pp 1102–1122. © Royal Aeronautical Society 2018 doi: 10.1017/aer.2018.49

# Investigations of synthetic jet control effects on helicopter rotor in forward flight based on the CFD method

## Q.-J. Zhao

zhaoqijun@nuaa.edu.cn

#### X. Chen and Y.-Yang Ma

National Key Laboratory of Science and Technology on Rotorcraft Aeromechanics Nanjing University of Aeronautics and Astronautics Nanjing China

## G.-Q. Zhao

General and Aerodynamic Institute the First Aircraft Institute of AVIC Xi'an China

# ABSTRACT

To investigate the control effect of the synthetic jet on the aerodynamic characteristic of rotors, a numerical simulation procedure for the rotor flowfield is established. First, a movingembedded grid method and an unsteady Reynolds Averaged Navier–Stokes (URANS) solver are established for predicting the complex flowfield of rotors. A velocity jet boundary condition over the jet actuator orifice is constructed, and a numerical method for simulating the active flow control on rotors is developed. Then, the effectiveness of the simulation method is validated by comparing the numerical results of jet control on NACA 0015 aerofoil with the experimental data. At last, the aerodynamic characteristic of rotors with synthetic jet actuators located on the suction surface of the blade in forward flight is calculated. The results indicate that the synthetic jet has the capability of improving the aerodynamic characteristic of rotors, especially in inhibiting the flow separation over the surface. In addition, the increase of the jet momentum coefficient and the jet angle can both enhance the lift coefficient in the retreating side. Compared with a single jet, jet arrays have better control effects on improving the aerodynamic characteristic of rotors in forward flight.

**Keywords:** Rotor; synthetic jet; aerodynamic characteristic; forward flight; helicopter; CFD method

Received 25 August 2017; revised 2 April 2018; accepted 4 April 2018.

# NOMENCLATURE

с	chord
$C_{u0}$	steady component of jet momentum coefficient
$C_{um}$	unsteady component of jet momentum coefficient
d	length of jet orifice along spanwise direction
$F^+$	jet frequency
$\mathbf{F}_{v}$	viscous fluxes
$\mathbf{F}_{c}$	convective fluxes
h	length of jet orifice along chordwise direction
$M_{\rm tip}$	mach number at the blade tip
Ma	mach number
<b>n</b> <sub>jet</sub>	direction vector of synthetic jet
Ř	radius of rotor
S	boundary of grid cell
t	real time
$t^*$	non-dimensional time
$U_{\rm tip}$	velocity of blade tip
$U_0$	amplitude of the steady components of velocity
$U_m$	amplitude of the unsteady components of velocity
W	conservation variables
τ	pseudo-time
Ω	volume of grid cell
ξ	chordwise direction
$\omega_{jet}$	excitation frequency of jet
$\omega_{rotor}$	angular velocity of rotor
η	cross-stream direction
ζ	spanwise direction
μ	advance ratio
$\Delta C_l$	lift coefficient increment
θ	synthetic jet angle
ψ	azimuthal angle

# **1.0 INTRODUCTION**

The aerodynamic characteristic of the rotor has an important impact on the flight performance of the helicopter. To avoid the large rolling moment and maintain the high lift of the blade, the retreating blade is always operated under a large Angle-of-Attack (AoA), while the AoA of the advancing side is much lower<sup>(1)</sup>. This periodic variation of AoA makes the rotor undergo dynamic stall in the retreating side, and this phenomenon may induce the flow separation on the surface of the rotor blade in forward flight<sup>(2)</sup>, resulting in the decrease of the lift and the increase of the vibration level. Therefore, the investigation on preventing the flow separation and improving the aerodynamic characteristic of the rotor has become one of the key and difficult points in the field of helicopter technology. In recent years, it has been observed that the active flow control was one of the most promising methods for improving the aerodynamic characteristic of the operating envelope of the rotorcraft<sup>(3-5)</sup>, and it is indicated to be one of the main design strategies of the advanced rotor in the future<sup>(6, 7)</sup>.

As a novel active flow-control method, the synthetic jet is an efficient technology to control the flow separation and the stall of the aerofoil<sup>(8-11)</sup>. In order to explore the control mechanism of the synthetic jet on the flow separation, Seifert<sup>(12)</sup> conducted active control experiments on NACA0015 aerofoil, and the capability of the synthetic jet on delaying the stall of the aerofoil was verified. Several experimental results indicated that the synthetic jet localised near the location where the flow separation formed can significantly enhance stall characteristics of the aerofoil, including the increment of the maximum lift and the stall angle<sup>(13-15)</sup>.

In the aspect of numerical simulations of the synthetic jet control, Kral<sup>(16)</sup> introduced a tangential synthetic jet at the leading edge of NACA0015 aerofoil, and used a Reynolds-averaged Navier-Stokes (RANS) equations to investigate the effect of the synthetic jet on the flow over the aerofoil, and it was demonstrated that the RANS approach could be used to effectively simulate the active flow control by employing unsteady jets. Lorber<sup>(17)</sup> conducted the investigations on delaying the retreating blade stall by the directed synthetic jet (DSJ) for aerofoil SC2110. Hassan<sup>(18)</sup> numerically simulated the effects of surface bowing and suction on the aerodynamic characteristics of the five-bladed MD-900 rotor in a low-speed descent flight by solving the unsteady 3D full-potential equations. Based on Euler equations, Dindar<sup>(19)</sup> numerically investigated the potential merits of using the transpiration flow control on the rotor in hovering flight. However, their work is difficult to simulate the rotor vortex flowfield and the detail of the jet flow due to the restriction of their methods.

In addition, there are still some questions which have not been overcome in the applications of synthetic jets<sup>(20)</sup>. The application of the synthetic jet on rotors is still theoretic because of numerous controlling parameters, such as the jet location, the jet momentum coefficient and the jet angle. Furthermore, the parametric analysis of the dual-jet control effect on the aerodynamic characteristic of rotors has not been carried out well.

