \rangle \geqslant \gamma_2 \|u\|_E^2, \ \forall u \in E_3.$$

$$(4.2)$$

Proof. On the one hand, for $u \in E_1 \oplus E_2$, we write $u = \sum_{\lambda_{nm} < b} \alpha_{nm} \varphi_n \phi_m$. A direct calculation yields

$$\langle (L-a-b)u, u \rangle$$

= $\sum_{\lambda_{nm} < b} (\lambda_{nm} - a) |\alpha_{nm}|^2 - b \sum_{\lambda_{nm} < b} |\alpha_{nm}|^2$
 $\leqslant \sum_{\lambda_{nm} < a} (\lambda_{nm} - a) |\alpha_{nm}|^2 + \sum_{a < \lambda_{nm} < b} (\lambda_{nm} - a) |\alpha_{nm}|^2 - b \sum_{a < \lambda_{nm} < b} |\alpha_{nm}|^2.$

In virtue of $a < \lambda_{nm} < b$, we have $0 < \lambda_{nm} - a < b^{-}$. Thus, we have

$$\frac{1}{b^{-}}\sum_{a<\lambda_{nm}< b}|\lambda_{nm}-a||\alpha_{nm}|^{2} \leq \sum_{a<\lambda_{nm}< b}|\alpha_{nm}|^{2}.$$

Therefore,

$$\langle (L-a-b)u, u \rangle$$

$$\leqslant -\sum_{\lambda_{nm} < a} |\lambda_{nm} - a| |\alpha_{nm}|^2 - \left(\frac{b}{b^-} - 1\right) \sum_{a < \lambda_{nm} < b} |\lambda_{nm} - a| |\alpha_{nm}|^2$$

$$\leqslant -\gamma_1 \|u\|_E^2,$$

where $\gamma_1 = \min\{1, (b/b^-) - 1\}$. Noting $0 < b^- < b$, it follows $\gamma_1 > 0$.

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On the other hand, for $u \in E_3$, we write $u = \sum_{\lambda_{nm} > b} \alpha_{nm} \varphi_n \phi_m$. Observing $\lambda_{nm} > b$, from (3.5), we obtain

$$\frac{1}{e}\sum_{\lambda_{nm}>b}|\lambda_{nm}-a||\alpha_{nm}|^2 \ge \sum_{\lambda_{nm}>b}|\alpha_{nm}|^2.$$

Therefore,

$$\langle (L-a-b)u, u \rangle = \sum_{\lambda_{nm} > b} |\lambda_{nm} - a| |\alpha_{nm}|^2 - b \sum_{\lambda_{nm} > b} |\alpha_{nm}|^2$$
$$\geqslant \left(1 - \frac{b}{e}\right) \sum_{\lambda_{nm} > b} |\lambda_{nm} - a| |\alpha_{nm}|^2.$$

Let $\gamma_2 = 1 - (b/e)$, by (3.4), it follows $\gamma_2 > 0$. Therefore

$$\langle (L-a-b)u, u \rangle \ge \gamma_2 ||u||_E^2$$

The proof is completed.

LEMMA 4.2. Let the assumptions in theorem 2.4 hold. If $\{u_i\} \subset E$ satisfies $\Phi(u_i) \rightarrow c$ and $\Phi'(u_i) \rightarrow 0$ as $i \rightarrow \infty$, then there exists a constant $\widetilde{C} > 0$ independent of i such that $||u_i||_E \leq \widetilde{C}$.

Proof. Split $u_i = u_i^+ + u_i^-$, where $u_i^+ \in E_3$ and $u_i^- \in E_1 \oplus E_2$, i = 1, 2, ...

(i) For $u_i^+ \in E_3$, since $\Phi'(u_i) \to 0$ as $i \to \infty$, from (3.7), we have

$$o(1)\|u_i^+\|_E \ge \langle \Phi'(u_i), u_i^+ \rangle = \langle (L-a)u_i^+, u_i^+ \rangle - \int_{\Omega} \rho \tilde{f}(t, x, u_i)u_i^+ dt dx$$
$$= \langle (L-a-b)u_i^+, u_i^+ \rangle - \int_{\Omega} \rho (\tilde{f}(t, x, u_i) - bu_i)u_i^+ dt dx.$$
(4.3)

From (4.2), we obtain

$$\langle (L-a-b)u_i^+, u_i^+ \rangle \ge \gamma_2 ||u_i^+||_E^2.$$
 (4.4)

On the other side, the condition (1.6) shows that for any $\varepsilon > 0$, there exists a constant $C = C(\varepsilon) > 0$ such that

$$|\tilde{f}(t,x,u_i) - bu_i| < \varepsilon |u_i| + C.$$
(4.5)

Therefore, by Hölder inequality and (3.2), (3.3), a direct calculation yields

$$\left| \int_{\Omega} \rho(\tilde{f}(t, x, u_{i}) - bu_{i})u_{i}^{+} dt dx \right| \leq \varepsilon ||u_{i}^{+}||_{L^{2}(\Omega)} ||u_{i}||_{L^{2}(\Omega)} + C ||u_{i}^{+}||_{L^{1}(\Omega)}$$
$$\leq \frac{\varepsilon}{2\delta} ||u_{i}^{+}||_{E}^{2} + \frac{\varepsilon}{2\delta} ||u_{i}||_{E}^{2} + C ||u_{i}^{+}||_{E}, \quad (4.6)$$

for some constant C independent of i.

