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Simulation of microwave optical links and proof of noise figure lower than electrical losses

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The operation of a microwave photonic link is thoroughly investigated both theoretically and experimentally. To this aim, we have developed a simulation tool based on an accurate physical model embedded in a radio frequency (RF) chain simulator. The theoretical predictions are tested on an intensity modulation-direct detection (IMDD) link we have specifically developed to this purpose. Our simulation tool takes into account both optical and electrical characteristics of the link components including the laser dynamics and impedance matching networks. It thus enables an accurate understanding of the different physical and electrical phenomena governing the link's performances even under unusual operation conditions. Specifically, we were able to isolate an unusual behavior and to confirm it experimentally. It is thereby clear that the noise figure of a microwave optical link can be lower than the electrical losses, such as a mismatched output passive electrical network. This state is reached when the optical losses are high enough and when the link's output impedance is mismatched, too.

Keywords: Microwave photonics, Modeling, Simulation and characterizations of devices and circuits, Microwave noise power measurements

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I. INTRODUCTION

The development of accurate models describing radio frequency (RF) optical links is of great importance for designers. Such accurate models should benefit to several fields of applications such as radio over fiber, analog remote links and networks as well as RF photonics architectures for defense applications. Current commercially available simulation tools [1] use time domain models that are well adapted to digital optical links, but show some limitations as far as analog optical links are considered. They fail, for instance, to accurately model the linear dynamic range and noise contributions. Moreover, they cannot be interfaced with available RF simulation softwares, like advanced design system, since these softwares use frequency domain models to determine standard characteristics as transfer function, 1 dB compression point, IP3 (third-order intercept point), and noise figure of the RF chain. In order to tackle this problem, several studies have already been published. Among these studies Garenaux et al. [2] and Bdeoui et al. [3] report on accurate evaluation of RF equivalent parameters including

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Corresponding author: A.-L. Billabert Email: anne-laure.billabert@cnam.fr close to carrier phase noise. Analytical models for the noise figure have also been proposed in [3–6]. We hereafter investigate theoretically and validate experimentally the operation of a directly modulated link with a focus on a peculiar behavior occurring when the link's output impedance is mismatched.

In the first section we present each component developed electrical model of the studied link, mainly its electrical-to-optical (E/O) and optical-to-electrical (O/E) signal conversion parts: the directly modulated distributed feedback (DFB) laser and the p.i.n. photodetector. We then describe the dynamic response of the link in the second section. In the third section, the simulation results, in terms of output noise power (ONP) and noise figure, are compared with the experimental results in the case of a full matched optical link. In particular, the last section focuses on noise characteristic simulations, compared with measurements, under high optical losses condition. Unusual noise behavior is predicted and confirmed experimentally.

II. SIMULATION OF OPTICAL MICROWAVE LINK ELEMENTS

The developed simulation tool is based on an accurate physical model embedded in a RF chain simulator. A more detailed description of this simulation tool can be found in [3]. As expressed before, the analyzed link consists of a laser diode, an optical fiber, and a photodetector. The simplified circuit model for each link element is reported in Fig. 1.

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Fig. 1. Simplified models of the direct modulated photonic link elements, including matching circuits and current noise sources.

A) Laser diode

The laser source is a DFB laser diode oscillating at 1.55μ m. The laser is simulated using a non linear model [3] based on the single-mode rate equations governing the carrier and photon evolution inside the laser cavity [7].

The generated optical carrier is directly intensity modulated through the injection current. At the output of the laser, the modulated optical power is photodetected by an ideal detector, i.e., perfectly matched for all operating frequencies and exhibiting 100% quantum efficiency, then providing an electrical current proportional to the optical power envelope. The relative intensity noise (RIN) is taken into account by including Langevin noise forces into the rate equations of the laser.

To model the laser, it is necessary to precisely assess its static response and also its intrinsic dynamic response. A straightforward method consists in measuring the RF gain of the link. It is important to notice that the measured global forward transmission S-parameter (S_{21}) of the link includes the connecting circuits of the laser diode. So, to obtain the exact intrinsic dynamic response of the laser one needs to extract from the global S_{21} these extrinsic circuits' responses. Difference between the measured and the theoretical transfer functions include the mismatching impedance information, which is equal to the access circuit transfer function. As the latter is computed rather than intrinsically measured, another method is preferred. This method relies on the measurement of the laser excess noise versus frequency.