In this paper, numerical simulations about the active flow control of the rotor by using the synthetic jet are conducted systematically. First, the solution of the rotor flowfield is obtained by solving the unsteady RANS equations with  $k-\omega$  SST turbulence model, and a moving-embedded grid method is employed. Surface points of the blade with a synthetic jet orifice are constructed, and grids around the section of the blade are quickly generated by solving Poisson equations. Additionally, a dimensionless velocity boundary condition is established to simulate flow control effects of the synthetic jet on the blade surface. By comparing the calculated results and the test data about the synthetic jet on the aerofoil, the feasibility and efficiency of the present numerical method are evaluated. Furthermore, parametric analyses of the synthetic jet are carried out to investigate the control mechanism of the synthetic jet on the aerodynamic characteristic of the rotor in forward flight, and some conclusions are obtained.

## 2.0 METHODOLOGIES

## 2.1 Flowfield solution method

The CFD methods<sup>(21)</sup> based on unsteady RANS equations are employed to simulate the flowfield of the rotor.

$$\frac{\partial}{\partial \tau} \int \mathbf{W} d\Omega + \frac{\partial}{\partial t} \int \mathbf{W} d\Omega + \oint (\mathbf{F}_c - \mathbf{F}_v) dS = 0 \qquad \dots (1)$$

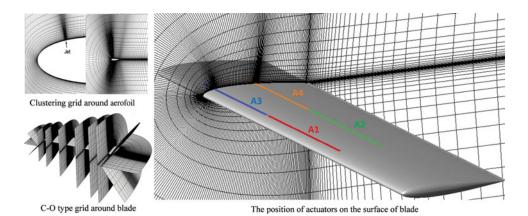


Figure 1. (Colour online) Body-fitted grids around rotor aerofoil and blade.

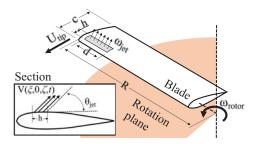


Figure 2. (Colour online) Jet boundary condition on the surface of the helicopter rotor.

where **W** are conservation variables,  $\mathbf{F}_{v}$  are the viscous fluxes, and  $\mathbf{F}_{c}$  are convective fluxes including the motion of the dynamic grid system.  $\tau$  and *t* are the pseudo time and the real time respectively.  $\Omega$  is the volume of the grid cell, and *S* is the boundary of the grid cell.

The dual-time method is used to simulate the unsteady flowfield of the rotor, and the subiteration is fulfilled by the implicit LU-SGS method<sup>(22)</sup> which can be implemented easily on vector and parallel computers. To accurately predict the non-linear and unsteady vortical flowfield of the rotor, the third-order Roe scheme and the MUSCL approach<sup>(23)</sup> are employed for the discretisation of convective fluxes. Trying to simulate the flow separation over the surface of the rotor blade, the k- $\omega$  SST turbulence model<sup>(24)</sup> is employed to calculate the turbulence viscosity.

#### 2.2 Grid generation method

The quality of grids around the rotor has a direct influence on the flowfield solution precision. There are three steps in the grid generation for active flow-control investigations on the rotor. First, the body-fitted and orthogonal C-type grids around rotor aerofoils are generated by solving Poisson equations. To capture the flow characteristics of the synthetic jet in detail, clustering grids over the orifice are generated. Second, C-O type grids around the rotor blade are generated automatically. At last, the moving-embedded grids are generated including two key points: (1) identification for the boundary cell of the hole; (2) searching of donor elements.

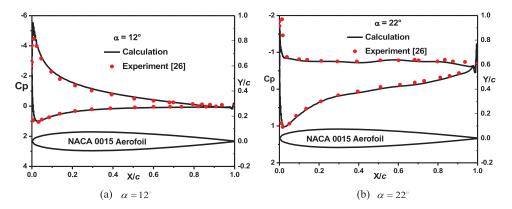


Figure 3. (Colour online) Pressure distributions of NACA0015 aerofoil at different angles of attack. (a)  $\alpha = 12^{\circ}$  (b)  $\alpha = 22^{\circ}$ .

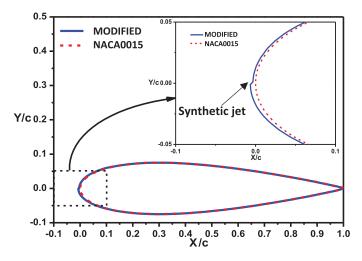


Figure 4. (Colour online) Modified NACA0015 aerofoil with the synthetic jet.

The 'top-map' method and the 'PSSDE' method<sup>(25)</sup> are adopted in the procedure of embedded grids. Body-fitted grids around rotors are shown in Fig. 1.

#### 2.3 Boundary condition of synthetic jet

The flow perturbation from the synthetic jet actuator on the rotor is modeled by a suction/blowing type boundary condition<sup>(26)</sup>.

The non-dimensional velocity of the synthetic jet at the actuator surface is introduced by:

$$\mathbf{V}(\xi, \eta = 0, \zeta, t) = \sqrt{\frac{Rc}{2hd}} U_{\text{tip}} \left[ \sqrt{C_{u0}} + \sqrt{2C_{um}} \sin\left(\omega_{\text{rotor}} F^+ t \frac{c}{U_{\text{tip}}}\right) \right] \mathbf{n}_{\text{jet}}, \qquad \dots (2)$$

where  $\xi$  denotes the chordwise direction,  $\eta$  denotes the cross-stream direction,  $\zeta$  denotes the spanwise direction,  $F^+$  denotes the jet frequency,  $C_{u0}$  is the steady component of jet momentum coefficient,  $C_{um}$  is the unsteady component of jet momentum coefficient,  $\omega_{\text{rotor}}$ 

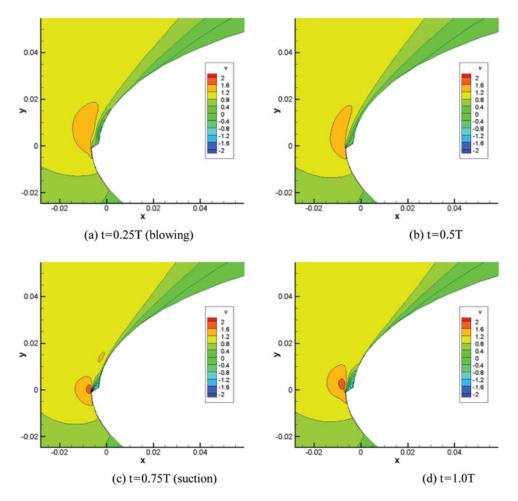


Figure 5. (Colour online) Vertical velocity profile in one synthetic jet period. (a) t = 0.25T (blowing) (b) t = 0.5T (c) t = 0.75T (suction) (d) t = 1.0 T.

is the angular velocity of rotor,  $\mathbf{n}_{jet}$  is the direction vector of synthetic jet, *h* is the length of the jet orifice along the chordwise direction, *d* is the length of the jet orifice along the spanwise direction, *R* is the radius of the rotor, and  $U_{tip}$  is the velocity of the blade tip, as shown in Fig. 2.