Inserting (4.4), (4.6) into (4.3), we obtain

$$\left(\gamma_2 - \frac{\varepsilon}{2\delta}\right) \|u_i^+\|_E^2 - \frac{\varepsilon}{2\delta} \|u_i\|_E^2 - C\|u_i^+\|_E \leqslant 0.$$

$$(4.7)$$

(ii) For $u_i^- \in E_1 \oplus E_2$, by (4.1), we have $\langle (L-a-b)u_i^-, u_i^- \rangle \leqslant -\gamma_1 ||u_i^-||_E^2$.

Moreover, noting that

$$\begin{split} o(1) \|u_i^-\|_E &\ge \langle -\Phi'(u_i), u_i^- \rangle \\ &= -\langle (L-a-b)u_i^-, u_i^- \rangle + \int_{\Omega} \rho(\tilde{f}(t, x, u_i) - bu_i)u_i^- \mathrm{d}t \mathrm{d}x, \end{split}$$

a similar calculation as the one in (4.6) yields

$$\left| \int_{\Omega} \rho(\tilde{f}(t,x,u_i) - bu_i) u_i^{-} \mathrm{d}t \mathrm{d}x \right| \leq \frac{\varepsilon}{2\delta} \|u_i^{-}\|_E^2 + \frac{\varepsilon}{2\delta} \|u_i\|_E^2 + C \|u_i^{-}\|_E.$$
(4.8)

Consequently,

$$\left(\gamma_1 - \frac{\varepsilon}{2\delta}\right) \|u_i^-\|_E^2 - \frac{\varepsilon}{2\delta} \|u_i\|_E^2 - C\|u_i^-\|_E \leqslant 0.$$

$$(4.9)$$

Let $\gamma_0 = \min\{\gamma_1, \gamma_2\}$. Since E_1 , E_2 and E_3 are orthogonal subspaces of E, we have $\|u_i\|_E^2 = \|u_i^+\|_E^2 + \|u_i^-\|_E^2$. Therefore, by the fact $\|u_i^+\|_E + \|u_i^-\|_E \leq \sqrt{2}\|u_i\|_E$, the sum of (4.7) and (4.9) yields

$$\left(\gamma_0 - \frac{3\varepsilon}{2\delta}\right) \|u_i\|_E^2 - C\|u_i\|_E \leqslant 0.$$
(4.10)

Taking $\varepsilon \in (0, (2\delta\gamma_0)/3)$ in (4.10), we know there exists a constant $\widetilde{C} > 0$ independent of *i* such that $||u_i||_E \leq \widetilde{C}$. We arrive at the result.

Since $E = E_1 \oplus E_2 \oplus E_3$, we can rewrite $E = E_1 \oplus E_1^{\perp}$ for simplicity, where

$$E_1^{\perp} = E_2 \oplus E_3 = \{ u \in E : \lambda_{nm} > a, \ n \in \mathbb{N}, m \in \mathbb{Z} \}.$$

Denote

$$E_0 = \{ u \in L^2(\Omega) : pn = q|m|, n \in \mathbb{N}, m \in \mathbb{Z} \}.$$

REMARK 4.3. Under the assumptions of theorem 2.4, the lemma 2.3 shows E_0 is an infinite dimensional subspace spanned by the eigenfunctions $\varphi_n \phi_m$ for pn = q|m|. Moreover, it is easy to see that $\dim(E_1^{\perp} \cap E_0) < \infty$ and $\dim(E_1 \cap E_0) = \infty$.

PROPOSITION 4.4. For $1 \leq r \leq 2$, the embedding

$$E \ominus E_0 \hookrightarrow L^r(\Omega), \tag{4.11}$$

is compact.

Proof. For $u \in E \ominus E_0$, it can be expanded as $u = \sum_{n,m} \alpha_{nm} \varphi_n \phi_m$ for $pn \neq q|m|$.

Existence of multiple periodic solutions to a semilinear wave equation 2597 The fact $||u||_E = (\sum_{n,m} |\lambda_{nm} - a| |\alpha_{nm}|^2)^{1/2}$ shows that the mapping

$$I_0: u = \sum_{n,m} \alpha_{nm} \varphi_n \phi_m \mapsto \{ |\lambda_{nm} - a|^{1/2} \alpha_{nm} \}$$

is continuous from $E \ominus E_0$ to l^2 .