Indeed, the noise power density of the laser is a signature of its transfer function since the spontaneous emission follows the laser dynamical response and is independent of RF access matching. Thus, the first step consists in measuring the laser RIN of the laser. Some of these measurements are presented in Fig. 2 for different bias currents.



Fig. 2. Measurements and simulations of laser RIN for different optical powers.

As the laser model relies on rate equations, all appropriate parameters correspond to physical data such as: the active layer volume, carrier and photon lifetimes, the differential gain coefficient, the carrier density at transparency, the confinement factor, the gain compression factor, and the spontaneous recombination factor. Some physical parameters can be easily estimated like the active layer length that can be obtained from the free spectral range of the laser. Other physical parameters have to be estimated using the rate equations of the laser [1] owing to knowledge of static and dynamic system parameters.

The approach consists in defining system parameters such as the optical wavelength, the threshold current, and the differential quantum efficiency, which characterize the static response (DC) of the laser, on the one hand, and the bandwidth and the relaxation oscillation frequency for a given bias current, which characterize its dynamic response, on the other hand. In our case the DC characteristic of the DFB laser is measured to be 47 mA threshold current and 0.29 W/A responsivity. As for the alternating current (AC) characteristic, for bias current of 117 mA (i.e. 13 dBm optical power), the relaxation oscillation frequency and the -3 dB bandwidth are, respectively, equal to 3.4 and 1.2 GHz.

The noise current has two contributions, thermal noise and RIN:

$$\begin{split} \langle i_{nG}^2 \rangle &= \frac{4 \, kTB}{R_o}, \\ \langle i_{nL}^2 \rangle &= \frac{(I_{bias} - I_{th})^2 RIN(\omega) B \eta_L}{2}. \end{split} \tag{1}$$

The different variables are: k the Boltzmann constant, T the temperature, B the spectral bandwidth, R_0 the 50 Ω generator impedance at the input of the link, I_{bias} the bias current, I_{th} the threshold current, and η_L the laser efficiency.

In Fig. 2, we plotted simulated and measured RIN spectra for several values of bias current. These curves are referenced with the optical power level at the output of the DFB. We notice a good match between simulation and measurement.

B) Optical fiber

The model of the optical fiber takes into account the optical losses and the chromatic dispersion. Input and output signals correspond to the envelope of the optical power as detailed in [3].

C) Photodetector

The photodiode exhibits a 5 GHz bandwidth. Its electrical model contains, in parallel, an optical power-controlled

current source, a dynamic resistance $R_D(1/g_D)$, a capacitance C_D , and a serial resistance R_S for access losses. Three current noise sources have been inserted in order to take into account the shot noise related to the DC photodetected current, the shot noise due to the dark current I_d , and the thermal noise related to the conductance g_D [3]:

$$\langle i_{ns}^2 \rangle = 2 q B (I_{bias} - I_{th}) \frac{\eta_L \eta_{PD}}{L_{OF}},$$

$$\langle i_{nd}^2 \rangle = 2 q B I_d,$$

$$\langle i_{nD}^2 \rangle = 4 k T B g_D,$$

$$(2)$$

where η_L and η_D are, respectively, the laser and the photodiode efficiencies and L_{OF} accounts for the additional optical losses.

III. RESPONSE OF THE MICROWAVE PHOTONIC LINK

We consider a directly modulated microwave photonic link that includes passive impedance matching networks to match the input laser series resistance (r_L) and the output photodiode conductance $(g_D = 1/R_D)$ to the 50 Ω impedance.

As Fig. 1 presents the different elements and noise current sources of the photonic microwave link, a simplified expression for the low-frequency gain G_{LK} of this matched photonic link is given by

$$G_{LK} = \frac{\eta_L^2 \eta_P^2}{4r_L g_D} \frac{1}{L_{OF}^2}.$$
 (3)

In our experiments L_{OF} is set to 14 dB, corresponding to a loss level usually encountered in microwave photonic links implemented in radar systems [8].

Figure 3 shows the measured gain of the link matched around 3 GHz of operating frequency, and for different optical powers corresponding to different laser bias currents. For these measurements, the achieved gain varies from -42 to -38 dB, which is equivalent to a link intrinsic gain between -14 and -10 dB.