The non-dimensional jet excitation frequency is determined by:

$$F^{+} = \frac{\frac{\omega_{jet}}{2\pi}}{\frac{\omega_{rotor}}{2\pi}} = \frac{\omega_{jet}}{\omega_{rotor}} \qquad \dots (3)$$

The instantaneous velocity of the synthetic jet and the non-dimensional jet momentum coefficient are determined by:

$$U(t) = U_0 + U_m \sin(\omega_{\text{jet}} t), \ C_{u0} = 2 \frac{hd}{Rc} \left(\frac{U_0}{U_{\text{tip}}}\right)^2, \ C_{um} = 2 \frac{hd}{Rc} \left(\frac{U_m}{\sqrt{2}U_{\text{tip}}}\right)^2, \qquad \dots (4)$$

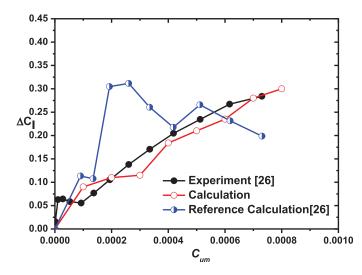


Figure 6. (Colour online) Effects of Cum on the lift coefficient increment of the aerofoil.

Table 1The Position of Actuators on the Surface of the Blade

Actuator	A1	A2	A3	A4
Spanwise position	50%-70%R	50%-70%R	70%-90%R	70%-90%R
Chordwise position	10%c	50%c	10%c	50%c

where  $U_0$  and  $U_m$  are the amplitude value of steady and unsteady components of the velocity, respectively.

The non-dimensional time and the definition of the angular velocity of the rotor are:

$$t^* = \frac{t}{c/U_{\text{tip}}}, \ \omega_{\text{rotor}} = \frac{U_{\text{tip}}}{R} \qquad \dots (5)$$

The boundary condition of the synthetic jet on the rotor can be converted to:

$$\mathbf{V}(\xi, 0, \zeta, t) = \sqrt{\frac{Rc}{2hd}} U_{\text{tip}} \left[ \sqrt{C_{u0}} + \sqrt{2C_{um}} \sin\left(\frac{c}{R}F^+t\right) \right] \mathbf{n}_{\text{jet}} \qquad \dots (6)$$

# **3.0 RESULTS AND ANALYSES**

## 3.1 Jet control of rotor aerofoil

The baseline case for NACA0015 under a condition of Ma = 0.15 and  $Re = 1.85 \times 10^6$  is calculated firstly. The grid has a resolution of 447 × 90. There are 185 points on the lower and upper surfaces of the aerofoil, respectively; 39 points on each of the wake cuts which extend from the aerofoil's trailing edge to the outflow boundary of the farfield; 90 points in the direction normal to the local surface of the aerofoil; and 36 points over the jet orifice. Fig. 3 shows the pressure distribution on the surface of NACA0015 aerofoil at different angles

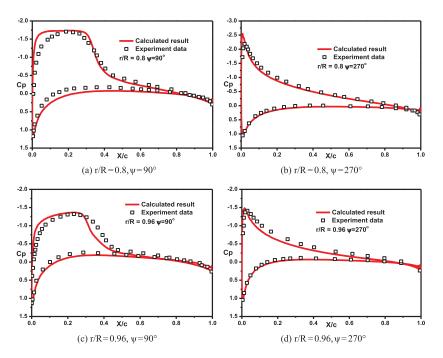


Figure 7. (Colour online) Pressure distributions on different blade sections. (a) r/R = 0.8,  $\psi = 90^{\circ}$ ; (b) r/R = 0.8,  $\psi = 270^{\circ}$ ; (c) r/R = 0.96,  $\psi = 90^{\circ}$ ; (d) r/R = 0.96,  $\psi = 270^{\circ}$ .

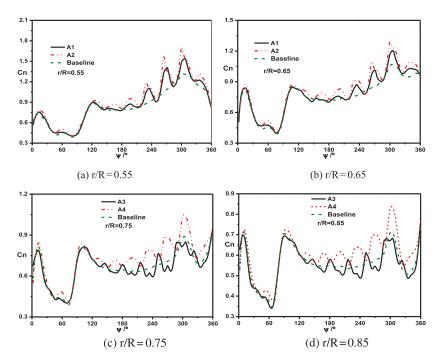


Figure 8. (Colour online) Comparisons of the control effects of different actuators on the sectional normal force. (a) r/R = 0.55, (b) r/R = 0.65, (c) r/R = 0.75, (d) r/R = 0.85.

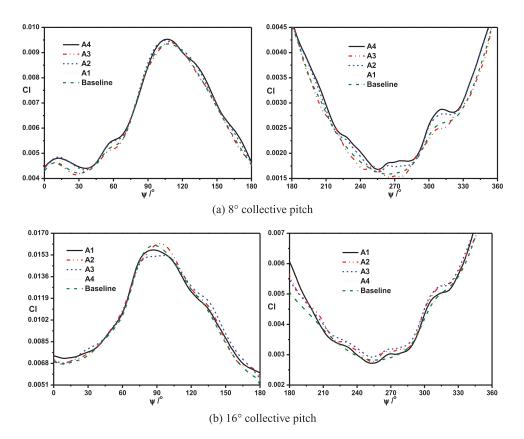


Figure 9. (Colour online) Comparisons of the control effects of different actuators on the lift coefficient of the blade. (a) 8° collective pitch, (b) 16° collective pitch.

of attack. As can be seen, numerical results are in good agreements with the experimental  $data^{(26)}$ .