Since $|\lambda_{nm} - a| \to \infty$ as $n, m \to \infty$, it follows that the mapping

$$I_1: \{|\lambda_{nm} - a|^{1/2}\alpha_{nm}\} \mapsto \{\alpha_{nm}\}$$

is compact from l^2 to l^2 .

Since $\varphi_n \phi_m$ is a complete orthonormal sequence of $L^2(\Omega)$, then the mapping

$$I_2: \{\alpha_{nm}\} \mapsto u = \sum_{n,m} \alpha_{nm} \varphi_n \phi_m$$

is continuous from l^2 to $L^2(\Omega)$.

Consequently, the mapping

$$I_2I_1I_0: E \ominus E_0 \to L^2(\Omega)$$

is compact. Furthermore, for $1 \leq r \leq 2$, the continuous embedding $L^2(\Omega) \hookrightarrow L^r(\Omega)$ implies that the embedding $E \ominus E_0 \hookrightarrow L^r(\Omega)$ is compact. \Box

LEMMA 4.5. Let the assumptions of theorem 2.4 hold. Then Φ satisfies the $(PS)_c$ condition for any $c \in \mathbb{R}$.

Proof. For any $c \in \mathbb{R}$, assume $\{u_i\} \subset E$ satisfies $\Phi(u_i) \to c$ and $\Phi'(u_i) \to 0$ as $i \to \infty$. Since E is a Hilbert space, by lemma 4.2, we have $u_i \rightharpoonup u$ as $i \to \infty$ for some $u \in E$. Decompose $u_i = v_i + y_i + w_i + z_i$ and u = v + y + w + z, where v, y, w, z are the weak limits of $\{v_i\}, \{y_i\}, \{w_i\}, \{z_i\}$ respectively, and $v_i, v \in E_1^{\perp} \ominus E_0, y_i, y \in E_1^{\perp} \cap E_0, w_i, w \in E_1 \ominus E_0, z_i, z \in E_1 \cap E_0$.

(i) For $v_i, v \in E_1^{\perp} \oplus E_0$, from (3.10) we have

$$\|v_i - v\|_E^2 = \langle (L-a)(v_i - v), v_i - v \rangle$$

= $\langle (L-a)v_i, v_i - v \rangle - \langle (L-a)v, v_i - v \rangle.$ (4.12)

In virtue of $v_i \rightarrow v$ in E and $E \rightarrow L^2(\Omega)$, we have $v_i \rightarrow v$ in $L^2(\Omega)$ along with a subsequence of $\{v_i\}$. In fact, by proposition 4.4, we also have $v_i \rightarrow v$ in $L^2(\Omega)$. We still use $\{v_i\}$ to denote the subsequence for convenience. Thus,

it follows

$$\langle (L-a)v, v_i - v \rangle \to 0, \text{ as } i \to \infty.$$
 (4.13)

On the other hand, noting $v_i, v \in E_1^{\perp} \ominus E_0$ and $u_i = v_i + y_i + w_i + z_i$, we have $u_i - v_i \in (E_1^{\perp} \ominus E_0)^{\perp}$. Thus

$$\langle (L-a)(u_i-v_i), v_i-v \rangle = 0.$$

By (3.7), we have

$$\langle (L-a)v_i, v_i - v \rangle = \langle (L-a)u_i, v_i - v \rangle$$
$$= \langle \Phi'(u_i), v_i - v \rangle + \int_{\Omega} \rho \tilde{f}(t, x, u_i)(v_i - v) dt dx.$$
(4.14)

Since $\Phi'(u_i) \to 0$ as $i \to \infty$, we have

$$\langle \Phi'(u_i), v_i - v \rangle \to 0, \quad \text{as } i \to \infty.$$
 (4.15)

By (4.5) and Hölder inequality, a direct calculation yields

$$\left| \int_{\Omega} \rho \tilde{f}(t, x, u_i)(v_i - v) \mathrm{d}t \mathrm{d}x \right|$$

$$\leq (b + \varepsilon) \|u_i\|_{L^2(\Omega)} \|v_i - v\|_{L^2(\Omega)} + C \|v_i - v\|_{L^1(\Omega)}.$$

In virtue of $v_i \to v$ in $L^2(\Omega)$ and the continuous embedding $L^2(\Omega, \rho) \hookrightarrow L^1(\Omega)$, we obtain $v_i \to v$ in $L^1(\Omega)$. Therefore,

$$\left| \int_{\Omega} \rho \tilde{f}(t, x, u_i)(v_i - v) dt dx \right| \to 0, \quad \text{as } i \to \infty.$$
(4.16)

Inserting (4.15), (4.16) into (4.14), we have

$$\langle (L-a)v_i, v_i-v \rangle \to 0, \text{ as } i \to \infty.$$
 (4.17)

Finally, substituting (4.17), (4.13) into (4.12), we have

$$\|v_i - v\|_E \to 0, \quad \text{as } i \to \infty.$$
(4.18)