Obviously, the observed changes of the gain with the laser bias current cannot be explained with the simple model of equation (3) which is independent of the laser optical power and then small signal. It is thus necessary to take into



Fig. 3. Gain measurements for the directly modulated microwave photonic link including passive impedance matching network (optical losses = 14 dB).

account the modification of the spectral response of the laser as its output power increases. This can be understood in Fig. 4 where the dynamical response of the directly modulated link is calculated for several bias currents.

For the sake of clarity, the passive impedance matching networks are not taken into account in this preliminary simulation. Figure 4 shows that the laser relaxation oscillation frequency becomes very close to the spectral band of interest around 3 GHz when the laser bias current is decreased. This explains the observed gain change (see Fig. 3) when the laser optical power is changed through its bias current.

Consequently, the dynamical behavior of the laser must also impact the noise characteristics of the link as well as its noise figure.

IV. ONP OF THE MATCHED MICROWAVE PHOTONIC LINK

In this section, the impact of the link characteristics on its intrinsic noise is simulated and compared to the measurements in the case of a matched link.

A) Output noise power

The measurement set up of the ONP is presented in Fig. 5.

In order to reduce the thermal noise power contribution of the load resistance, a low noise amplifier (LNA), followed by a second amplifier, are used in the measurement set-up. The first stage LNA has a 1 dB noise figure and the cascade twostage amplifier provides an electrical gain of 64.7 dB. Thus, the output noise power of the photodetector (ONP_{LK}) can be recovered from the noise power measured at the output of the second stage amplifier (ONP_{LK} amplified). Indeed, ONP_{LK} depends on the LNA noise figure NF_{LNA} and amplified link RF gain $G_{(LNA+Amplifier)}$ as

$$ONP_{LK} = \frac{ONP_{LK \ amplified}}{G_{(LNA+Amplifier)}} - (NF_{LNA} - 1)INP_{LK}.$$
 (4)

With the input noise power equals to $INP_{LK} = kTB$.

The measured and simulated ONP at 3 GHz of the intensity modulation-direct detection (IMDD) link described in Fig. 1 versus the DC photodetected current are plotted in



Fig. 4. Simulation of the magnitude of S_{21} for output optical power varying between 11 and 17 dBm.



Fig. 5. Setup for measurement of ONP at a single frequency.

Fig. 6 for four bias currents. One can notice that the simulation results are matching very closely the measurements.

As previously mentioned, three factors contribute to the ONP (equations (1) and (2)). The first one depends on the RIN and the squared photo-detected current, the second term corresponds to shot noise related to the photo-detected current. Finally, the third term, independent of DC photodetected current, corresponds to the shot noise due to the dark current and thermal noise due to the conductance of the photodiode. Consequently, three different coefficients can appear for high optical power (or high laser bias current). In Fig. 6 the slope of the RIN versus the photodetected current is 20 dB/dec whereas the slope of the photodetector's shot noise is 10 dB/dec. Hence, the RIN contribution is usually dominant when the photocurrent is high. However, it can also be predominant for low photocurrents since the RIN level increases when the laser bias current is reduced. Finally, the balance between shot noise and RIN contributions depends on the modulation frequency at which the link is operated. The excess noise due to the RIN contribution becomes significant when the modulation frequency is close to the oscillation relaxation frequency as shown in Fig. 2. Consequently, all these aspects must be taken into account in order to properly estimate the link ONP.

For instance, we plotted in Fig. 6, the evolution of the ONP density as a function of the photodetected current when the input laser optical power is changed. At 3 GHz, the RIN contribution decreases at a given detected photocurrent, when the laser bias current is increased.



Fig. 6. Simulated (solid lines) and measured (dots) ONP of matched link at a frequency of $_3\ {\rm GHz}.$

B) Noise figure

Considering the five noise current sources $\langle i_{nG} \,^2 \rangle$, $\langle i_{nL}^2 \rangle$, $\langle i_{nS}^2 \rangle$, $\langle i_{nD}^2 \rangle$ present at the link output, the simplified expression of the noise figure of the link can be written [3] as

$$NF_{LK} = 1 + \frac{(R_0 + r_L)^2}{R_0} \left[\frac{(I_{bias} - I_{TH})^2 RIN(\omega)}{8kT} + \frac{(I_{bias} - I_{TH})q}{2kT} \frac{L_{OF}}{\eta_L \eta_P} + \left(\frac{qI_d}{2kT} + g_D\right) \left(\frac{L_{OF}}{\eta_L \eta_P}\right)^2 \right].$$
(5)