Figure 4 shows the modified aerofoil of NACA0015 with a synthetic jet. The synthetic jet actuator is implemented at the leading edge of the aerofoil with the width of 0.14% chord length. In the experiment, the direction of the jet is normal to the surface of the aerofoil, and AoA of the aerofoil is 22°, and Re is  $1.2 \times 10^6$ . Fig. 5 shows the vertical velocity profile in one synthetic jet period under a condition of  $C_{u0} = 0$  and  $C_{um} = 0.001$  ( $U_0/U_{\infty} = 0.6$  and  $U_m/U_{\infty} = 0.6$ ). Under the influence of the periodic movement of the synthetic jet, the velocity near the leading edge changes obviously with the time. The mixing between the internal and external layers of the flow in the boundary layer was heightened, and the energy of the boundary layer was strengthened at the same time. These two effects are useful for the inhibition of the flow separation and the increment of the lift force.

Figure 6 shows the effect of  $C_{um}$  on the lift coefficient increment ( $\Delta C_l$ ). As can be seen, the lift coefficient increment increases with the increase of  $C_{um}$ , and it is similar to the variation of experimental data. In addition, the numerical results are better than those from reference<sup>(26)</sup>.

## 3.2 Verification of aerodynamic characteristic of rotor in forward flight

The Caradonna-Tung (CT)  $rotor^{(27)}$  has two rectangular blades with a conventional NACA 0012 aerofoil, and the aspect ratio is 6. The CT rotor is used to investigate the effects of the

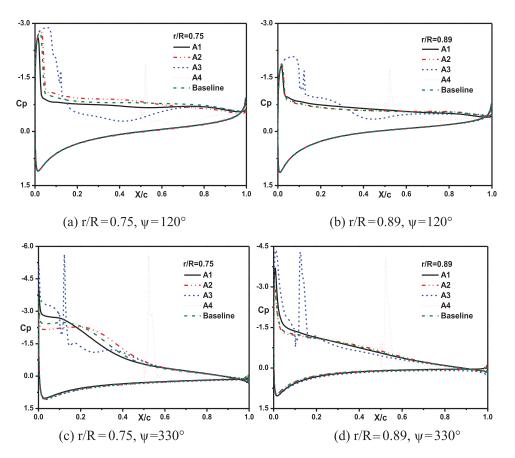


Figure 10. (Colour online) Comparisons of the sectional pressure coefficient with different actuators at a collective pitch of 16°. (a) r/R = 0.75,  $\psi = 120^{\circ}$ ; (b) r/R = 0.89,  $\psi = 120^{\circ}$ ; (c) r/R = 0.75,  $\psi = 330^{\circ}$ ; (d) r/R = 0.89,  $\psi = 330^{\circ}$ .

synthetic jet on the rotor blade in forward flight. Four jet actuators are fitted on the upper surface at different positions as shown in Fig. 1. The actuators (A1, A3) are placed at 10%c on the suction surface while the other ones (A2, A4) are placed at 50%c, and the width of the jet orifice is 1%c as indicated in Table 1.

The grid of the blade has a resolution of  $343 \times 40 \times 131$ , and there are 137 points on the lower and upper surfaces of the blade section, 11 points over the jet orifice. The baseline case for this rotor under a condition of  $M_{tip} = 0.628$  and  $\mu = 0.3$  is calculated first. Figure 7 shows the pressure distribution on the blade section at a collective pitch of 8°. As can be seen, the numerical results present a good agreement with the data obtained from Ref. (28).

#### 3.3 Effects of synthetic jet location

The control effects of the synthetic jet on the aerodynamic characteristic of rotors with different synthetic jet locations are investigated ( $M_{tip} = 0.628$ ,  $\mu = 0.3$ ), and the excitation frequency is  $F^+ = 10$ .

Figure 8 shows the control effects of the synthetic jet on the normal force coefficient at a collective pitch of 8°. The steady component of the momentum coefficient is  $C_{u0} = 0$ , and the

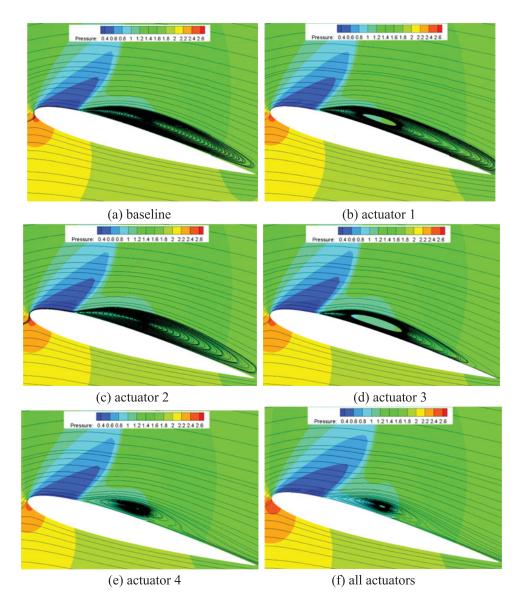


Figure 11. (Colour online) Streamlines over the r = 0.8 R blade section at 90° azimuthal angle. (a) baseline, (b) actuator 1, (c) actuator 2, (d) actuator 3, (e) actuator 4, (f) all actuators.

unsteady component is  $C_{um} = 0.004 (U_m/U_{tip}=1)$ , and the jet angle is  $\theta_{jet} = 25^\circ$ . As seen in the figure, the synthetic jet has a significant effect on the sectional normal force coefficient of blades. The synthetic jets induced by actuators A2 and A4 are more efficient than those by actuators A1 and A3.

The control effects of the synthetic jet on the lift coefficient of the blade at different collective pitches are illustrated in Fig. 9. The synthetic jet induced by actuator A4 is more efficient on enhancing lift than other actuators at the collective pitch of  $8^{\circ}$ , since there is no flow separation on the blade surface. When the collective pitch is  $16^{\circ}$ , the control effect of actuator A3 is the best, followed by A2. As a result, the synthetic jet located in the middle

ZHAO ET AL

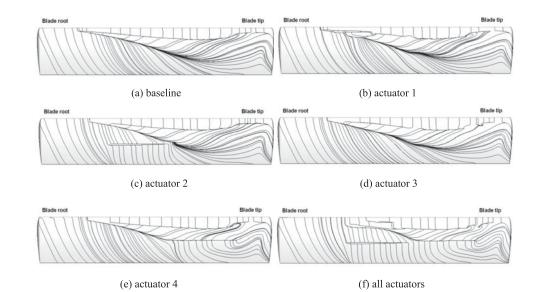


Figure 12. Streamlines on the blade surface under the control of different actuators at 90° azimuthal angle. (a) baseline, (b) actuator 1, (c) actuator 2, (d) actuator 3, (e) actuator 4, (f) all actuators.