- (ii) For $y_i, y \in E_1^{\perp} \cap E_0$, since dim $(E_1^{\perp} \cap E_0) < \infty$ and $y_i \rightharpoonup y$ in E, then there exists a subsequence of $\{y_i\}$ which strongly converges to y in E. The subsequence is still denoted by $\{y_i\}$.
- (iii) For $w_i, w \in E_1 \ominus E_0$, by (3.7) and (3.9) we obtain

$$\begin{split} \|w_i - w\|_E^2 &= -\langle (L-a)(w_i - w), w_i - w \rangle \\ &= -\langle \Phi'(u_i), w_i - w \rangle - \int_{\Omega} \rho \tilde{f}(t, x, u_i)(w_i - w) \mathrm{d}t \mathrm{d}x \\ &+ \langle (L-a)w, w_i - w \rangle. \end{split}$$

A similar calculation as the one in (i) yields

$$||w_i - w||_E \to 0, \quad \text{as } i \to \infty.$$
 (4.19)

(iv) Since $z_i \in E_1 \cap E_0$, then the compact embedding (4.11) is invalid. In addition, since dim $(E_1 \cap E_0) = \infty$, we can not extract a strong convergence subsequence of $\{z_i\}$ similar to $\{y_i\}$. In what follows, with the aid of the monotone method, we prove $z_i \to z$ as $i \to \infty$ in E.

Since $\Phi'(u_i) \to 0$ and $z_i \rightharpoonup z$ in E, we have

$$||z_i - z||_E^2 = -\langle (L - a)(z_i - z), z_i - z \rangle$$

= $-\langle \Phi'(u_i), z_i - z \rangle - \int_{\Omega} \rho \tilde{f}(t, x, u_i)(z_i - z) dt dx + \langle (L - a)z, z_i - z \rangle$
 $\leqslant - \int_{\Omega} \rho \tilde{f}(t, x, u_i)(z_i - z) dt dx + o(1).$ (4.20)

Denote $\tilde{f}(u_i) = \tilde{f}(t, x, u_i)$ for simplicity. We rewrite $\int_{\Omega} \rho \tilde{f}(t, x, u_i)(z_i - z) dt dx$ in the form of inner product and decompose it as follows

$$\int_{\Omega} \rho \tilde{f}(t, x, u_i)(z_i - z) dt dx$$

$$= \langle \tilde{f}(u_i), z_i - z \rangle$$

$$= \langle \tilde{f}(u_i) - \tilde{f}(\tilde{u}_i + z), z_i - z \rangle + \langle \tilde{f}(\tilde{u}_i + z) - \tilde{f}(u), z_i - z \rangle$$

$$+ \langle \tilde{f}(u), z_i - z \rangle,$$
(4.21)
(4.21)

where $\tilde{u}_i = v_i + y_i + w_i$.

Since \tilde{f} is increasing in u, then

$$\langle \tilde{f}(u_i) - \tilde{f}(\tilde{u}_i + z), z_i - z \rangle \ge 0.$$
(4.23)

To continue the discussion, by (4.5), we have $\tilde{f}: u \mapsto \tilde{f}(t, x, u)$ is continuous from $L^2(\Omega)$ to $L^2(\Omega)$. Moreover, from the proof of (i)–(iii), we have $\tilde{u}_i \to \tilde{u}$ in E, where $\tilde{u} = v + y + w$. The inequality (3.2) shows $\tilde{u}_i \to \tilde{u}$ in $L^2(\Omega)$. Thus

$$\langle \tilde{f}(\tilde{u}_i+z) - \tilde{f}(u), z_i-z \rangle \to 0, \quad \text{as } i \to \infty.$$
 (4.24)

Consequently, since $z_i \rightharpoonup z$ in $L^2(\Omega)$, by (4.20), (4.23) and (4.24), we have

$$|z_i - z||_E \to 0$$
, as $i \to \infty$.

We arrive at the result.

5. Bounds of the reduced functional

The following three lemmas are concerned with the bounds of the reduced functional $\widetilde{\Phi}$ and play an important role in the proof of theorem 2.4.

LEMMA 5.1. Let the assumptions of theorem 2.4 hold. Then

- (i) there exists a constant M > 0, such that $\widetilde{\Phi}(u) < M$, $\forall u \in E_2$;
- (ii) there exists a constant $\widetilde{R} > 0$ such that $\widetilde{\Phi}(u) \leq 0$ for $u \in E_2$ with $||u||_E \geq \widetilde{R}$.