The DC photodetected current can be expressed as

$$I_{ph} = \frac{\eta_L \eta_P}{L_{OF}} (I_{bias} - I_{TH}).$$
(6)

To compare the noise figure to the electrical losses (inverse of the link gain), we plot them on the same figure. The measured noise figure of the IMDD link (NF_{LK}) can be extracted from to the output noise power:

$$ONP_{LK} = INP_{LK} G_{LK} NF_{LK}$$
(7)

The theoretical and experimental results shown in Fig. 7 are obtained at the RF frequency of 3 GHz for a matched link and when the optical power of the laser is set to 14 dBm. One can notice that, similarly to a pure RF link, when the optical link is RF matched at its ends, the noise figure level remains above the electrical noise value.



Fig. 7. Simulated level of noise figure compared to electrical losses.

V. ONP OF AN UNMATCHED MICROWAVE PHOTONIC LINK

A) Output noise power

In order to investigate the situation where the optical link is unmatched, we now turn the frequency to 2.7 GHz.

At this frequency, the opto-microwave link is mismatched only at the output of the photodetector. We verified experimentally this assumption by measuring at 2.7 GHz the input return loss of both DFB and LNA. They are found to be, respectively, equal to 16 and 14.2 dB. On the other hand, the photodetector output return loss is measured to be 1.4 dB.

Because of this mismatching load condition, photodetector thermal noise due to the conductance g_D gives rise to a "new" reference level of thermal noise at the photonic microwave link output. The low value of g_D , as compared to the inverse of generator resistance, implies that the ONP at 2.7 GHz is lower than the matched thermal noise, *kTB*. This unusual phenomenon is depicted in Fig. 8 for an optical power of 14 dBm.

B) Noise figure

Let focus on this peculiar situation where the link output impedance is mismatched. As with the matched link, we compute and measure the noise figure and the electrical losses versus detected photocurrent. The results are presented in Fig. 9.

One can notice, on the left-hand side, that the noise figure becomes lower than the electrical losses when the DC photocurrent is below 0.3 mA. This result is thereby obtained theoretically and confirmed experimentally. It clearly evidences that an unmatched optical microwave link can present a lower noise figure than an electrical link exhibiting similar losses as already theoretically foreseen by Frigyes [6]. It is important to notice that in the field of pure RF, the noise figure can be lower than the inverse of the electrical transducer gain (or electrical losses) as soon as the passive attenuator presents a mismatched output.

Consequently, the ONP is lower than the input noise power, indeed, the link presents a gain-noise factor product less than 1. That is, if RIN, I_d and g_D are sufficiently low, the main contribution to the noise figure in relation (5) is related to the shot noise, which is proportional to optical



Fig. 8. Simulations (full lines) and measurements (dots) of ONPs in the case of an impedance matched link (at 3 GHz) and unmatched link (at 2.7 GHz).



Fig. 9. Simulated (solid lines) and measured (dots) level of noise figure compared to electrical losses.

losses, unlike the electrical signal which experiences the square of the optical losses. Furthermore, the noise figure can be defined as the ratio of the sum of all square noise current sources at the receiver of the link by the square noise current source at the receiver due to the thermal noise of the resistance of the generator. As g_D is lower than the inverse of R_o , the noise figure can become lower than the electrical loss of the link (Fig. 9). This is the first time, to our knowledge, that this result is experimentally demonstrated and theoretically validated on a real photonic microwave link.

VI. CONCLUSION

In conclusion, we presented a simulation tool based on an accurate physical model embedded in a RF chain simulator. This simulation tool takes into account both optical and electrical characteristics of the link components such as the laser dynamics and impedance matching networks at the input and output of the optical link. It enabled us to get an accurate description of the optical link even under unusual operation conditions. In particular, we were able to expose an unusual behavior in which the noise figure of a microwave optical link can be lower than the electrical losses like in pure microwave links. This highly unusual situation is reached when the optical losses are high enough and when the optical link is impedance mismatched at its output. All the theoretical predictions are tested and confirmed on an IMDD link that we have specifically developed for this purpose. Further refinements include the development of accurate physical models taking into account the second- and third-order nonlinear behavior of the link.

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