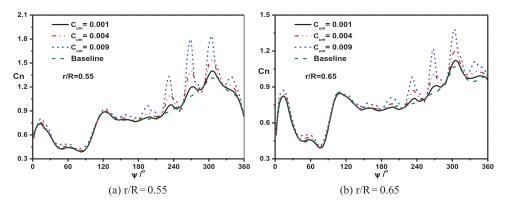


Figure 13. (Colour online) Comparisons of the sectional normal force coefficient with different jet momentum coefficients of A1 actuator at a collective pitch of 8°. (a) r/R = 0.55, (b) r/R = 0.65.

along the chordwise direction has a better effect at a smaller collective pitch, and the synthetic jet near the leading edge has a better effect at a bigger collective pitch.

Figure 10 shows the sectional pressure coefficient of the rotor blade with different actuators at a collective pitch of 16°. The synthetic jets (A1 and A3) near the leading edge of the blade have a more obvious effect on the disturbance of the pressure gradient. They can effectively weaken the strength of the shock wave. The synthetic jets (A2 and A4) located in the middle along the chordwise direction affect the pressure coefficient in a wider scope. They may decrease the adverse pressure gradient, and the flow separation can be delayed.

The streamlines over the r = 0.8 R blade section at the 90° azimuthal angle are illustrated in Fig. 11. It can be seen that there is a large flow separation at the upper surface due to the increase of the adverse pressure gradient caused by the shock wave. Under the control of

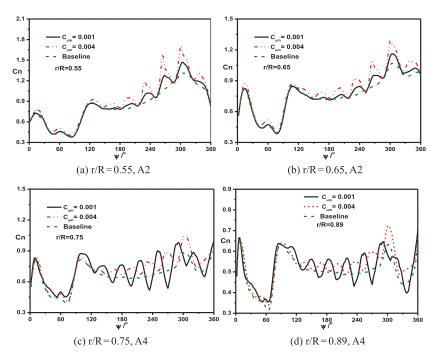


Figure 14. (Colour online) Comparisons of the sectional normal force coefficient with different jet momentum coefficients at a collective pitch of 8°. (a) r/R=0.55, A2; (b) r/R=0.65, A2; (c) r/R=0.75, A4; (d) r/R=0.89, A4.

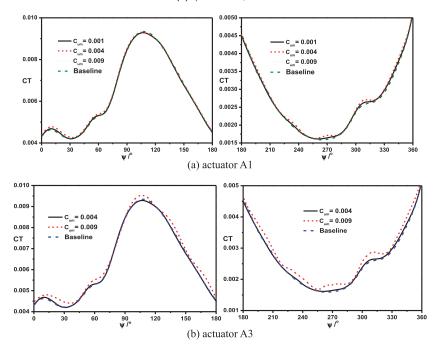


Figure 15. (Colour online) Comparisons of the lift coefficient with different jet momentum coefficients at a collective pitch of 8°. (a) actuator A1, (b) actuator A3.

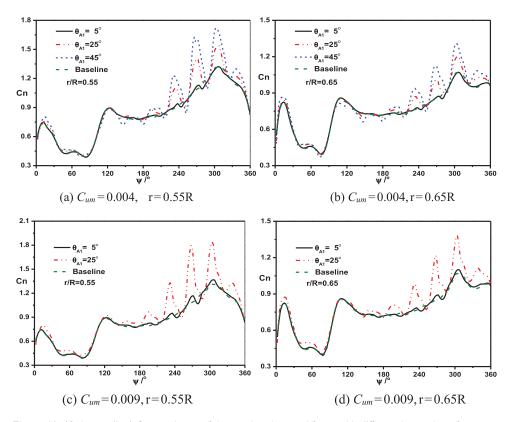


Figure 16. (Colour online) Comparisons of the sectional normal force with different jet angles of actuator A1 at collective pitch of 8°. (a)  $C_{um} = 0.004$ , r = 0.55R; (b)  $C_{um} = 0.004$ , r = 0.65R; (c)  $C_{um} = 0.009$ , r = 0.55R; (d)  $C_{um} = 0.009$ , r = 0.65R.

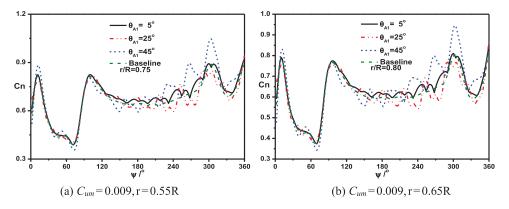


Figure 17. (Colour online) Comparisons of the sectional normal force of the blade with different jet angles of actuator A1 at a collective pitch of  $16^{\circ}$ . (a)  $C_{um} = 0.009$ , r = 0.55R; (b)  $C_{um} = 0.009$ , r = 0.65R.

actuator A1 (or A2), there is still a flow separation on the upper surface, since 0.8 R section is outside of its control area. Under the control of actuator A3, the area of flow separation region is reduced. Under the control of actuator A4, the synthetic jet has a better effect on preventing the flow separation. It is because that the actuator A4 is located in the flow separation region. In

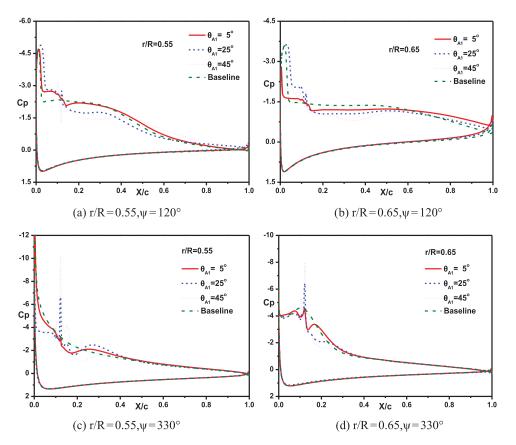


Figure 18. (Colour online) Comparisons of the blade sectional pressure coefficient with different jet angles of actuator A1 at a collective pitch of 16°. (a) r/R = 0.55,  $\psi = 120^{\circ}$ ; (b) r/R = 0.65,  $\psi = 120^{\circ}$ ; (c) r/R = 0.55,  $\psi = 330^{\circ}$ ; (d) r/R = 0.65,  $\psi = 330^{\circ}$ .

addition, the best control effect on preventing flow separation could be obtained if all actuators work.