Proof. For $u \in E_2$, from (3.15), we have

$$\tilde{\Phi}(u) = \min_{w \in E_3} \max_{v \in E_1} \Phi(u + v + w) \leqslant \max_{v \in E_1} \Phi(u + v),$$

where

$$\Phi(u+v) = \frac{1}{2} \langle (L-a-b)(u+v), u+v \rangle - \int_{\Omega} \rho \left(\widetilde{F}(t,x,u+v) - \frac{b}{2}(u+v)^2 \right) \mathrm{d}t \mathrm{d}x.$$
(5.1)

By (4.5), it follows

$$|\widetilde{F}(t, x, u+v) - \frac{b}{2}(u+v)^2| \le \varepsilon |u+v|^2 + C|u+v|.$$
(5.2)

Inserting (5.2) into (5.1), by (4.1), it follows

$$\begin{split} \Phi(u+v) \leqslant -\frac{\gamma_1}{2} \|u+v\|_E^2 + \int_{\Omega} \rho(\varepsilon |u+v|^2 + C|u+v|) \mathrm{d}t \mathrm{d}x \\ \leqslant -\frac{\gamma_1}{2} \|u+v\|_E^2 + \varepsilon \|u+v\|_{L^2(\Omega)}^2 + C\|u+v\|_{L^1(\Omega)}, \end{split}$$

for some constant C depending on ε .

Taking $\varepsilon = (\delta \gamma_1/4)$ in above inequality, by (3.2) and (3.3), we obtain

$$\Phi(u+v) \leqslant -\frac{\gamma_1}{4} \|u+v\|_E^2 + C\|u+v\|_E.$$
(5.3)

Therefore, the estimate (5.3) shows there exists M > 0 such that $\Phi(u + v) \leq M$ and the assertion (i) is proved.

By the fact $||u + v||_E^2 = ||u||_E^2 + ||v||_E^2$ and $||u + v||_E \leq ||u||_E + ||v||_E$, from (5.3), we have

$$\Phi(u+v) \leqslant -\frac{\gamma_1}{4} \|u\|_E^2 + C \|u\|_E + \left(-\frac{\gamma_1}{4} \|v\|_E^2 + C \|v\|_E\right)$$

$$\leqslant -\frac{\gamma_1}{4} \|u\|_E^2 + C \|u\|_E + C_0,$$

where $C_0 = \max_{s \ge 0} \{-(\gamma_1/4)s^2 + Cs\}$. Thus, there exists a constant $\widetilde{R} > 0$ such that $\Phi(u+v) \le 0$ for $||u||_E \ge \widetilde{R}$. We conclude the assertion (ii).

LEMMA 5.2. Let the assumptions of theorem 2.4 hold. Then for any $R_1 > 0$, there exists a constant τ_1 depending on R_1 such that

$$\Phi(u) \geqslant \tau_1, \quad \forall u \in E_2 \cap B_{R_1},$$

where $B_{R_1} = \{ u \in E : ||u||_E < R_1 \}.$

Existence of multiple periodic solutions to a semilinear wave equation 2601 Proof. For $u \in E_2$, by (3.15), we have

$$\widetilde{\Phi}(u) = \max_{v \in E_1} \min_{w \in E_3} \Phi(u + v + w) \ge \min_{w \in E_3} \Phi(u + w),$$

where

$$\Phi(u+w) = \frac{1}{2} \langle (L-a-b)(u+w), u+w \rangle$$
$$-\int_{\Omega} \rho\left(\widetilde{F}(t,x,u+w) - \frac{b}{2}(u+w)^2\right) dt dx.$$
(5.4)

Since E_2 and E_3 are orthogonal subspaces of E, we have

$$\langle (L-a-b)(u+w), u+w \rangle = \langle (L-a-b)u, u \rangle + \langle (L-a-b)w, w \rangle.$$

By (4.2), we have $\langle (L-a-b)w, w \rangle \ge \gamma_2 ||w||_E^2$. Moreover, noting $\langle (L-a)u, u \rangle = ||u||_E^2$, by (3.2), we obtain

$$|\langle (L-a-b)u, u \rangle| \leq ||u||_E^2 + b||u||_{L^2(\Omega)}^2 \leq \left(1 + \frac{b}{\delta}\right) ||u||_E^2 = C_1 ||u||_E^2,$$

where $C_1 = 1 + (b/\delta)$. Thus, we have

$$\langle (L-a-b)(u+w), u+w \rangle \ge \gamma_2 ||w||_E^2 - C_1 ||u||_E^2.$$
 (5.5)

On the other hand, a similar calculation as the one in (5.2) yields

$$\left|\widetilde{F}(t,x,u+w) - \frac{b}{2}(u+w)^2\right| \leqslant \varepsilon |u+w|^2 + C|u+w|.$$
(5.6)