The streamlines on the blade surface at  $90^{\circ}$  azimuthal angle are shown in Fig. 12. It can be seen that the flow separation is obvious near the  $90^{\circ}$  azimuthal angle. Under the control of the synthetic jet, the stability of the boundary layer increases, and the local adverse pressure gradient decrease, and then the flow separation is inhibited.

## 3.4 Effect of momentum coefficient

With the different momentum coefficient of actuators, the control effects of the synthetic jet on the aerodynamic characteristic of the rotor are investigated. Three typical momentum coefficients,  $C_{um} = 0.001 \ (U_m/U_{tip} = 0.5), C_{um} = 0.004 \ (U_m/U_{tip} = 1)$  and  $C_{um} = 0.009 \ (U_m/U_{tip} = 1.5)$ , are employed. The jet angle is constantly equal to 25°.

Figure 13 shows the sectional normal force coefficient under the control of actuator A1 at a collective pitch of 8°. The synthetic jet has an obvious effect on the normal force in the retreating side, and the amplitude value and the average value of the sectional normal force increase with the increase of momentum coefficient.

Table 2The Combinations of Jet Arrays

Array number	array 1	array 2	array 3	array 4
Jet angle of A1 and A3	20°	20°	40°	40°
Jet angle of A2 and A4	20°	40°	20°	40°

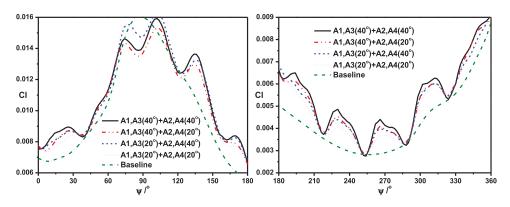


Figure 19. (Colour online) Comparisons of the lift coefficient of the blade with different jet array combinations.

Figure 14 shows the sectional normal force coefficient under the control of actuators A2 and A4 at a collective pitch of 8°. Under the control of actuators A2 and A4, the average value of the sectional normal force significantly increases with the increase of the momentum coefficient. However, under the control of actuator A4, the oscillation of the sectional normal force decreases with the increase of the momentum coefficient.

Figure 15 shows the lift coefficient under the control of actuators A1 and A3 at a collective pitch of 8°. The variation of the momentum coefficient has a noticeable effect on the lift coefficient in the retreating side, and the lift coefficient obviously increases with the increase of momentum coefficient.

#### 3.5 Effects of synthetic jet angle

To investigate the influence of the synthetic jet angle on the control effect, the jet angles are selected as  $5^{\circ}$ ,  $25^{\circ}$  and  $45^{\circ}$ . Figure 16 shows the sectional normal force of the blade with different jet angles under the control of actuator A1 at a collective pitch of  $8^{\circ}$ . There is a significant variation in the sectional normal force of the blade under the control of the synthetic jet, especially in the retreating side. In the retreating side, the average value and the amplitude value of the normal force obviously increase with the increase of the jet angle. As the jet angle changes from  $5^{\circ}$  to  $45^{\circ}$ , there is more energy of the synthetic jet into the boundary-layer flow, and the mixing phenomenon is more obvious.

Figure 17 shows the sectional normal force of the blade with different jet angles under the control of actuator A1 at a collective pitch of 16°. When the jet angle is 45°, there is the most obvious effect on the normal force in the retreating side. Since the collective pitch is higher than that in Fig. 16, the phenomenon of the spanwise flow near the blade tip is more obvious,

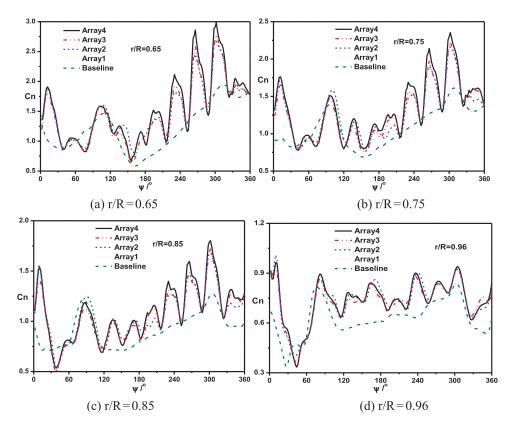


Figure 20. (Colour online) Comparisons of the blade sectional normal force with different jet arrays. (a) r/R = 0.65, (b) r/R = 0.75, (c) r/R = 0.85, (d) r/R = 0.96.

and the flow separation is possible in the retreating side. Therefore, a larger jet angle has a better control effect on the aerodynamic force of the blade at a large collective pitch.

Figure 18 shows the comparisons of the sectional pressure coefficient with different jet angles under the control of actuator A1 at a collective pitch of 16°. When the jet angle increases in a certain extent, the control effect on the sectional pressure coefficient is more obvious.

## 3.6 Effects of synthetic jet array

Since the jet angle has an important effect on the aerodynamic characteristic of rotors, four jet arrays with different jet angles are investigated. The collective pitch is  $16^{\circ}$  and the jet momentum coefficient is  $C_{u0} = 0$ ,  $C_{um} = 0.009 (U_m/U_{tip} = 1.5)$ . The different combinations of jet arrays are shown in Table 2.

Figure 19 shows the lift coefficient of the blade with different jet arrays. Compared with the control effect under a single actuator, the control effect of a jet array on the lift coefficient is more obvious. All arrays can effectively enhance the lift coefficient of the blade, especially in the retreating side. In addition, array 4 has the best control effect on enhancing the lift coefficient, followed by array 2 and array 3.