Inserting (5.5), (5.6) into (5.4) and taking $\varepsilon = (\delta \gamma_2/4)$, from (3.2) and (3.3), we have

$$\begin{split} \Phi(u+w) &\ge \frac{\gamma_2}{2} \|w\|_E^2 - \frac{C_1}{2} \|u\|_E^2 - \frac{\delta\gamma_2}{4} \|u+w\|_{L^2(\Omega)}^2 - C\|u+w\|_{L^1(\Omega)} \\ &\ge \frac{\gamma_2}{2} \|w\|_E^2 - \frac{C_1}{2} \|u\|_E^2 - \frac{\gamma_2}{4} (\|u\|_E^2 + \|w\|_E^2) - C(\|u\|_E + \|w\|_E) \\ &\ge - \left(\frac{C_1}{2} + \frac{\gamma_2}{4}\right) \|u\|_E^2 - C\|u\|_E + C_2, \end{split}$$

where $C_2 = \min_{s \ge 0} \{ (\gamma_2/4) s^2 - Cs \}.$

For any $R_1 > 0$, the above estimate shows $\widetilde{\Phi}(u) \ge \tau_1 = -((C_1/2) + (\gamma_2/4))R_1^2 - CR_1 + C_2$ for any $||u||_E < R_1$. The proof is completed.

LEMMA 5.3. Let the assumptions of theorem 2.4 hold. Then there exist two constants $\tau_2 > 0$ and $R_2 > 0$ such that $\widetilde{\Phi}(u) \ge \tau_2$, for any $u \in E_2$ with $||u||_E = R_2$.

Proof. Firstly, for any $u \in E_2$, $w \in E_3$, by (3.6), we have

$$\Phi(u+w) = \frac{1}{2} \langle (L-a)(u+w), u+w \rangle - \int_{\Omega} \rho \widetilde{F}(t, x, u+w) dt dx.$$
 (5.7)

Since E_2 and E_3 are orthogonal subspaces of E, from (3.10), we obtain

$$\langle (L-a)(u+w), u+w \rangle = ||u||_E^2 + ||w||_E^2.$$
 (5.8)

In addition, it is easy to see

$$\int_0^1 \int_0^1 s \frac{\partial \tilde{f}}{\partial \xi} (u + s\theta w) w^2 d\theta ds = \int_0^1 w \tilde{f}(u + sw) ds - \tilde{f}(u) w$$
$$= \tilde{F}(u + w) - \tilde{F}(u) - \tilde{f}(u) w,$$

where $\tilde{f}(\xi) = \tilde{f}(t, x, \xi)$ and $\tilde{F}(\xi) = \tilde{F}(t, x, \xi)$. Thus,

$$\widetilde{F}(u+w) = \int_0^1 \int_0^1 s \frac{\partial \widetilde{f}}{\partial \xi} (u+s\theta w) w^2 \mathrm{d}\theta \mathrm{d}s + \widetilde{f}(u)w + \widetilde{F}(u).$$
(5.9)

In what follows, we estimate the upper bound of $\int_{\Omega} \rho \widetilde{F}(u+w) dt dx$. The equation (5.9) shows that it needs to estimate the upper bounds of the following three terms:

$$\int_{\Omega} \rho\left(\int_{0}^{1} \int_{0}^{1} s \frac{\partial \tilde{f}}{\partial \xi}(u+s\theta w) w^{2} \mathrm{d}\theta \mathrm{d}s\right) \mathrm{d}t \mathrm{d}x, \quad \int_{\Omega} \rho \tilde{f}(u) w \mathrm{d}t \mathrm{d}x, \quad \int_{\Omega} \rho \widetilde{F}(u) \mathrm{d}t \mathrm{d}x.$$

(i) For $w \in E_3$, the fact $|\lambda_{nm} - a| \ge e$ shows $||w||_{L^2(\Omega)}^2 \le (1/e) ||w||_E^2$. Observing $0 \le (\partial \tilde{f}/\partial u)(t, x, u) \le e - \kappa$, we obtain

$$0 \leqslant \int_{\Omega} \rho\left(\int_{0}^{1} \int_{0}^{1} s \frac{\partial \tilde{f}}{\partial \xi} (u + s\theta w) w^{2} \mathrm{d}\theta \mathrm{d}s\right) \mathrm{d}t \mathrm{d}x \leqslant \frac{e - \kappa}{2e} \|w\|_{E}^{2}.$$
(5.10)

(ii) Fix r > 1, by the assumptions (1.5) and (1.6), we obtain that for any $\varepsilon > 0$, there exists a constant $C = C(\varepsilon, r) > 0$ such that

$$|\tilde{f}(u)| \leq \varepsilon |u| + C|u|^r, \quad \forall \ (t, x, u) \in \Omega \times \mathbb{R}.$$
(5.11)

Therefore, a direct calculation yields

$$\left| \int_{\Omega} \rho \tilde{f}(u) w dt dx \right| \leq \varepsilon ||u||_{L^{2}(\Omega)} ||w||_{L^{2}(\Omega)} + C ||u||_{L^{2r}(\Omega)}^{r} ||w||_{L^{2}(\Omega)}$$
$$\leq \frac{\varepsilon}{2} ||u||_{L^{2}(\Omega)}^{2} + C ||u||_{L^{2r}(\Omega)}^{2r} + \varepsilon ||w||_{L^{2}(\Omega)}^{2}.$$