Figure 20 shows the blade sectional normal force with different jet arrays. The normal forces at different sections are all enhanced by jet arrays. When the section is closer to the

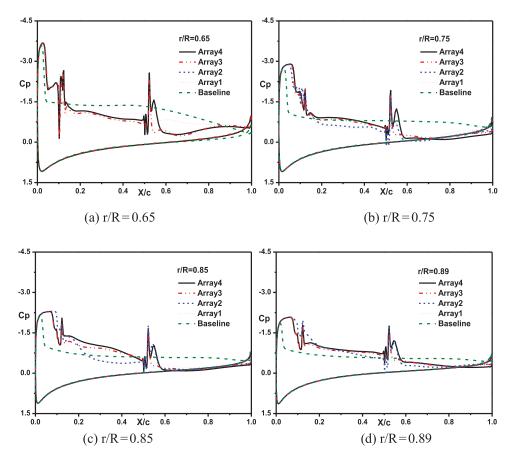


Figure 21. (Colour online) Comparisons of the pressure coefficient of the blade section with different jet arrays at  $120^{\circ}$  azimuthal angle. (a) r/R = 0.65, (b) r/R = 0.75, (c) r/R = 0.85, (d) r/R = 0.89.

blade root, the ratio of the jet velocity to the sectional relative velocity is larger, and then the oscillation amplitude value of the normal force is larger. At the same time, array 4 has the best control effect on enhancing the normal force coefficient of the blade, which is coincident with that in Fig. 19.

Figure 21 shows the pressure coefficient of the blade with different jet arrays at 120° azimuthal angle. As seen in the figure, there is an obvious variation of the pressure coefficient on the upper surface under the control of jet arrays. Since the adverse pressure gradient on the upper surface decreases, the flow separation may be inhibited, and then the blade sectional normal force increases. The jet caused by actuator A1 or A3 can effectively control the adverse pressure gradient near the leading edge, and the jet caused by actuator A2 or A4 can enlarge the pressure coefficient near the trailing edge.

Figure 22 shows comparisons of the pressure coefficient of the blade with different jet arrays at 330° azimuthal angle. Compared with the advancing side, the jet caused by A2 or A4 has a better control effect on the pressure coefficient of the blade surface, and the pressure coefficient near the trailing edge is almost the same as the baseline.

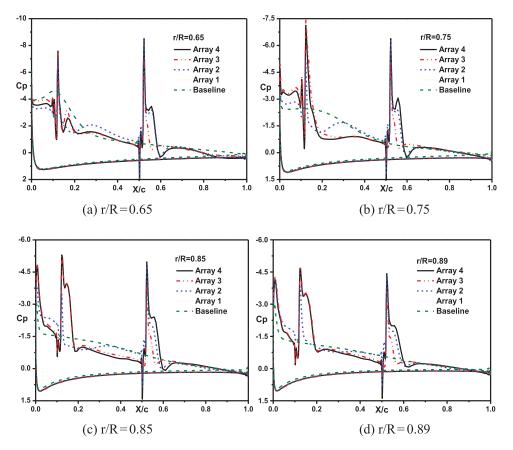


Figure 22. (Colour online) Comparisons of the pressure coefficient of the blade section with different jet arrays at  $330^{\circ}$  azimuthal angle. (a) r/R = 0.65, (b) r/R = 0.75, (c) r/R = 0.85, (d) r/R = 0.89.

# 4.0 CONCLUSIONS

Based on a numerical method for predicting the unsteady flowfield of the rotor, the control effects of the synthetic jet on the aerodynamic characteristic of the rotor in forward flight are effectively simulated and then parametric analyses are conducted. Some valuable conclusions are obtained as follows:

- (1) The synthetic jet has the capability of improving the aerodynamic characteristic of the rotor, especially in inhibiting flow separation over the blade surface. The synthetic jet has a better control effect on the lift coefficient in the retreating side. The synthetic jet located in the middle of the aerofoil along the chordwise direction has a better effect on improving aerodynamic characteristics at a small collective pitch and the synthetic jet near the leading edge has a better effect at a large collective pitch.
- (2) When the synthetic jet is in the middle of the blade, the oscillation amplitude value and the average value of the normal force both increase with the increase of momentum coefficient. When the synthetic jet is near the blade tip, the oscillation amplitude value decreases with the increase of momentum coefficient.

- (3) In the advancing side, there is no obvious difference in the normal force coefficient among different jet angles. In the retreating side, the average value and oscillation amplitude value of the normal force coefficient increase with the increase of jet angle.
- (4) Comparing with the single synthetic jet actuator, jet arrays have a better control effect on the aerodynamic characteristic of the rotor. In the retreating side, the lift coefficient of the blade increases with the increase of jet angles (15°-45°). The pressure coefficient of the suction surface could be totally affected by the jet array. The synthetic jet near the blade leading edge has an effective control on the adverse pressure gradient, resulting in the weakening of the shock wave. The synthetic jet located in the middle along the chordwise direction can increase the pressure coefficient near the blade trailing edge.