Since dim $(E_2) < \infty$, then all norms of E_2 are equivalent. Thus there exists a constant C > 0 such that $||u||_{L^{2r}(\Omega)}^{2r} \leq C ||u||_E^{2r}$. Since $||w||_{L^2(\Omega)}^2 \leq (1/e) ||w||_E^2$, by (3.2), we have

$$\left| \int_{\Omega} \rho \tilde{f}(u) w \mathrm{d}t \mathrm{d}x \right| \leq \frac{\varepsilon}{2\delta} \|u\|_{E}^{2} + C \|u\|_{E}^{2r} + \frac{\varepsilon}{e} \|w\|_{E}^{2}.$$
(5.12)

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(iii) Since dim $(E_2) < \infty$, there exists a constant C > 0 such that $||u||_{L^{r+1}(\Omega)}^{r+1} \leq C||u||_{E}^{r+1}$. Thus, by (3.2) and (5.11), a direct calculation yields

$$\left| \int_{\Omega} \rho \widetilde{F}(u) \mathrm{d}t \mathrm{d}x \right| \leqslant \frac{\varepsilon}{2\delta} \|u\|_{E}^{2} + C \|u\|_{E}^{r+1}, \tag{5.13}$$

for some constant C depending on ε and r.

Therefore, the sum of (5.10), (5.12) and (5.13) yields

$$\left| \int_{\Omega} \rho \widetilde{F}(u+w) \mathrm{d}t \mathrm{d}x \right| \leqslant \frac{\varepsilon}{\delta} \|u\|_{E}^{2} + C \|u\|_{E}^{r+1} + C \|u\|_{E}^{2r} + \frac{e-\kappa+2\varepsilon}{2e} \|w\|_{E}^{2}.$$
(5.14)

Consequently, inserting (5.8) and (5.14) into (5.7) and taking $\varepsilon = \min\{(\delta/4), (\kappa/4)\}$, we have

$$\begin{split} \Phi(u+w) \geqslant \left(\frac{1}{2} - \frac{\varepsilon}{\delta}\right) \|u\|_{E}^{2} - C\|u\|_{E}^{r+1} - C\|u\|_{E}^{2r} + \left(\frac{1}{2} - \frac{e - \kappa + 2\varepsilon}{2e}\right) \|w\|_{E}^{2} \\ &= \left(\frac{1}{2} - \frac{\varepsilon}{\delta}\right) \|u\|_{E}^{2} - C\|u\|_{E}^{r+1} - C\|u\|_{E}^{2r} + \left(\frac{\kappa - 2\varepsilon}{2e}\right) \|w\|_{E}^{2}. \\ &\geqslant \frac{1}{4} \|u\|_{E}^{2} - C\|u\|_{E}^{r+1} - C\|u\|_{E}^{2r}. \end{split}$$
(5.15)

Finally, since r > 1, then $\phi(s) = (1/4)s^2 - Cs^{r+1} - Cs^{2r}$ attains the local minimum at s = 0. Therefore there exist two constants $R_2 > 0$ and $\tau_2 > 0$ such that $\phi(R_2) \ge \tau_2$. Since $\tilde{\Phi}(u) \ge \min_{w \in E_3} \Phi(u+w)$, from (5.15), we obtain $\tilde{\Phi}(u) \ge \tau_2$ with $\|u\|_E = R_2$.

6. Proof of theorem 2.4

With the above lemmas in hand, we give the proof of theorem 2.4.

Proof. Let R_2 and τ_2 be the constants in lemma 5.3 and $B_{R_2} = \{u \in E_2 : ||u||_E < R_2\}.$

Firstly, by lemma 5.1, we know $\widetilde{\Phi}$ is bounded from above. Let $\beta_1 = \sup_{u \in E_2} \widetilde{\Phi}(u)$. By lemmas 3.1 and 4.5, $\widetilde{\Phi}$ satisfies $(PS)_{\beta_1}$ condition. Thus, there exists a critical point $u_1 \in E_2$ such that $\widetilde{\Phi}(u_1) = \beta_1$ and $\widetilde{\Phi}'(u_1) = 0$.

Secondly, recall $\widetilde{F}(t,r,u) \ge 0$ for any $u \in \mathbb{R}$, then for any $v \in E_1$, from (3.9), we obtain

$$\Phi(v) = \frac{1}{2} \langle (L-a)v, v \rangle - \int_{\Omega} \rho \widetilde{F}(t, x, v) dt dx \leqslant 0.$$

Therefore, we have

$$\widetilde{\Phi}(0) = \min_{w \in E_3} \max_{v \in E_1} \Phi(v+w) \leqslant \max_{v \in E_1} \Phi(v) \leqslant 0.$$
(6.1)

Noting that $0 \in B_{R_2}$ and taking $R_1 = R_2$ in lemma 5.2, in virtue of lemmas 5.2, 5.3 and $\tilde{\Phi}(0) \leq 0$, we obtain the reduced functional $\tilde{\Phi}$ that attains its infimum in B_{R_2} . Let $\beta_2 = \inf_{u \in B_{R_2}} \widetilde{\Phi}(u)$, then lemmas 3.1 and 4.5 show that $\widetilde{\Phi}$ satisfies the $(PS)_{\beta_2}$ condition. Thus, there exists $u_2 \in E_2$ such that $\widetilde{\Phi}'(u_2) = 0$ and $\widetilde{\Phi}(u_2) = \beta_2$.

Now we claim u_1 and u_2 are two different critical points in E_2 . Observing that $0 \in B_{R_2}$, the following inequality

$$\beta_2 = \inf_{u \in B_{R_2}} \widetilde{\Phi}(u) \leqslant \widetilde{\Phi}(0) \leqslant 0 < \tau_2 \leqslant \inf_{\|u\|_E = R_2} \widetilde{\Phi}(u) \leqslant \sup_{u \in E_2} \widetilde{\Phi}(u) = \beta_1 \quad (6.2)$$

shows $u_1 \neq u_2$.

If Φ possesses another local maximum point which is different from u_1 , then there exist at least three critical points of $\tilde{\Phi}$.

Otherwise, if u_1 is the unique local maximum point of Φ . We shall prove that there exists a critical point of mountain-pass type which is different from u_1 and u_2 .

Taking $u_0 \in E_2$ with $||u_0||_E = 1$, by lemma 5.1, there exists $\widetilde{R} > R_2$ such that $\widetilde{\Phi}(\widetilde{R}u_0) \leq 0$. Moreover, one of the fact holds: for any $s \in [0, 1]$, either $s\widetilde{R}u_0 \neq u_1$ or $-s\widetilde{R}u_0 \neq u_1$.

Case 1: if $s\widetilde{R}u_0 \neq u_1$ for any $s \in [0, 1]$, by lemma 5.3 and (6.1), we obtain

$$\max\{\widetilde{\Phi}(0), \widetilde{\Phi}(\widetilde{R}u_0)\} \leqslant 0 < \tau_2 \leqslant \inf_{\|u\|_E = R_2} \widetilde{\Phi}(u).$$

Let

$$c^{+} = \inf_{g \in \Delta^{+}} \max_{s \in [0,1]} \widetilde{\Phi}(g(s)),$$

where $\Delta^+ = \{g \in C([0,1], E_2) : g(0) = 0, g(1) = \widetilde{R}u_0\}$. In addition, we have $\widetilde{\Phi}$ which is C^1 functional and satisfies $(PS)_{c^+}$ condition. By mountain pass lemma [11], we obtain a critical point u_3^+ satisfying $\widetilde{\Phi}(u_3^+) = c^+$ and $\widetilde{\Phi}'(u_3^+) = 0$. Obviously, $c^+ \ge \tau_2 > 0$.

It is easy to see that $s\widetilde{R}u_0 \in \Delta^+$ for any $s \in [0, 1]$. Therefore, the assumption $s\widetilde{R}u_0 \neq u_1$ implies

$$c^+ \leqslant \max_{s \in [0,1]} \widetilde{\Phi}(s\widetilde{R}u_0) < \sup_{u \in E_2} \widetilde{\Phi}(u) = \beta_1.$$

Thus, the following inequalities

$$\beta_2 = \inf_{u \in B_{R_2}} \widetilde{\Phi}(u) \leqslant 0 < \tau_2 \leqslant c^+ < \sup_{u \in E_2} \widetilde{\Phi}(u) = \beta_1$$

show $u_1 \neq u_2 \neq u_3^+$.

Case 2: similarly, if $-s\widetilde{R}u_0 \neq u_1$ for any $s \in [0, 1]$, we have

$$c^{-} = \inf_{g \in \Delta^{-}} \max_{s \in [0,1]} \widetilde{\Phi}(g(s))$$

is the critical value of $\widetilde{\Phi}$, where $\Delta^- = \{g \in C([0,1], E_2) : g(0) = 0, g(1) = -\widetilde{R}u_0\}$. Therefore, there exists $u_3^- \in E_2$ satisfying $\widetilde{\Phi}(u_3^-) = c^-$ and $\widetilde{\Phi}'(u_3^-) = 0$. Furthermore, we have $u_1 \neq u_2 \neq u_3^-$.

Consequently, by lemma 3.1, we obtain the functional Φ that has at least three critical points. The proof is completed.

Funding

This work was supported by NSFC Grants (nos. 11671071, 11701077 and 11871140), the Fundamental Research Funds for the Central Universities of Jilin University (no. 2017TD-18), and the Special Funds of Provincial Industrial Innovation in Jilin Province (no. 2017C028-1).

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