# REFERENCES

- 1. LEISHMAN, J.G. *Principles of Helicopter Aerodynamics*, Cambridge University Press, 2000, New York, US, Chapter 8.
- 2. YU, Y.H., LEE, S., MCALISTER, K.W., TUNG, C. and WANG, C.M. Dynamic stall control for advanced rotorcraft application, *AIAA J*, 1995, **33**, (2), pp 289-295.
- 3. LIIPFERT, E., POTTLER, K. and ULMER, S. Parabolic trough optical performance analysis techniques, *J Sol Energ-T ASME*, 2007, **147**, (2), pp 147-152.
- 4. MA, X.Y., GUO, H.T. and FAN, Z.L. Investigating of simulation methods for synthetic jet, *Procedia Engineering*, 2012, **31**, pp 416-421.
- 5. XIA, Q.F. and ZHONG, S. Enhancement of laminar flow mixing using a pair of staggered lateral synthetic jets, *Sensors and Actuators A: Physical*, 2014, **207**, pp 75-83.
- 6. NAGIB, H., GREENBLATT, D. and KIEDAISCH, J. Effective flow control for rotorcraft applications at flight Mach number, 31st AIAA Fluid Dynamics Conference and Exhibit, 11–14 June 2001, Anaheim, California, US.
- 7. Woo, G.T., CRITTENDEN, T. and GLEZER, A. Transitory control of dynamic stall over a stalled airfoil, 39th AIAA Fluid Dynamics Conference, 22–25 June 2009, San Antonio, Texas, US.
- 8. SMITH, B.L. and GLEZER, A. The formation and evolution of synthetic jets, *Physics of Fluids*, 1998, **10**, (9), pp 2281-2297.
- DURRANI, D. and HAIDER, B.A. Study of stall delay over a generic airfoil using synthetic jet actuator, 49th AIAA Aerospace Sciences Meeting Including the New Horizons Forum and Aerospace Exposition, 4–7 January 2011, Orlando, Florida, US. AIAA No 2011-943, 2011. doi:10.2514/6.2011-943.
- SANDRA, U. Experimental analysis and analytical modeling of synthetic jet cross flow interactions, Dissertation, University of Maryland, 2007.
- 11. HASSAN, A.A. Oscillatory and pulsed jets for improved airfoil aerodynamics- a numerical simulation, 42nd AIAA Aerospace Sciences Meeting and Exhibit, 5–8 January 2004, Reno, Nevada, US. doi:10.2514/6.2004-227.
- 12. SEIFERT, A., DARABI, A. and WYGNANSKI, I. Delay of airfoil stall by periodic excitation, *AIAA J*, 1999, **33**, (4), pp 691-707.
- GILARRANZ, J.L., TRAUB, L.W. and REDINIOTIS, O.K. Characterization of a compact, high-power synthetic jet actuator for flow separation control, 40th AIAA Aerospace Sciences Meeting & Exhibit, 14–17 January 2002, Reno, Nevada, US. AIAA No. 2002-127, 2002. doi:10.2514/6.2002-127.
- 14. ZHANG, P.F., YAN, B. and DAI, C.F. Lift enhancement method by synthetic jet circulation control, *Science China Technological Sciences*, 2012, **55**, (9), pp 2585-2592.
- MONIR, E.H., TADJFAR, M. and BAKHTIAN, A. Tangential synthetic jets for separation control, J Fluids and Structures, 2014, 45, pp 50-65.
- KRAL, L.D., DONOVAN, J.F. and CAIN, A.B. Numerical simulation of synthetic jet actuators, 4th Shear Flow Control Conference, 29 June–2 July 1997, Snowmass village, Colorado, US. AIAA No. 1997-1824, 1997. doi:10.2514/6.1997-1824.

- LORBER, P., MCCORMICK, D. and ANDERSON, B.W. Rotorcraft retreating blade stall control, Fluids 2000 Conference and Exhibit, 19–22 June 2000, Denver, Colorado, US. AIAA No. 2000-2475, 2000. doi:10.2514/6.2000-2475.
- HASSAN, A.A., STRAUB, F.K. and CHARLES, B.D. Effects of surface blowing/suction on the aerodynamics of helicopter rotor blade-vortex interactions (BVI)- a numerical simulation, Proceedings of 52nd Annual Forum of AHS, 4–6 June 1996, Washington, DC, US.
- DINDAR, M., JANSEN, K. and HASSAN, A.A. Effect of transpiration flow control on hovering rotor blades, 17th Applied Aerodynamics Conference, 28 June–1 July 1999, Norfolk, Virginia, US. AIAA No. 1999-3192, 1999. doi:10.2514/6.1999-3192.
- CAIN, A.B., KRAL, L.D., DONOVAN, J.F. and SMITH, T.D. Numerical simulation of compressible synthetic jet flows, 36th AIAA Aerospace Sciences Meeting and Exhibit, 12–15 January 1998, Reno, Nevada, US. doi:10.2514/6.1998-84.
- ZHAO, Q.J., ZHAO, G.Q., WANG, B., WANG, Q., SHI, Y.J. and XU, G.H. Robust Navier-Stokes method for predicting unsteady flowfield and aerodynamic characteristics of helicopter rotor, *Chinese J Aeronautics*, 2018, 31, (2), pp 214-224.
- SHAROV, D. and NAKAHASHI, K. Low speed preconditioning and LU-SGS scheme for 3D viscous flow computations on unstructured grids, 36th AIAA Aerospace Sciences Meeting and Exhibit, 12–15 January 1998, Reno, Nevada, US. AIAA No. 1998-0614, 1998. doi:10.2514/6.1998-614.
- 23. EDWARDS, J.R., FRANKLIN, R.K. and LIOU, M.S. Low-diffusion flux-splitting methods for real fluid flows with phase transitions, *AIAA J*, 2000, **38**, (9), pp 1624-1633.
- 24. MENTER, F.R. Two-equation eddy-viscosity turbulence models for engineering applications, *AIAA J*, 1994, **32**, (8), pp 1598-1605.
- 25. WANG, B., ZHAO, Q.J. and XU, G.H. A new moving-embedded grid method for numerical simulation of unsteady flowfield of the helicopter rotor in forward flight, *Acta Aerodynamica Sinica*, 2012, **30**, (1), pp 14-21. (in Chinese)
- DONOVAN, J.F., KRAL, L.D. and CARY, A.W. Active flow control applied to an airfoil, 36th AIAA Aerospace Sciences Meeting and Exhibit, 12–15 January 1998, Reno, Nevada, US. AIAA No. 1998-16119, 1998. doi:10.2514/6.1998-210.
- CARADONNA, T.X., LAUB, G.H. and TUNG, C. An experimental investigation of the parallel blade-vortex interactions, Netherlands Association of Aeronautical Engineers and Technische Hogeschool te Delft, 10th European Rotorcraft Forum, 28–31 August, 1984, The Hague, Netherlands. NASA No. TM 86005, 1984.
- SHEFFER, S.G., ALONSO, J.J., MARTINELLI, L. and JAMESON, A. Time-accurate simulation of helicopter rotor flows including aeroelastic effects, 35th Aerospace Sciences Meeting and Exhibit, 6–9 January 1997, Reno, Nevada, US. AIAA No. 1997-0399, 1997. doi:10.2514/6.1997-399.